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# Interaction between bosonic dark matter and stars

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R.Brito, V. Cardoso & H. Okawa, Phys.Rev.Lett. 115 (2015) no.11, 111301 arXiv: 1508.04773  
R. Brito, V. Cardoso, H. Okawa, C. Macedo & C.Palenzuela, Phys.Rev. D93 (2016) no.4, 044045,  
arXiv: 1512.00466



More info at <http://blackholes.ist.utl.pt>

**FCT**

Fundação para a Ciência e a Tecnologia  
MINISTÉRIO DA CIÊNCIA, INOVAÇÃO E DO ENSINO SUPERIOR

# Gravity & fundamental (bosonic) fields

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## **The never ending search for dark matter...**

Signatures of dark matter in gravitating systems?

Can we use it to constrain (or even detect) dark matter candidates?

Explore the rich phenomenology of fundamentals fields within full General Relativity.

$$S = \int d^4x \sqrt{-g} \left( \frac{R}{\kappa} - \frac{1}{4} F^{\mu\nu} \bar{F}_{\mu\nu} - \frac{\mu_V^2}{2} A_\nu \bar{A}^\nu - \frac{1}{2} g^{\mu\nu} \Phi_{,\mu}^* \Phi_{,\nu} - \frac{\mu_S^2 |\Phi|^2}{2} + \mathcal{L}_{\text{matter}} \right)$$

# Bosonic stars

D.J. Kaup '68; Ruffini & Bonazzola '69; Seidel & Suen '91;  
 Brito, Cardoso, Herdeiro & Radu '15; Brito, Cardoso & Okawa '15

$$T_{\text{scalar}}^{\mu\nu} = -\frac{1}{4}g^{\mu\nu} (\Phi_{,\alpha}^* \Phi^{,\alpha} + \mu_S^2 \Phi^* \Phi) + \frac{1}{4} (\Phi^{*,\mu} \Phi^{,\nu} + \Phi^{,\mu} \Phi^{*,\nu})$$

$$T_{\text{vector}}^{\mu\nu} = F_{\alpha}^{(\mu} \bar{F}^{\nu)\alpha} - \frac{1}{4} \bar{F}^{\alpha\beta} F_{\alpha\beta} g^{\mu\nu} - \frac{1}{2} \mu_V^2 A_{\alpha} \bar{A}^{\alpha} g^{\mu\nu} + \mu_V^2 A^{(\mu} A^{\nu)}$$

**Spherical symmetry:**  $ds^2 = -F(t, r)dt^2 + B(t, r)dr^2 + r^2 d\Omega^2$

**Complex fields - Boson stars:**

$$\begin{aligned} N^i(t, r) &= N(r) \\ \Phi(t, r) &= \phi(r) e^{-i\omega t} \end{aligned}$$

**Real fields - Oscillatons:**

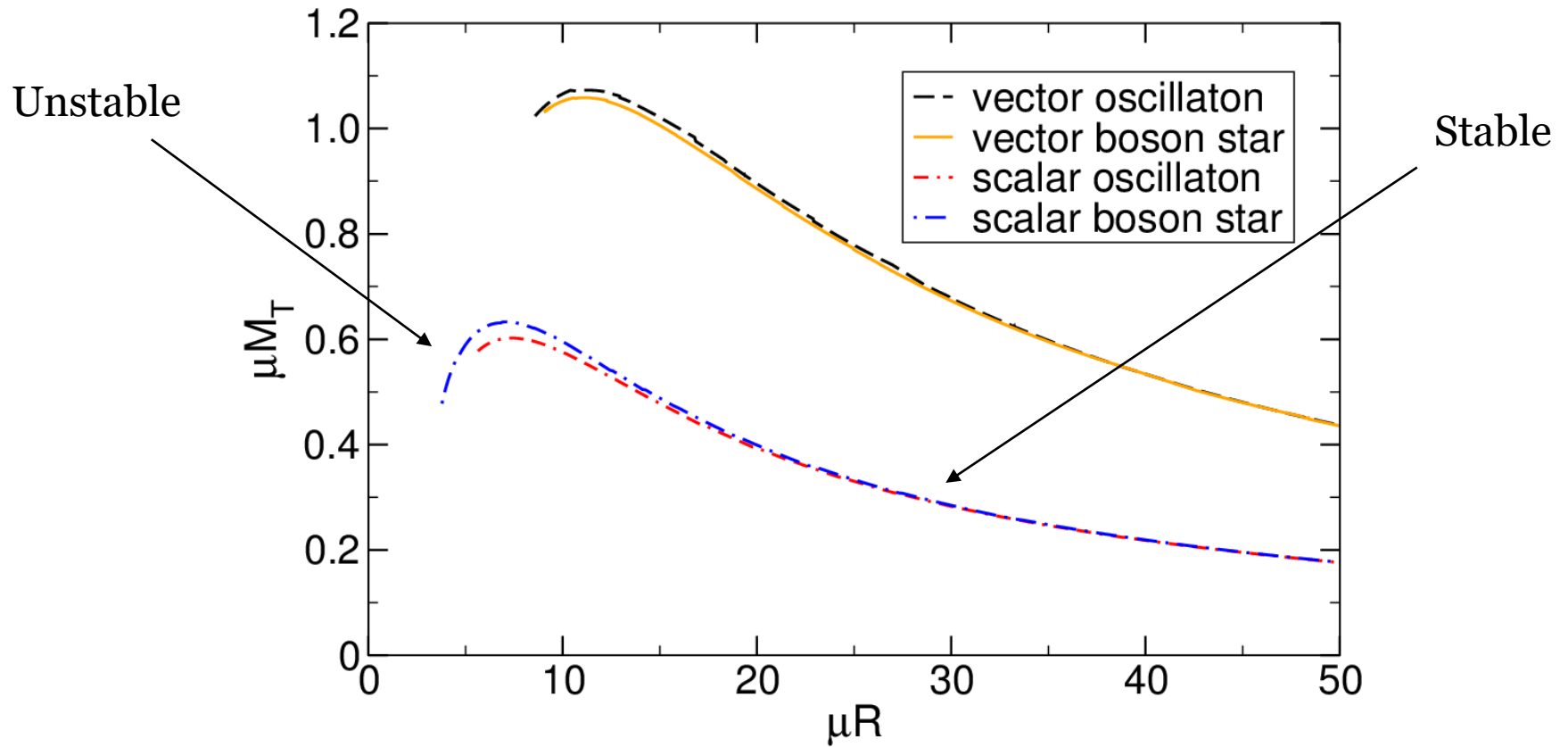
$$\begin{aligned} N^i(t, r) &= \sum_{j=0}^{\infty} N_{2j}^i(r) \cos(2j\omega t) \\ \Phi(t, r) &= \sum_{j=0}^{\infty} \phi_{2j+1}(r) \cos[(2j+1)\omega t] \end{aligned}$$

$$N^i = (B, F)$$

**For rotating boson/Proca stars see:**  
 Yoshida, Eriguchi '97; Schunck, Mielke '98  
 Brito, Cardoso, Herdeiro & Radu '15

# Oscillatons and Boson stars

D.J. Kaup '68; Ruffini & Bonazzola '69; Seidel & Suen '91;  
Brito, Cardoso, Herdeiro & Radu '15; Brito, Cardoso & Okawa '15



$$\frac{M_{\max}}{M_{\odot}} = 8 \times 10^{-11} \left( \frac{\text{eV}}{m_B c^2} \right)$$

$$\omega \sim \mu \implies f = 2.5 \times 10^{14} \left( \frac{m_B c^2}{\text{eV}} \right) \text{ Hz}$$

# Accretion onto stars: boson-fluid stars

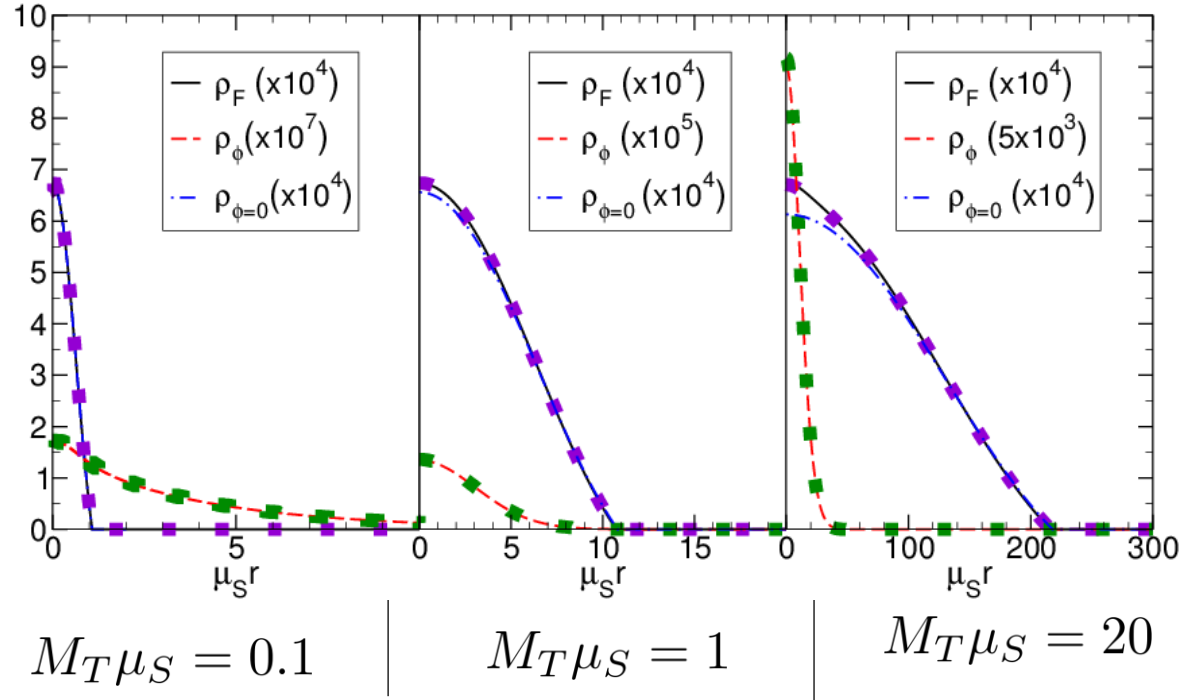
A. Henriques, A. R. Liddle & R. Moorhouse '89;  
 Brito, Cardoso & Okawa '15; Brito, Cardoso, Macedo, Okawa & Palenzuela, '16

Do Einstein's equations allow for stable solutions describing a star with a bosonic core?

**Perfect fluid star:**  $T_{\text{fluid}}^{\mu\nu} = (\rho_F + P) u^\mu u^\nu + P g^{\mu\nu}$

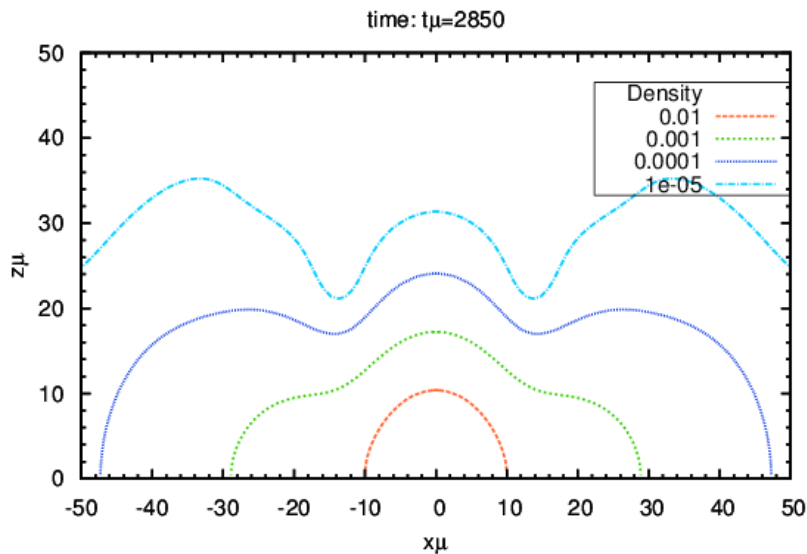
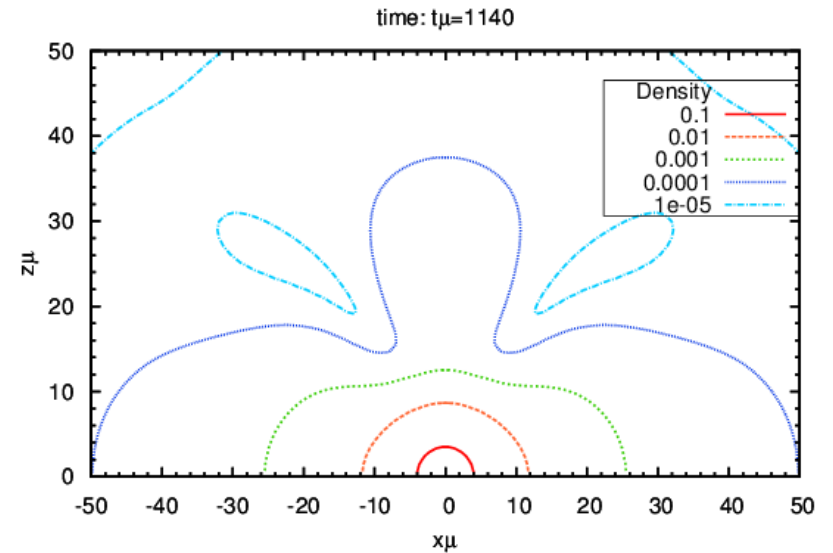
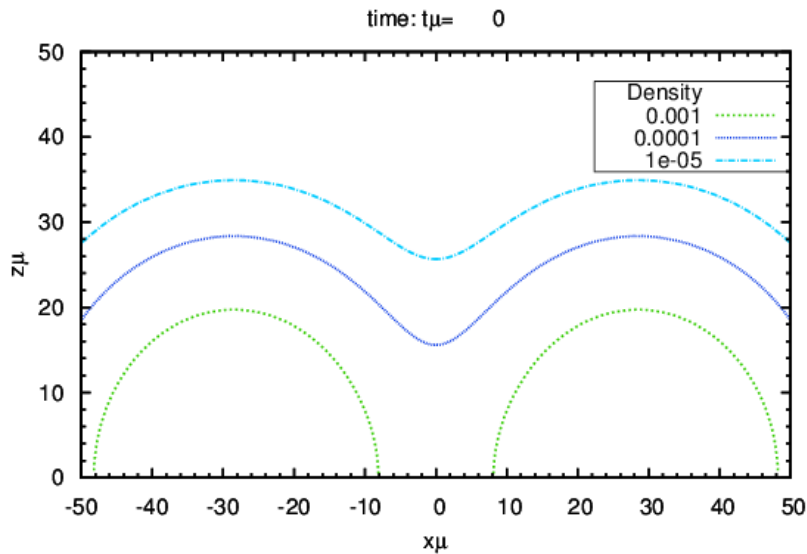
For real bosonic fields, equations imply that the star's material must oscillate:

$$\rho_F = \sum_{j=0}^{\infty} \rho_{F2j}(r) \cos(2j\omega t)$$



# Growth of bosonic structures

Brito, Cardoso & Okawa '15; Brito, Cardoso, Macedo, Okawa & Palenzuela, '16

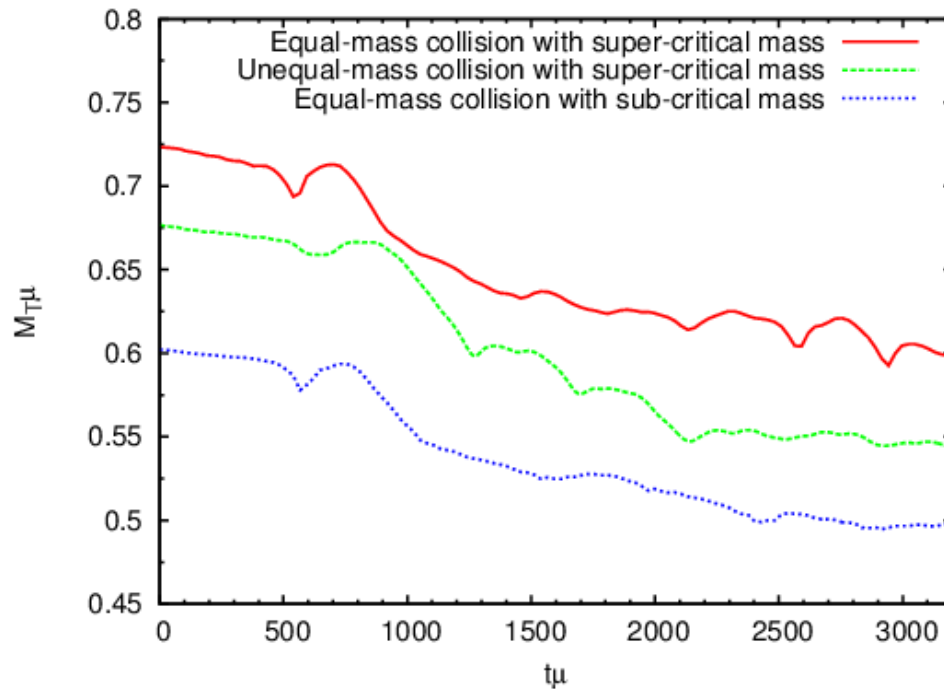


$$M_{\text{initial}} \sim 0.3/\mu_S, R_{\text{initial}} \sim 20/\mu_S$$

Mass of final object larger than initial  
mass of each oscillaton  $M\mu \sim 0.5$ .

# Collision of bosonic structures

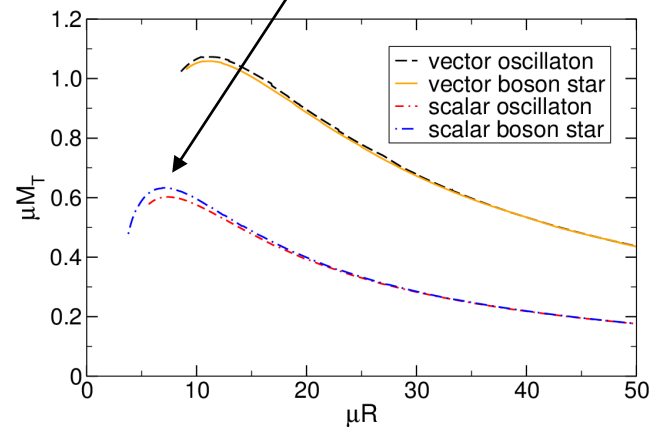
Brito, Cardoso & Okawa '15; Brito, Cardoso, Macedo, Okawa & Palenzuela, '16



Collapse to a black hole for  $M_\mu > 0.6$  can be avoided due to a “gravitational cooling” mechanism.

Bosonic structures can grow through mergers.

Growth continues till the threshold mass  $M_\mu \sim 0.6$ .



# Accretion onto stars

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- ❖ Accumulation stage and thermalization.

- ❖ DM core collapses, after becoming self-gravitating or when the DM core reaches the threshold  $M\mu \sim 0.6$ .

(Goldman and Nussinov PRD40, 3221 (1989); Bertone and Fairbairn PRD77, 043515 (2008); Bramante, PRL115, 141301 (2015); Kurita and Nakano, arXiv:1510.00893...etc)

- ❖ Lack of rigorous support for this picture.

- ❖ For Compton wavelengths smaller than size of star, bosonic core behaves as isolated oscillaton.

- ❖ We just showed cases where the core does *not* collapse to a black hole when  $M\mu > 0.6$ .

- ❖ Stable configurations with self-gravitating DM cores can be constructed...



Thank you

# Backup Slides

# Accretion onto stars

A. Henriques, A. R. Liddle & R. Moorhouse '89;

Brito, Cardoso & Okawa '15; Brito, Cardoso, Macedo, Okawa & Palenzuela, arXiv: 1512.00466

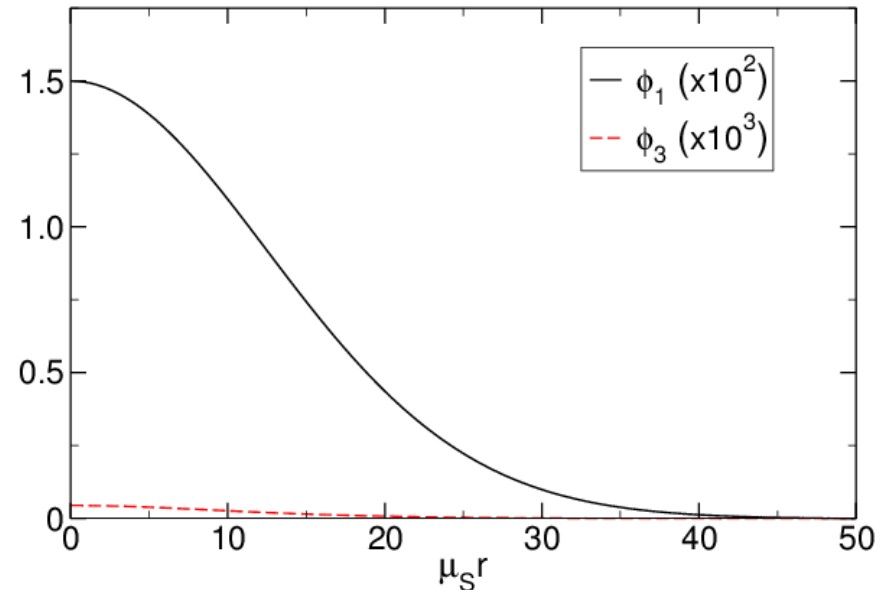
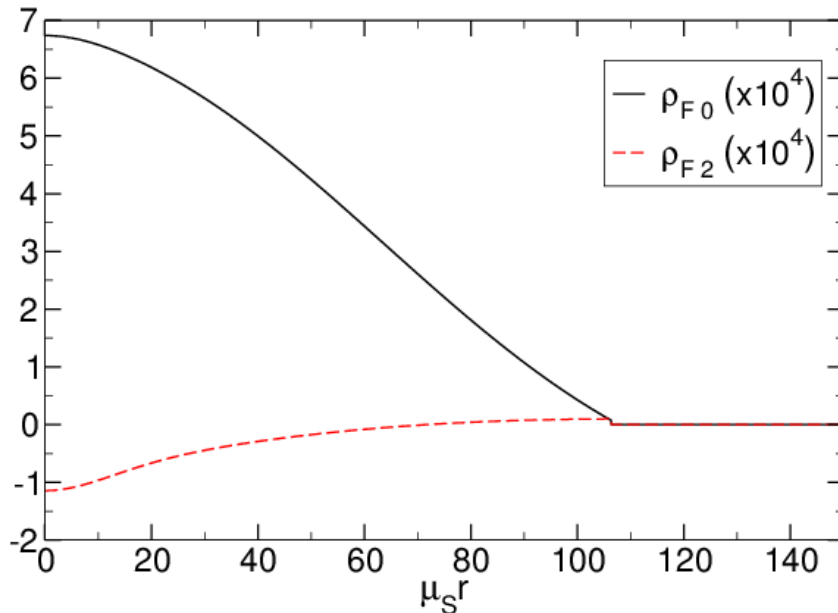
**Perfect fluid star:**  $T_{\text{fluid}}^{\mu\nu} = (\rho_F + P) u^\mu u^\nu + P g^{\mu\nu}$

$$u^\mu = \frac{\Gamma}{\sqrt{-g_{tt}}} (1, V(t, r), 0, 0)$$

$$N^i = \sum_{j=0}^{\infty} N_{2j}^i(r) \cos(2j\omega t)$$

$$N^i = (g_{tt}, g_{rr}, n_F, \rho_F, P, V)$$

$$\Phi(t, r) = \sum_{j=0}^{\infty} \phi_{2j+1}(r) \cos[(2j+1)\omega t]$$

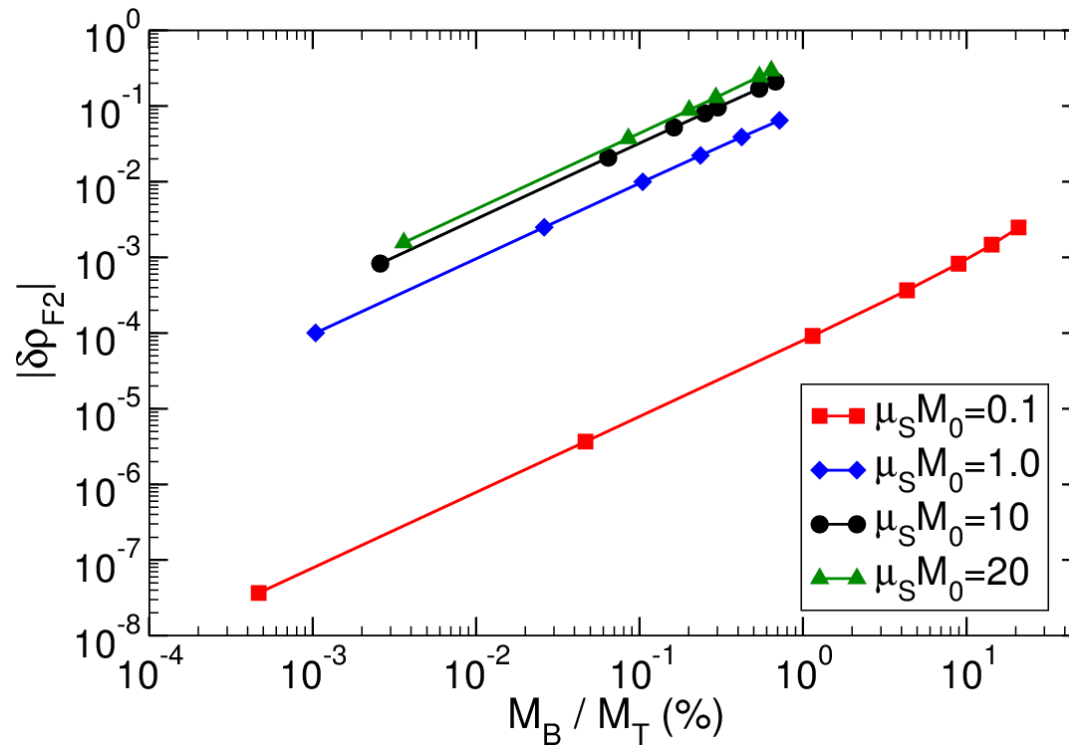


# Fluid oscillations

Brito, Cardoso & Okawa '15;

Brito, Cardoso, Macedo, Okawa & Palenzuela, '15

Equations imply that the star material must oscillate:  $\rho_F = \sum_{j=0}^{\infty} \rho_{F2j}(r) \cos(2j\omega t)$

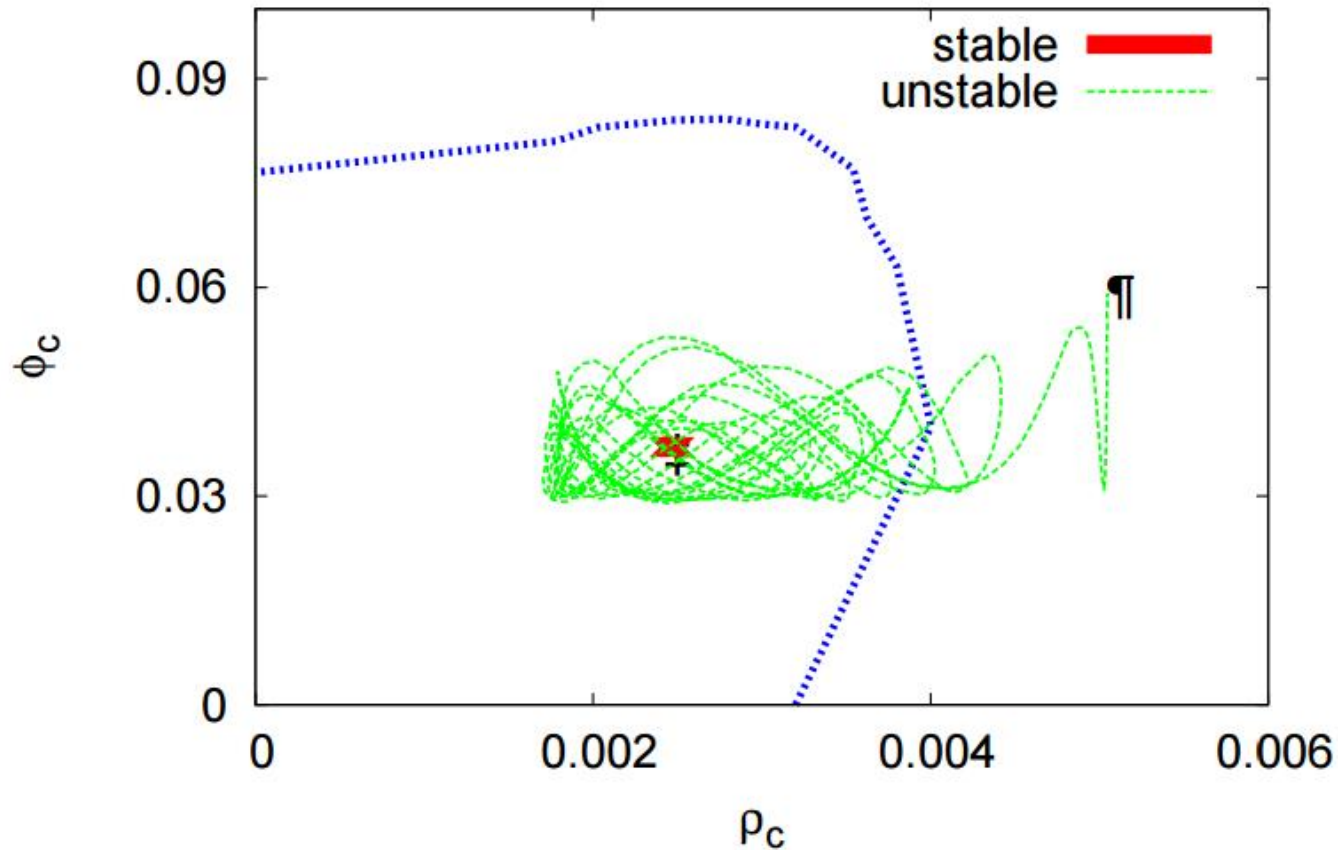


$$\omega \sim \mu_S \implies f = 2.5 \times 10^{14} \left( \frac{m_B c^2}{eV} \right) \text{ Hz}$$

$$\mu_S^{-1} = \lambda_c \rightarrow \text{Compton wavelength}$$

# Stability?

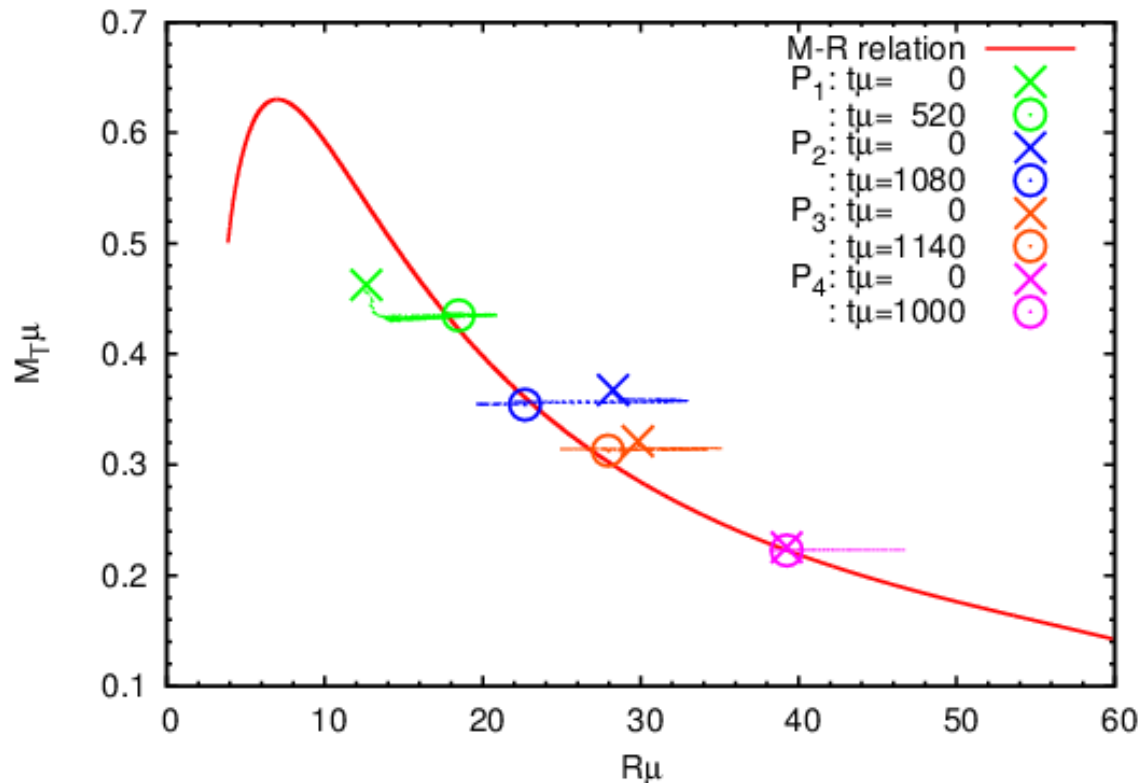
A. Henriques, A. R. Liddle & R. Moorhouse '90 ; Valdez-Alvarado, Palenzuela, Alic & Ureña-López '13;  
Brito, Cardoso, Macedo, Okawa & Palenzuela, '16



From: Valdez-Alvarado, Palenzuela, Alic & Ureña-López, Phys.Rev. D87 (2013) 8, 084040

**Stability?** For sufficiently small scalar composites (or vice-versa) stability analysis of the host star is still valid. (A. Henriques, A. R. Liddle & R. Moorhouse PLB B251, 511 (1990))

# Do they ever form?



Purely bosonic states do.

Two channels for composite fluid-boson stars:

- Gravitational collapse in a bosonic environment;
- Capture and accretion of DM into the core of compact stars. Collapse to a black hole?

# Accretion onto stars

A. Henriques, A. R. Liddle & R. Moorhouse '89;  
 Brito, Cardoso & Okawa '15; Brito, Cardoso, Macedo, Okawa & Palenzuela, '16

**Perfect fluid star:**  $T_{\text{fluid}}^{\mu\nu} = (\rho_F + P) u^\mu u^\nu + P g^{\mu\nu}$

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$$\Phi(t, r) = \sum_{j=0}^{\infty} \phi_{2j+1}(r) \cos[(2j+1)\omega t]$$

