

Dynamics of Superrotations and Supertranslations in 2+1 Dimensions

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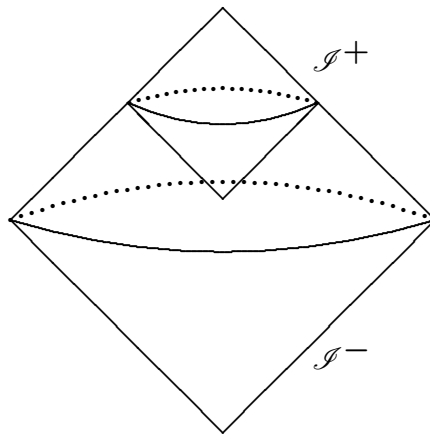
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The setting



- Asymptotically flat spacetime has large symmetry group at \mathcal{I}^\pm
- Strominger et al.: BMS group is spontaneously broken
 \Rightarrow Goldstone modes? Degenerate vacuum?
- This *might* help with black hole information loss problem

Can one extract a dynamical description of the new degrees of freedom?

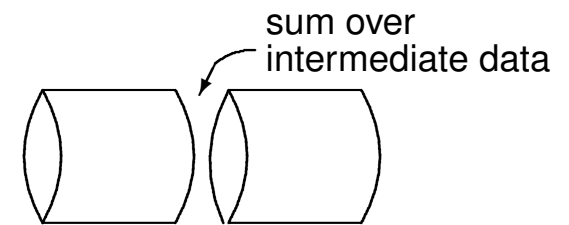
The idea in a nutshell

- Action on manifold with boundary has two pieces:

$$I = I_{bulk} + I_{bdry} \quad \text{with} \quad I_{bulk} = \int \mathcal{L} d^n x$$

Boundary piece needed

- Classically: to allow extrema
- Quantum mechanically: to ensure proper “sewing” of path integrals

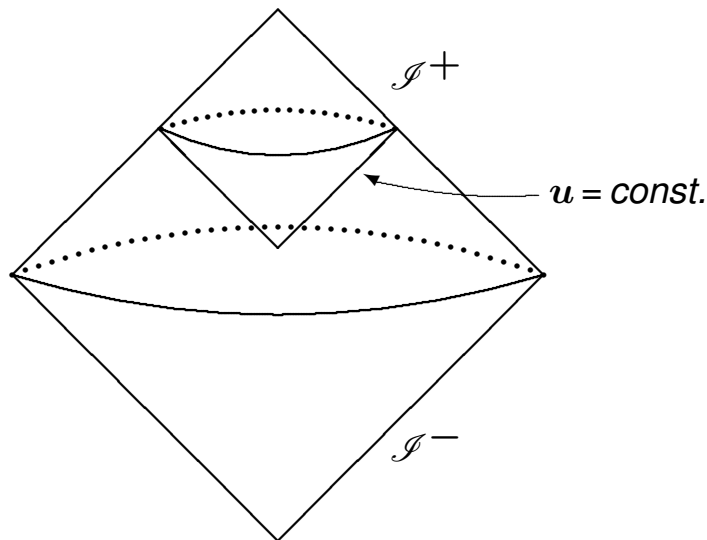


- Gauge symmetries of I_{bulk} will typically be broken by I_{bdry}
 - Formerly nonphysical degrees of freedom become dynamical at boundary
- Action for new degrees of freedom is induced from I_{bdry}
- Results for (2+1)-dimensional gravity:
 - Asymptotically AdS: Liouville action
 - Asymptotically flat: relation to chiral Liouville theory

$(2 + 1)$ -dimensional asymptotically flat gravity

Partially gauge-fix metric: “Bondi” coordinates

$$ds^2 = -2dudr + g_{uu} du^2 + 2g_{u\phi} dud\phi + r^2 e^{2\varphi} d\phi^2$$



$u = \text{const.}$: outgoing null surfaces

r is an affine parameter along null geodesics $u = \text{const.}$, $\phi = \text{const.}$

Behavior of metric near \mathcal{I}^\pm

Symmetries: supertranslations $\sim u \rightarrow u + e^\varphi T(\phi)$

superrotations $\sim \phi \rightarrow \phi + Y(\phi)$

From field equations (Barnich & Troessaert)

$$g_{uu} = -2r\partial_u\varphi + e^{-2\varphi} \left[-(\partial_\phi\varphi)^2 + 2\partial_\phi^2\varphi + \Theta \right]$$

$$g_{u\phi} = e^{-\varphi} \left[\Xi + \int^u d\tilde{u} \left\{ \frac{1}{2}\partial_\phi\Theta - \partial_\phi\varphi[\Theta - (\partial_\phi\varphi)^2 + 3\partial_\phi^2\varphi] + \partial_\phi^3\varphi \right\} \right]$$

$$\text{with } \partial_u\Theta = \partial_u\Xi = 0$$

Θ and Ξ are charges for supertranslations and superrotations

$$\mathcal{Q} \sim \frac{1}{16\pi G} \int d\phi (\Theta T + \Xi Y)$$

Boundary terms

$$\delta I_{grav} = \text{bulk piece} \\ + \frac{1}{16\pi G} \int_{\partial M} d^2x [-\partial_r(r e^\varphi \delta g_{uu}) - 2g_{uu} \partial_r(r e^\varphi \delta \varphi)] + \mathcal{O}\left(\frac{1}{r}\right)$$

Use asymptotic expression to integrate: find

$$\delta I = -\frac{1}{8\pi G} \delta \int_{\partial M} d^2x e^{\tilde{\varphi}} \tilde{g}_{uu} + \frac{1}{16\pi G} \int_{\partial M} d^2x e^{\tilde{\varphi}} \delta \tilde{g}_{uu}$$

where \tilde{X} means the $\mathcal{O}(1)$ part of X

Fix either \tilde{g}_{uu} (i.e., Θ) or $e^{\tilde{\varphi}}$ (or some combination)

(Same result from looking directly at canonically conjugate variables)

Fix $\tilde{\varphi}$:

$$I_{bdry} = \frac{1}{16\pi G} \int_{\partial M} d^2x e^{\tilde{\varphi}} \tilde{g}_{uu}$$

Not quite extrinsic curvature of $r = \text{const.}$:

- unit normal n^a to $r = \text{const.}$
- induced metric q_{ab} on $r = \text{const.}$
- null vector $\ell \propto du$, normalized so $n_a \ell^a = 1$

$$I_{bdry} = \frac{1}{16\pi G} \int_{\partial M} d^2x \sqrt{q} \nabla_a \ell^a$$

Relation to diffeomorphisms

Start with flat base metric

$$ds^2 = -2d\bar{u}d\bar{r} + d\bar{u}^2 + \bar{r}^2 d\bar{\phi}^2$$

Diffeomorphism

$$\bar{u} = u_0 + \frac{u_1}{r} + \dots, \quad \bar{\phi} = \phi_0 + \frac{\phi_1}{r} + \dots, \quad \bar{r} = ar + b_0 + \frac{b_1}{r} + \dots$$

Form invariance of metric \Rightarrow relations among coefficients u_0, ϕ_0, b_0

Confirm form $g_{uu} = -2r\partial_u\varphi + e^{-2\varphi} \left[-(\partial_\phi\varphi)^2 + 2\partial_\phi^2\varphi + \Theta \right]$

with $e^\varphi = \frac{\partial_\phi\phi_0}{\partial_u u_0}, \quad \Theta = -(\partial_\phi\phi_0)^2 - 2\{\phi_0; \phi\}$

Schwarzian derivative $\{f; z\} = \frac{f'''}{f'} - \frac{3}{2} \left(\frac{f''}{f'} \right)^2$

An action for supertranslations and superrotations

Boundary action (for fixed φ) is then

$$I_{bdry} = \frac{1}{16\pi G} \int_{\partial M} d^2x e^{-\varphi} \left[-(\partial_\phi \varphi)^2 + 2\partial_\phi^2 \varphi - (\partial_\phi \phi_0)^2 - 2\{\phi_0; \phi\} \right]$$

What can we say about this action?

- Depends on ϕ_0 (superrotation parameter)
- But $e^\varphi = (\partial_\phi \phi_0)/(\partial_u u_0)$ fixed,
so equivalently depends on u_0 (supertranslation parameter)
- Schwarzian derivative in $\Theta \Leftrightarrow \Theta$ transforms as stress-energy tensor
- Many connections to Liouville theory and Virasoro group

- Vary φ : Hill's equation

$$\partial_\phi^2 \chi + \frac{6}{c} T \chi = 0 \quad \text{with } \chi = \frac{1}{\sqrt{2\pi G}} e^{-\varphi/2}, \quad T = \frac{c}{24} (\partial_\phi \phi_0)^2 + \frac{c}{12} \{\phi_0; \phi\}$$

– Solutions χ_1, χ_2 give orbits of Virasoro group:

$$- \delta_\varepsilon T = -\frac{c}{12} \varepsilon''' - 2\varepsilon' T - \varepsilon T' = 0 \quad \text{with } \varepsilon = \chi_1^2, \chi_1 \chi_2, \chi_2^2$$

- Action for ϕ_0 related to Alexe'ev-Shatashvili action for coadjoint quantization

- φ fixed \Rightarrow either ϕ_0 or u_0 can be treated as independent:

$$- \partial_\phi \phi_0 = e^\sigma \Rightarrow \text{chiral Liouville action for } \sigma$$

$$- \partial_u u_0 = e^{\tilde{\sigma}} \Rightarrow \text{chiral Liouville action for } \tilde{\sigma}$$

Clearly a strong connection to CFT (or BMS/Galilean theory?)

Generalization to 3+1 dimensions?