



UNIVERSITY OF TRENTO - Italy



Trento Institute for
Fundamental Physics
and Applications

Short gamma-ray bursts from binary neutron star mergers forming long-lived neutron stars

RICCARDO CIOLFI

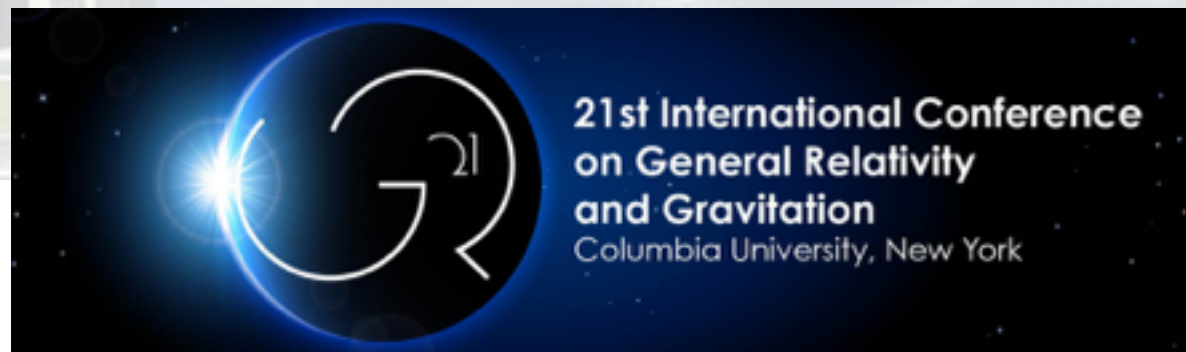
Ciolfi R. & Siegel D.M. (2015a) ApJ Letters 798, L36

Ciolfi R. & Siegel D.M. (2015b), in “*Swift: 10 Years of Discovery*”, PoS(SWIFT 10)108

Siegel D.M. & Ciolfi R. (2016a), ApJ 819, 14

Siegel D.M. & Ciolfi R. (2016b), ApJ 819, 15

Ciolfi R. (2016), submitted to ApJ Letters, ArXiv:1606.01743



21st International Conference
on General Relativity
and Gravitation
Columbia University, New York

13th July 2016

 COLUMBIA UNIVERSITY
IN THE CITY OF NEW YORK

Short Gamma-Ray Bursts

GAMMA-RAY BURSTS

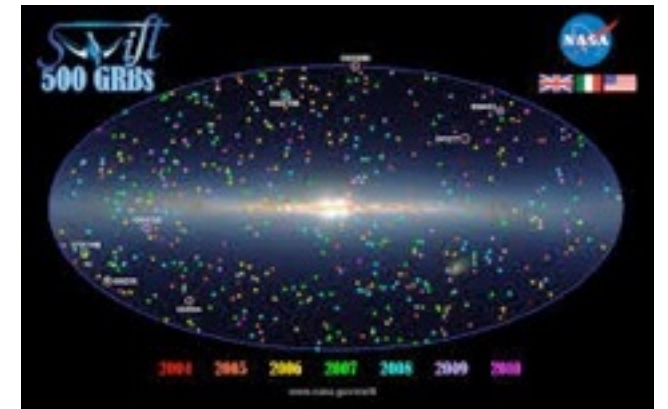
- **bright gamma-ray flashes** with extragalactic origin, **huge luminosities!**
- followed by **afterglows** in **X-ray, optical** and **radio band**
- divided into **long** and **short** (prompt burst **longer/shorter than 2 s**)



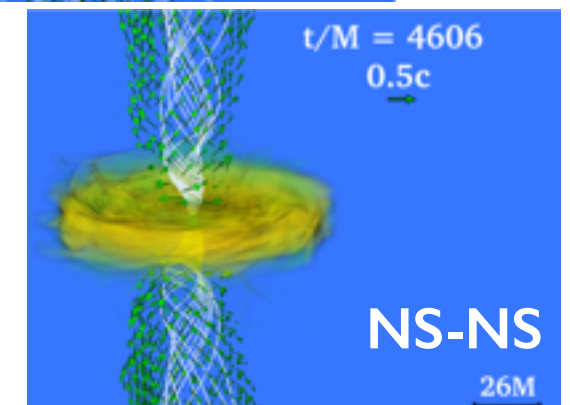
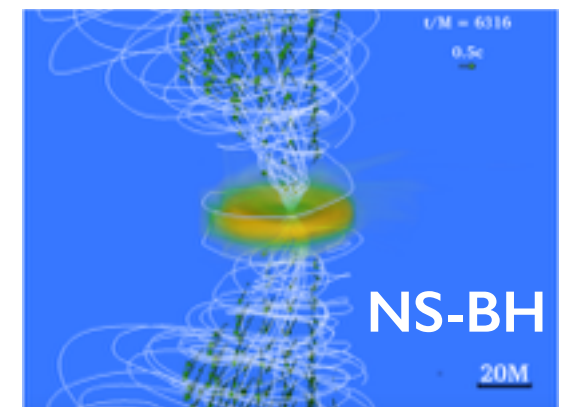
“standard” model of SGRBs

central engine is a **black hole**
surrounded by a massive **accretion**
torus → end result of a
BNS or NS-BH binary merger

Paczynski 1986, Eichler et al. 1989, Narayan et al. 1992,
Barthelmy et al. 2005, Fox et al. 2005, Gehrels et al. 2005, ...



Paschalidis et al. 2015



Ruiz et al. 2016

X-ray afterglows of SGRBs

- SWIFT revealed that most SGRBs are accompanied by long-duration ($\sim 10^2 - 10^5$ s) and high-luminosity ($10^{46} - 10^{51}$ erg/s) X-ray afterglows
- total energy can be higher than the SGRB itself
- hardly produced by BH-torus system - they suggest ongoing energy injection from a **long-lived NS**

MAGNETAR MODEL

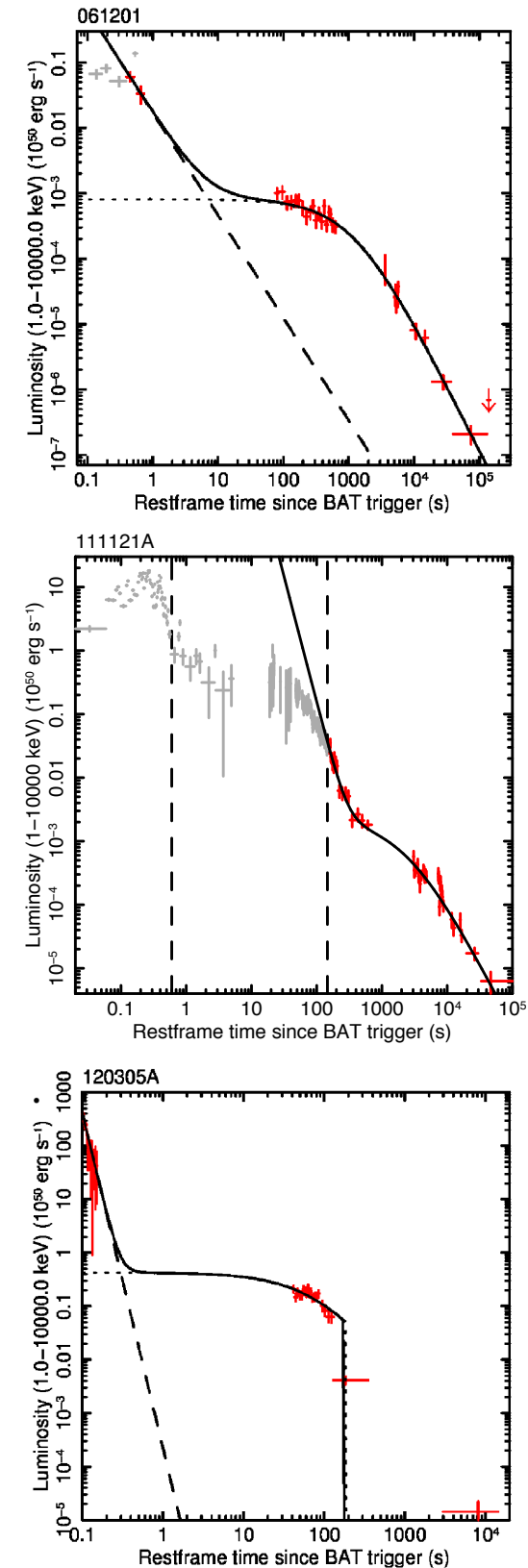
Zhang & Meszaros 2001

Metzger et al. 2008

X-ray emission \rightarrow spindown of a **uniformly rotating NS** with a strong surface magnetic field
 $\gtrsim 10^{14} - 10^{15}$ G

**dipole
spindown**

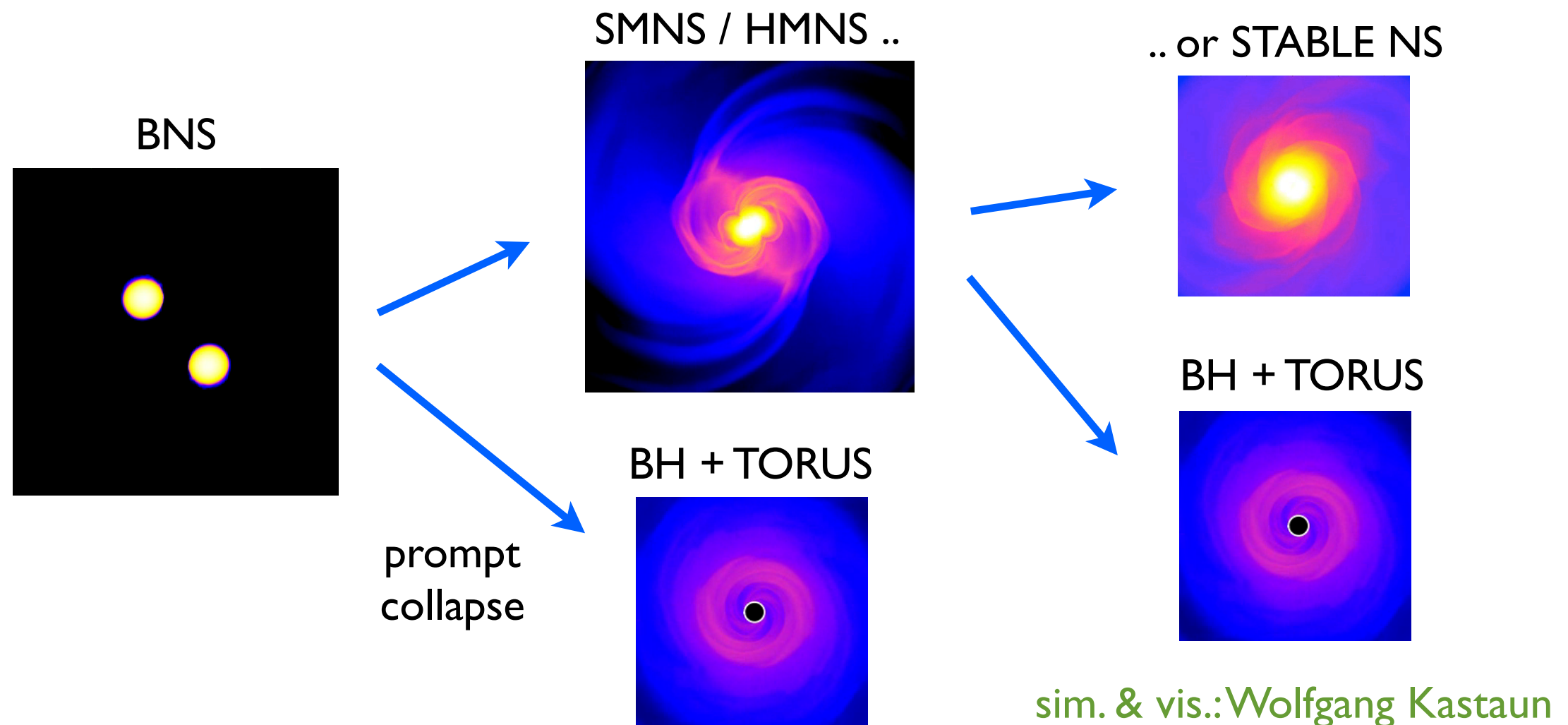
$$L_{\text{sd}}(t) \sim B^2 R^6 \Omega_0^4 \left(1 + \frac{t}{t_{\text{sd}}}\right)^{-2}$$



Gompertz et al. 2013

Rowlinson et al. 2013

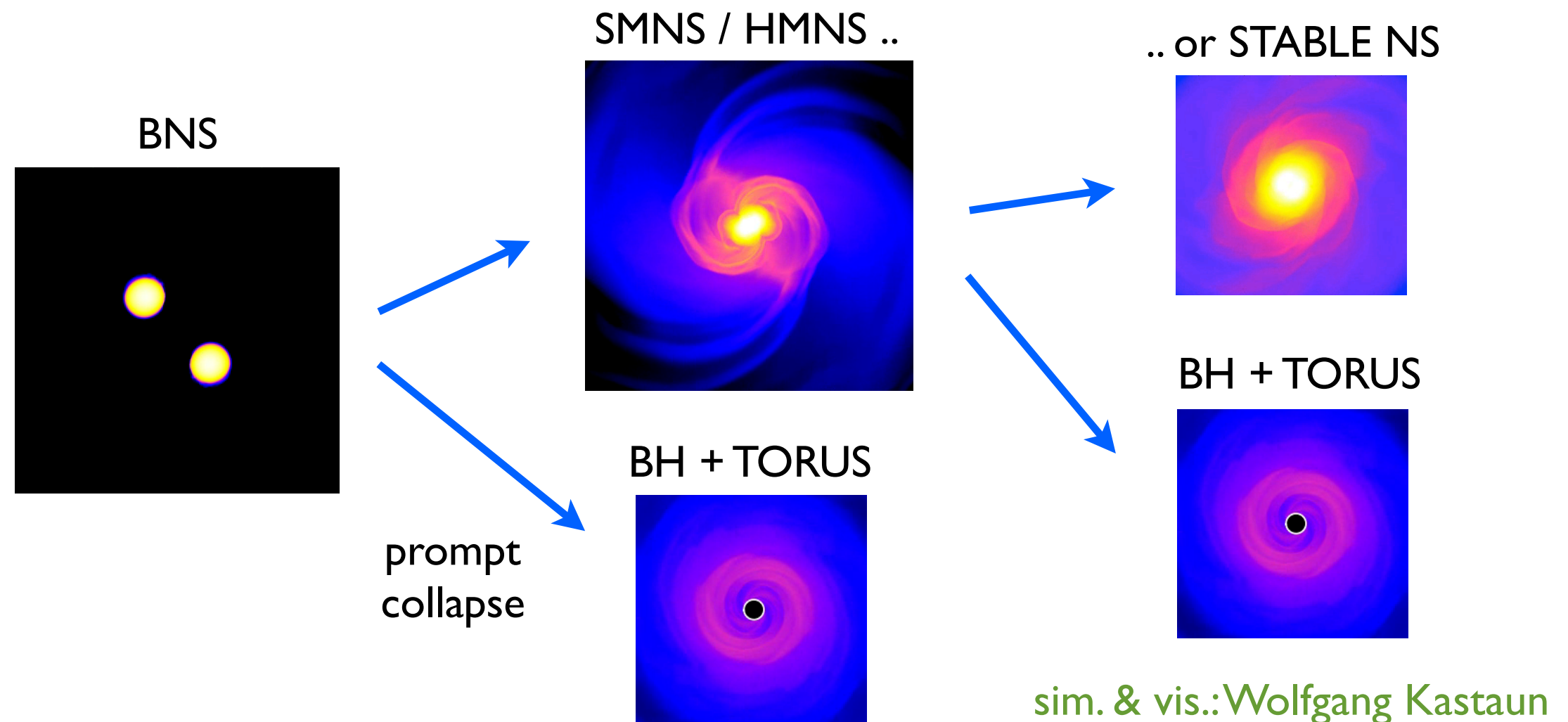
Product of BNS mergers



LONG-LIVED NS IS A VERY LIKELY OUTCOME OF THE MERGER

- observation of $\sim 2 M_{\odot}$ NSs Demorest et al. 2010
Antoniadis et al. 2013
- progenitor masses peak around $1.3 - 1.4 M_{\odot} \rightarrow$ BMP mass likely $< 2.5 M_{\odot}$ Belczynski et al. 2008
- stable NS obtained in GR BNS merger simulations Giacomazzo & Perna 2013

Product of BNS mergers



PROBLEM OF THE LONG-LIVED NS MODEL :

strong baryon pollution can choke the
formation of a relativistic jet

→ **HARD TO EXPLAIN THE SGRB PROMPT EMISSION**

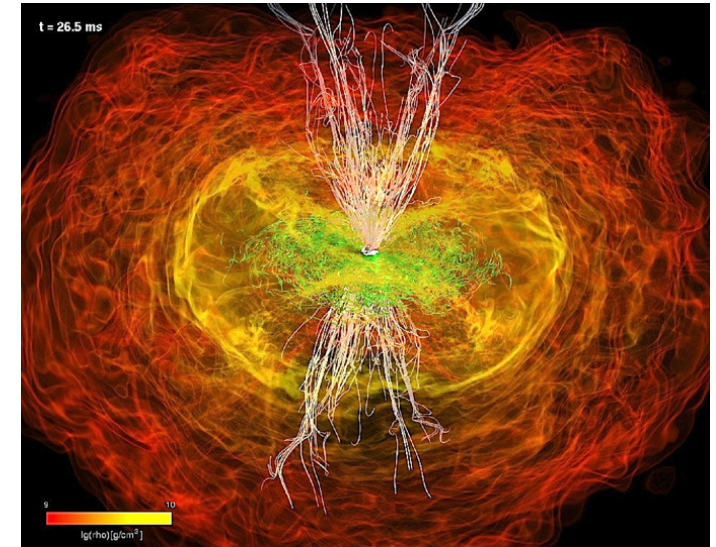
e.g., Dessart et al. 2009, Hotokezaka et al. 2013, Siegel et al. 2014

The SGRB dichotomy

- Numerical relativity picture: prompt BH-torus formation

→ can explain prompt SGRB emission ✓

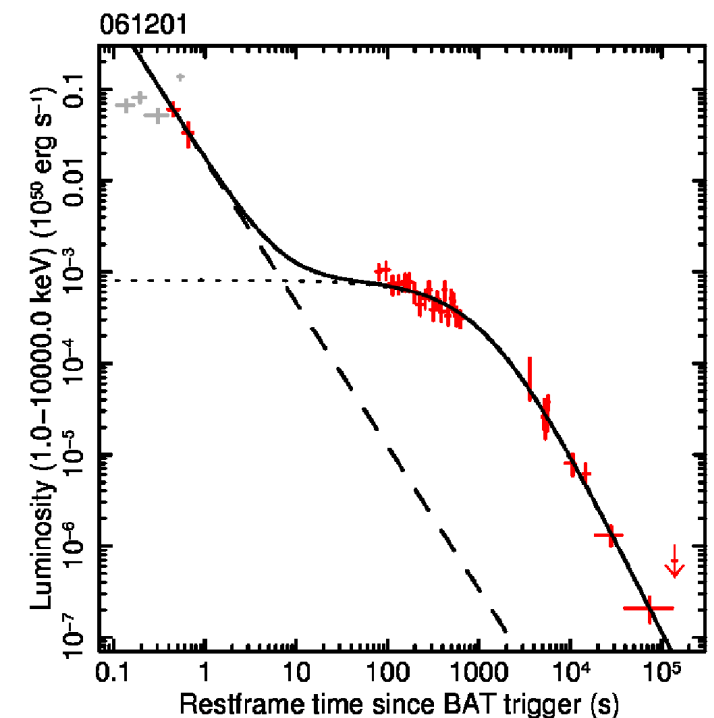
→ cannot explain X-ray afterglows ✗



- Observational picture: magnetar model

→ cannot explain prompt SGRB emission ✗

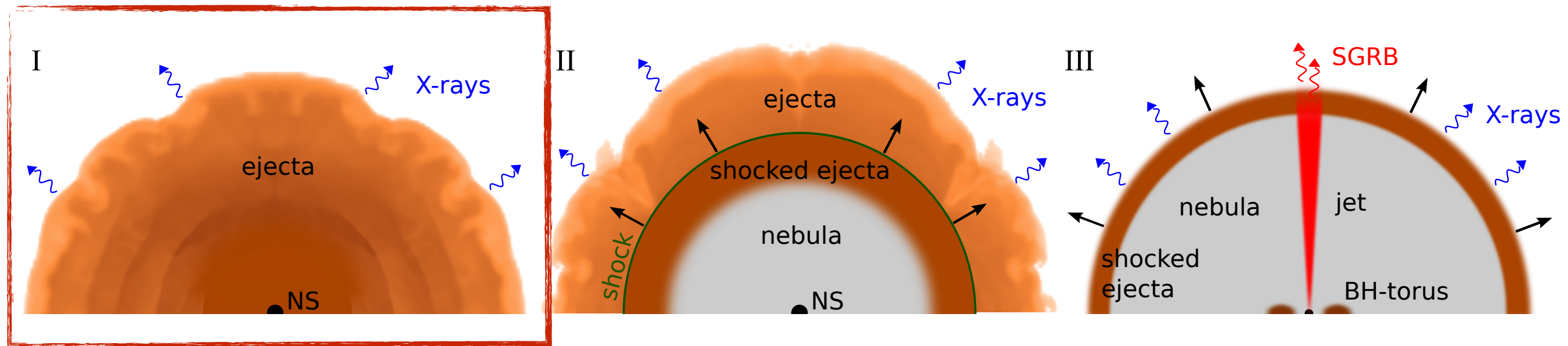
→ can explain X-ray afterglows ✓



..possible alternative solutions?

“Time-reversal” phenomenology

Cioffi & Siegel 2015a, ApJ Letters 798, L36



- ➡ (I) The differentially rotating, supramassive NS (SMNS) ejects a baryon-loaded and highly isotropic wind

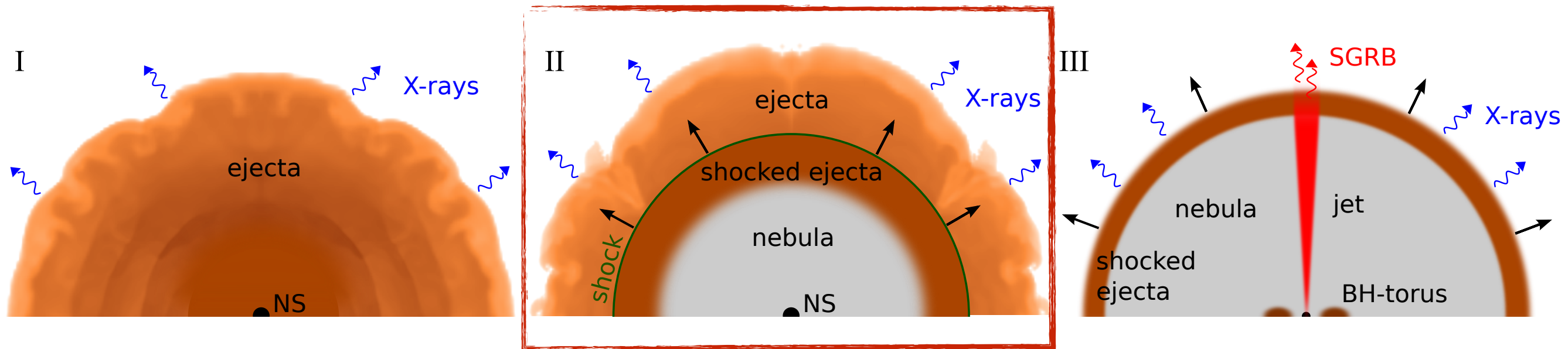
Siegel et al. 2014

(see also Siegel & Cioffi 2015)

Dessart et al. 2009

“Time-reversal” phenomenology

Cioffi & Siegel 2015a, ApJ Letters 798, L36

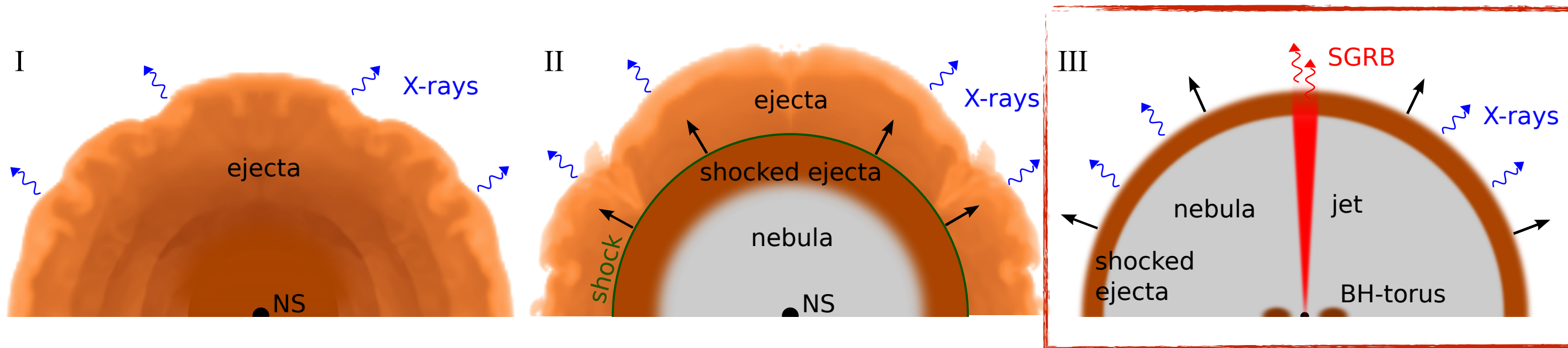


(I) The differentially rotating, supramassive NS (SMNS) ejects a baryon-loaded and highly isotropic wind

➡ (II) The cooled-down and uniformly rotating NS emits spin-down radiation inflating a photon-pair nebula that drives a shock through the ejecta

“Time-reversal” phenomenology

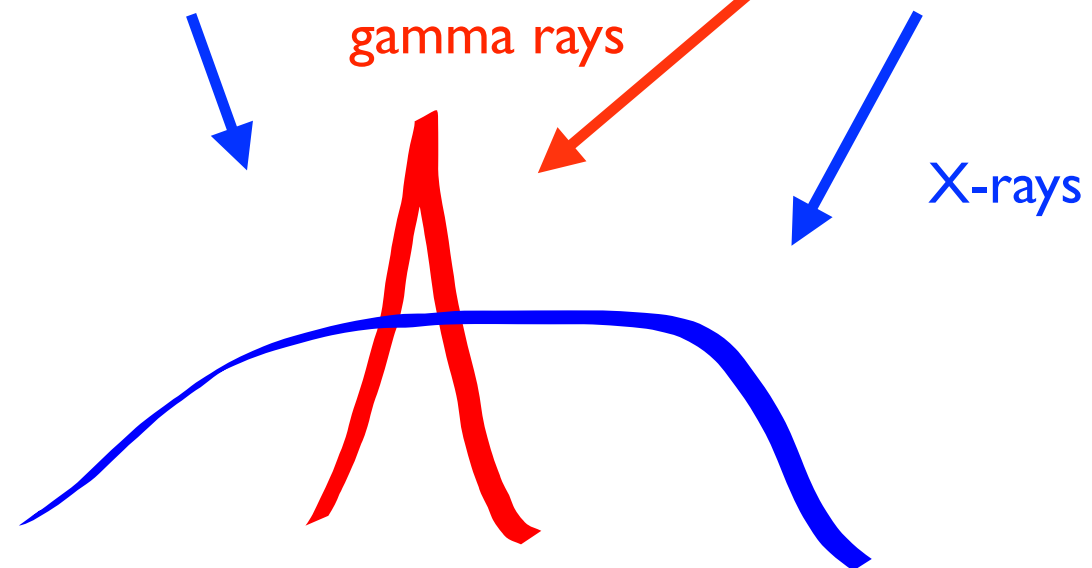
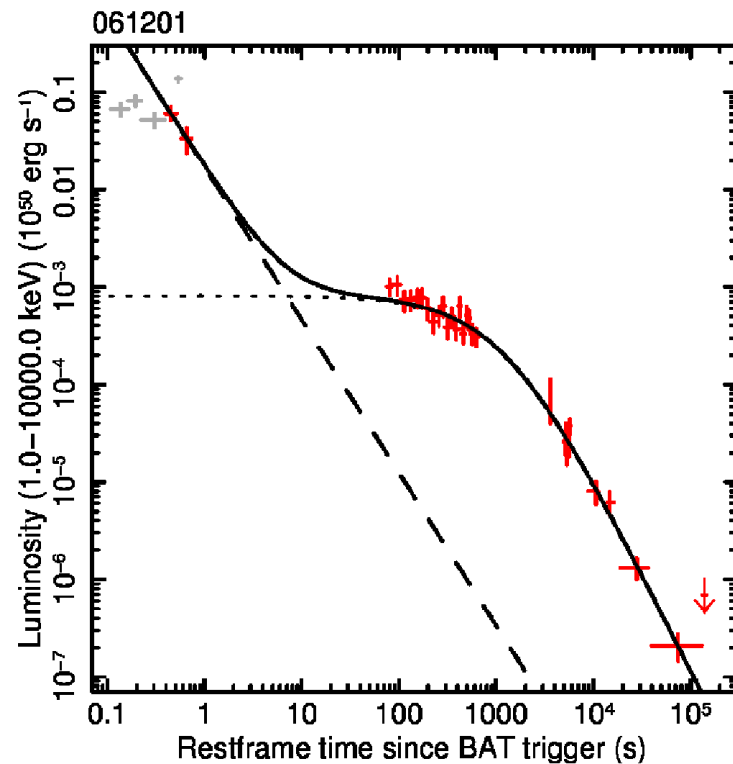
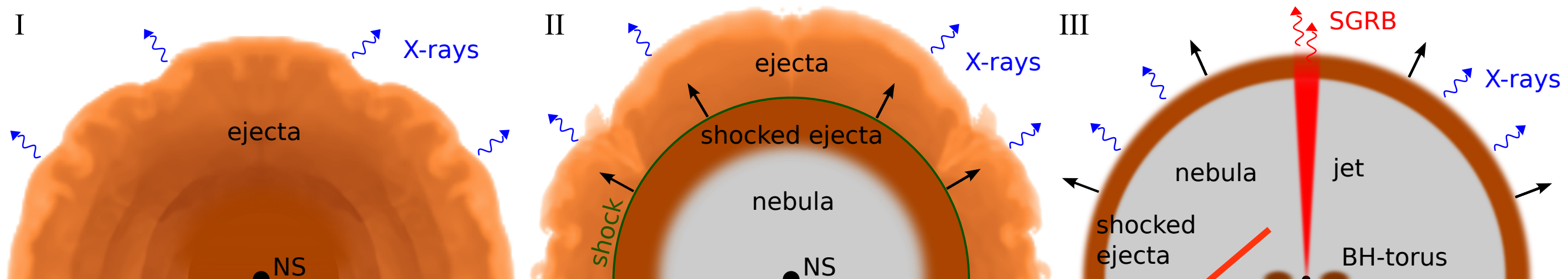
Ciolfi & Siegel 2015a, ApJ Letters 798, L36



- (I) The differentially rotating, supramassive NS (SMNS) ejects a baryon-loaded and highly isotropic wind
- (II) The cooled-down and uniformly rotating NS emits spin-down radiation inflating a photon-pair nebula that drives a shock through the ejecta
- ➡ (III) The NS collapses to a black hole (BH), a relativistic jet drills through the nebula and the ejecta shell and produces the prompt SGRB, while spin-down emission diffuses outwards on a much longer timescale, producing the X-ray afterglow

(but see Margalit et al. 2015)

Electromagnetic emission



The spin-down emission is **given off before** but (in part) **observed after** the prompt SGRB radiation

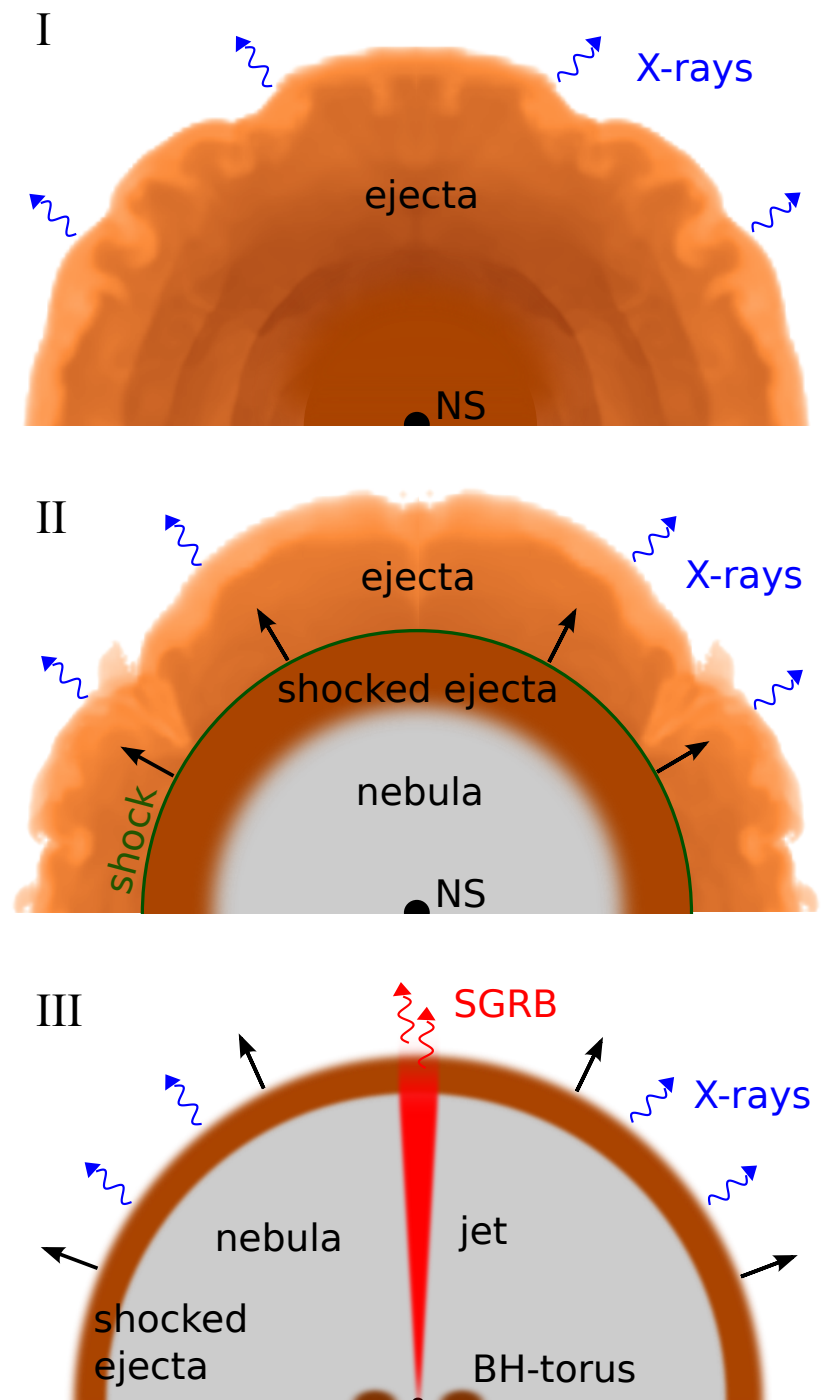
TR scenario: evidence

Ciolfi & Siegel 2015a, ApJ Letters 798, L36

- proposed new scenario to solve SGRB-X-ray afterglow dichotomy → “time-reversal” scenario
- delay times can explain observed X-ray afterglow durations
→ attractive alternative to current models

Evidence:

- potential observation of X-ray plateau with SGRB in between
→ indication of time reversal
- potential observation of an orphan event without SGRB
→ isotropy of afterglow



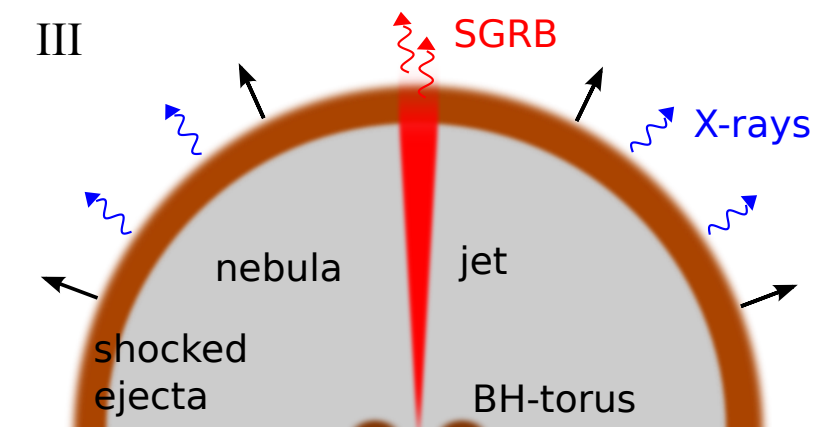
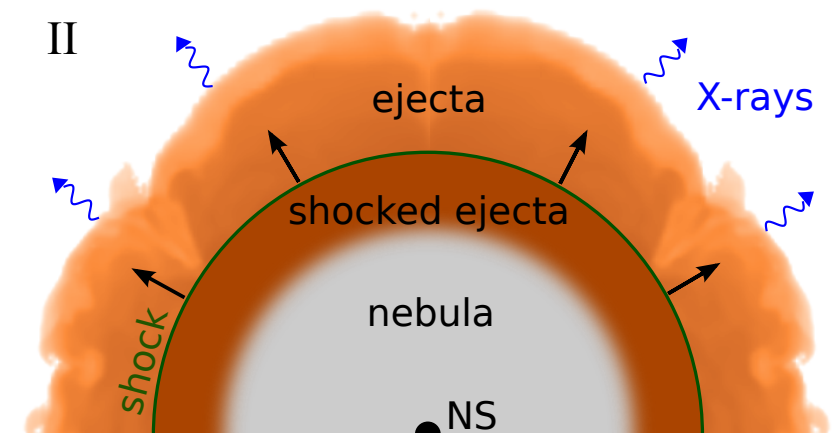
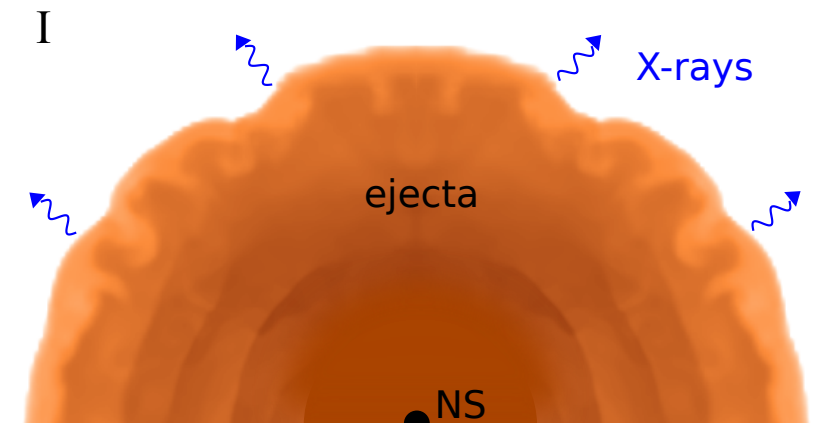
TR scenario: implications

Ciolfi & Siegel 2015a, ApJ Letters 798, L36

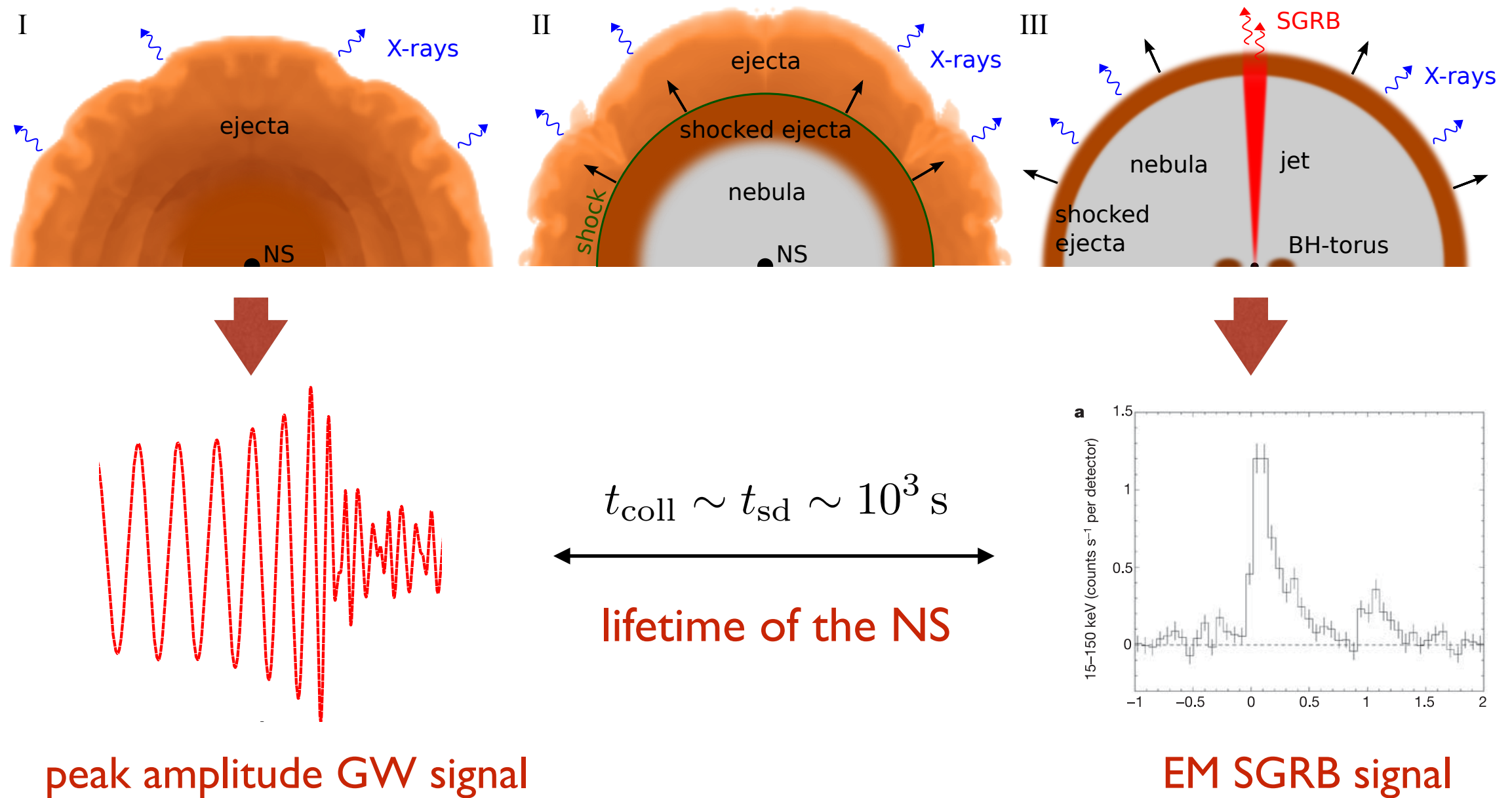
Implications:

- SGRBs with X-ray afterglows (majority of observed events) originate from **BNS mergers** → **no BH-NS progenitors**
- **SMNS constraints on EOS** in combination with a mass estimate
- peak amplitude of **GW emission** separated from **SGRB** by **lifetime of the NS**

Ciolfi & Siegel 2015b



GW and EM observations

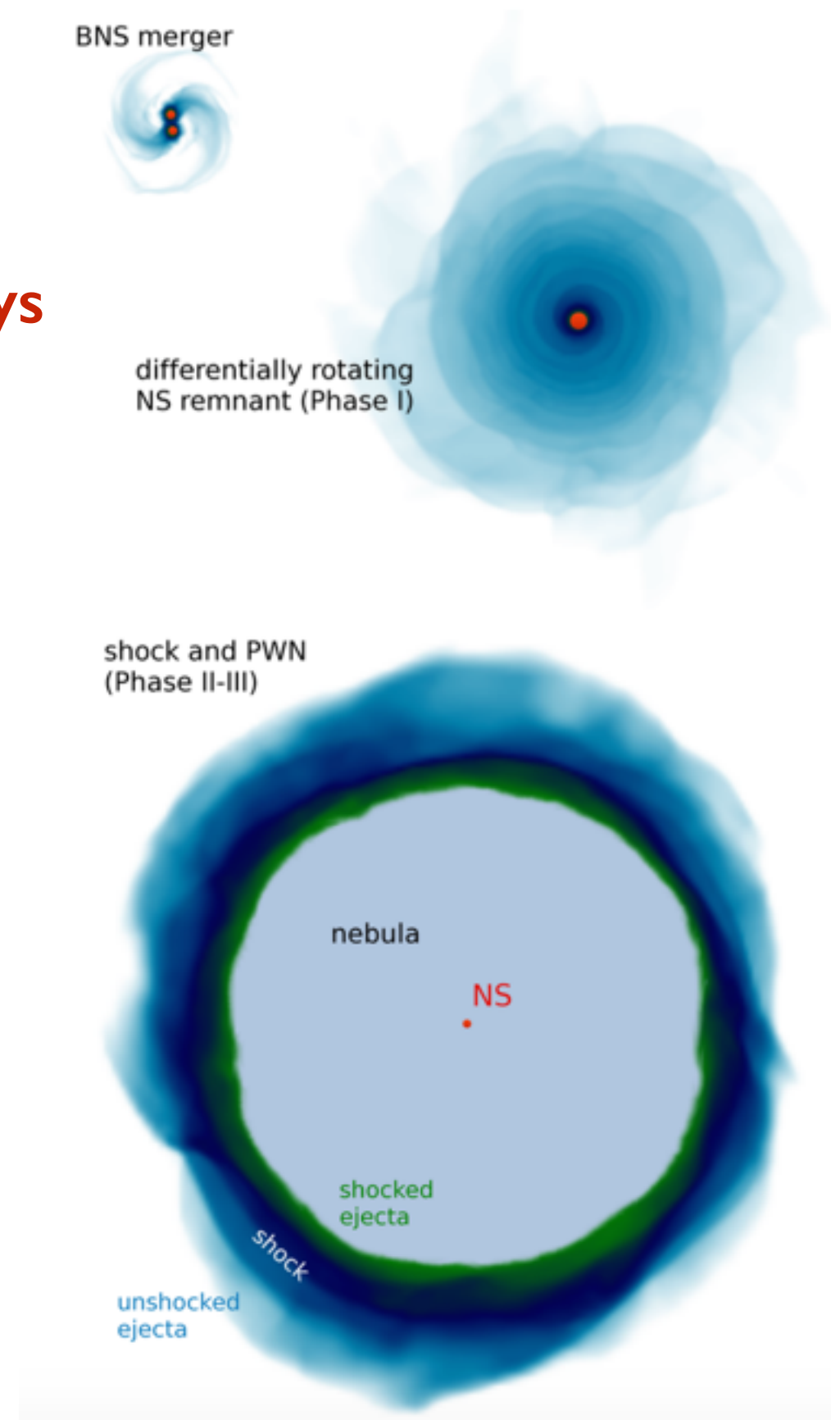


- GW observations ideal trigger for EM observations
- peak amplitude of **GW emission separated from SGRB** by **lifetime of the NS**
 → **very precise measurement of the NS lifetime!**

EM emission from the long-lived NS remnant

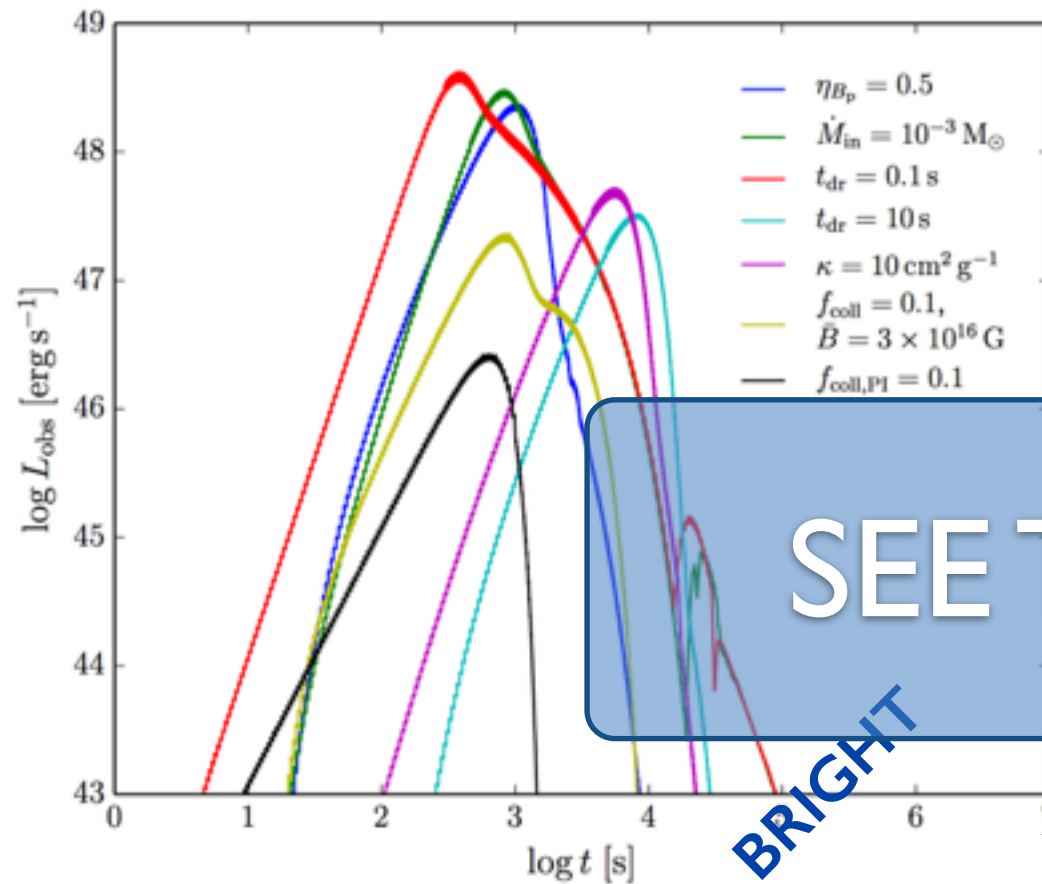
Siegel & Ciolfi 2016a, 2016b

- detail model **evolving the post-merger system** (NS remnant and surrounding environment) **on large scales and up to ~days**
→ far beyond the reach of numerical relativity simulations
- takes into account the relevant **radiative processes** (thomson, compton, synchrotron, pair production/annihilation, ..) and provides **lightcurves and spectra**
- first/only **self-consistent** model of its kind
- the model **can also account for the collapse** of the NS to a BH at any time during the evolution



EM emission from the long-lived NS remnant

Siegel & Ciolfi 2016a, 2016b



- signal peaks at 10^2 - 10^4 s (similar range for duration), with ~ 10 - 100 s delayed onset
- luminosities 10^{46} - 10^{49} erg/s

SEE TALK BY SIEGEL

BRIGHT

ISOTROPIC

LONG-LASTING

HIGH OCCURRENCE

DISTINGUISH BNS AND NS-BH

SGRB



KILONOVA



NS SPIN-DOWN

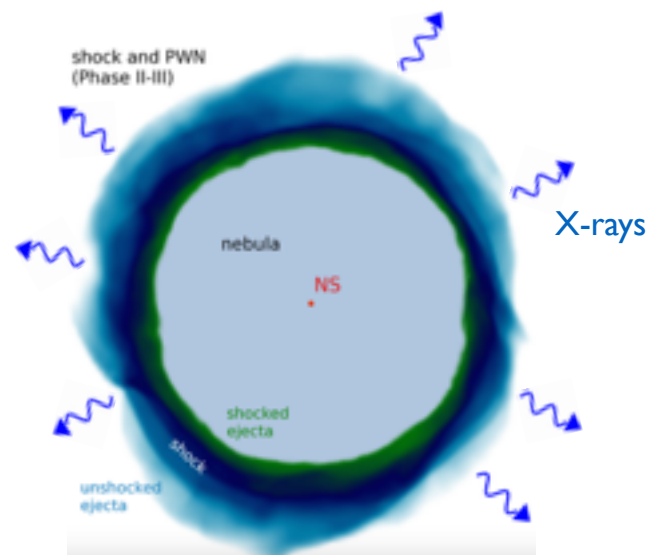
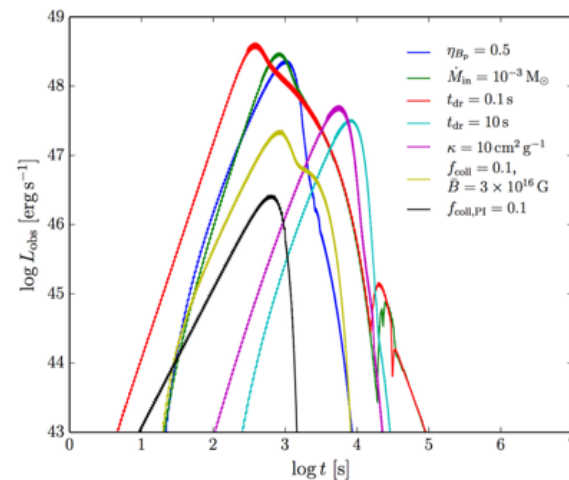


very promising EM counterpart!

X-ray flashes powered by NS spindown

Ciolfi 2016 (ArXiv:1606.01743)

- are XRFs really a subclass of long GRBs?



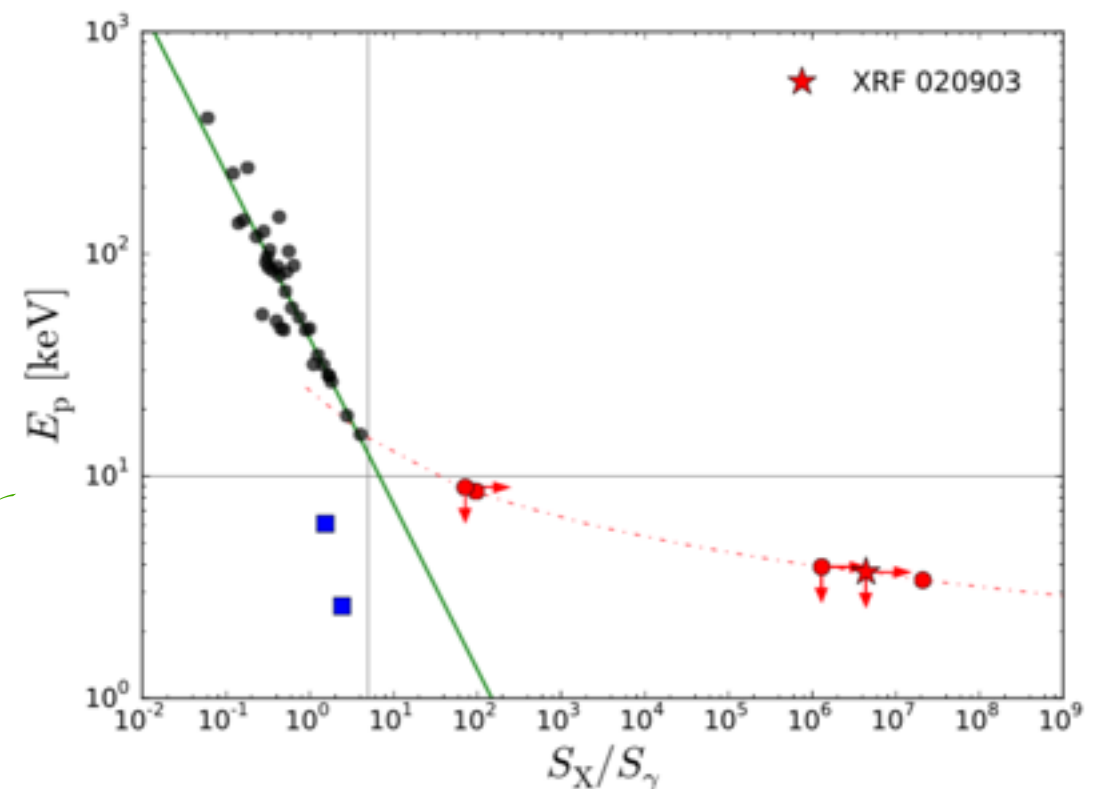
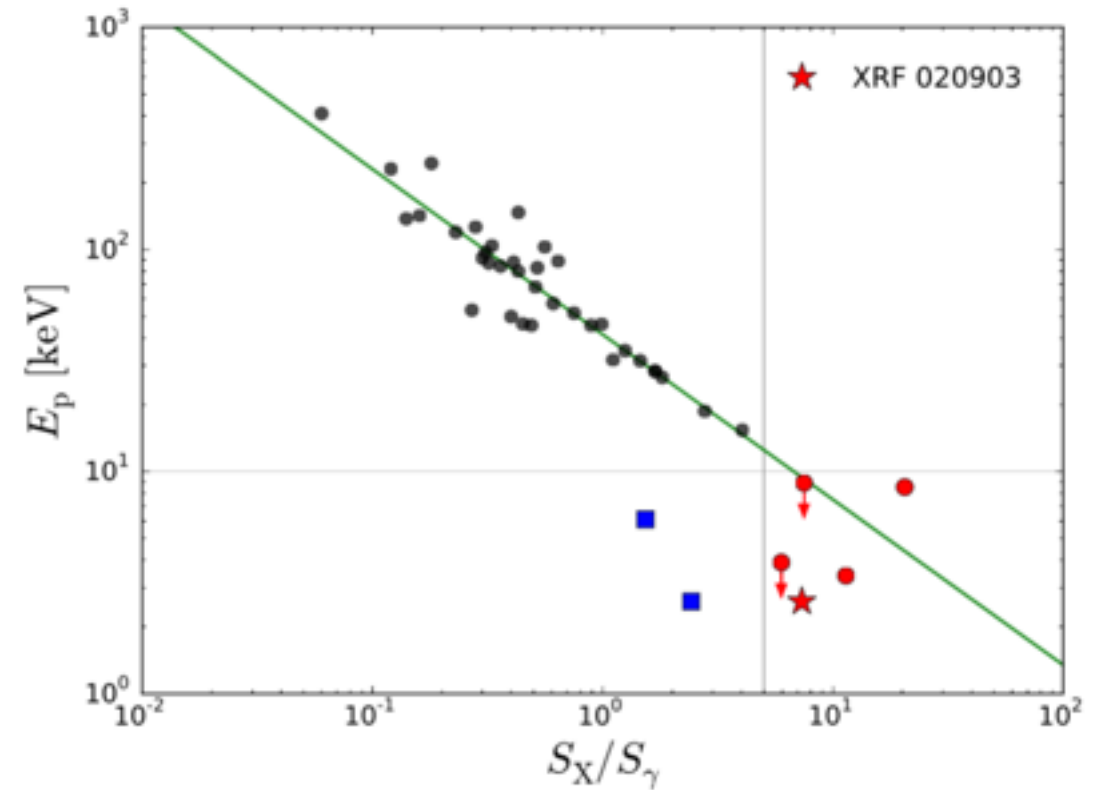
- spindown-powered X-ray emission from long-lived NSs **matches the high-energy emission of soft XRFs** (those emitting no gamma-rays)

spectral properties

- lack of gamma-rays ✓
- mainly thermal ✓
- black body T and its evolution ✓

lightcurve properties

- luminosity ✓
- duration ✓



WORKSHOP

SHORT GAMMA-RAY BURSTS

FROM OBSERVATION TO NUMERICAL SIMULATIONS

University of Trento, 8-9 September 2016
www.unitn.it/event/sgrb2016

INVITED SPEAKERS

Maria Grazia Bernardini (U. Montpellier)
Marica Branchesi (U. Urbino)
Pablo Cerda-Duran (U. Valencia)
Paolo D'avanzo (INAF Brera)
Bruno Giacomazzo (U. Trento)
Wolfgang Kastaun (U. Trento)
Kenta Kiuchi (U. Kyoto)
Paul O'Brien (U. Leicester)

Tsvi Piran (Hebrew U. Jerusalem)
Stephan Rosswog (U. Stockholm)
Ruben Salvaterra (INAF IASF)
Daniel Siegel (Columbia)
Giulia Stratta (U. Urbino)

ORGANIZERS

Riccardo Ciolfi (U. Trento)
Giancarlo Ghirlanda (INAF Brera)

Summary

- **long-lived NS** is a **very likely outcome of BNS mergers** and its EM emission can be extremely valuable to understand NS properties and SGRBs

see talk by Siegel on Thursday, C2 session

EM counterpart point of view

- X-ray signal powered by the spin-down of the long-lived NS remnant represents an **ideal counterpart to the GW signal of BNS mergers**

SGRB point of view

- long-lasting X-ray afterglows challenge the standard BH-torus picture and baryon pollution challenges the alternative magnetar model
- the time-reversal scenario offers a possible way out
- at the same time it gives a better match for the X-ray afterglow light curves

References

Ciolfi R. & Siegel D.M. (2015a) ApJ Letters 798, L36

Short gamma-ray bursts in the ‘time-reversal’ scenario

Ciolfi R. & Siegel D.M. (2015b) in “Swift: 10 Years of Discovery”, PoS(SWIFT 10)108

Short gamma-ray bursts from binary neutron star mergers: the time-reversal scenario

Siegel D.M., Ciolfi R., Rezzolla L. (2014) ApJ Letters 785, L6

Magnetically driven winds from differentially rotating neutron stars and X-ray afterglows of short gamma-ray bursts

Siegel D.M. & Ciolfi R. (2015) in “Swift: 10 Years of Discovery”, PoS(SWIFT 10)169

Magnetically-induced outflows from binary neutron star merger remnants

Siegel D.M. & Ciolfi R. (2016a) ApJ 189, 14

Electromagnetic emission from long-lived binary neutron star merger remnants I: formulation of the problem

Siegel D.M. & Ciolfi R. (2016b) ApJ 189, 15

Electromagnetic emission from long-lived binary neutron star merger remnants II: light curves and spectra

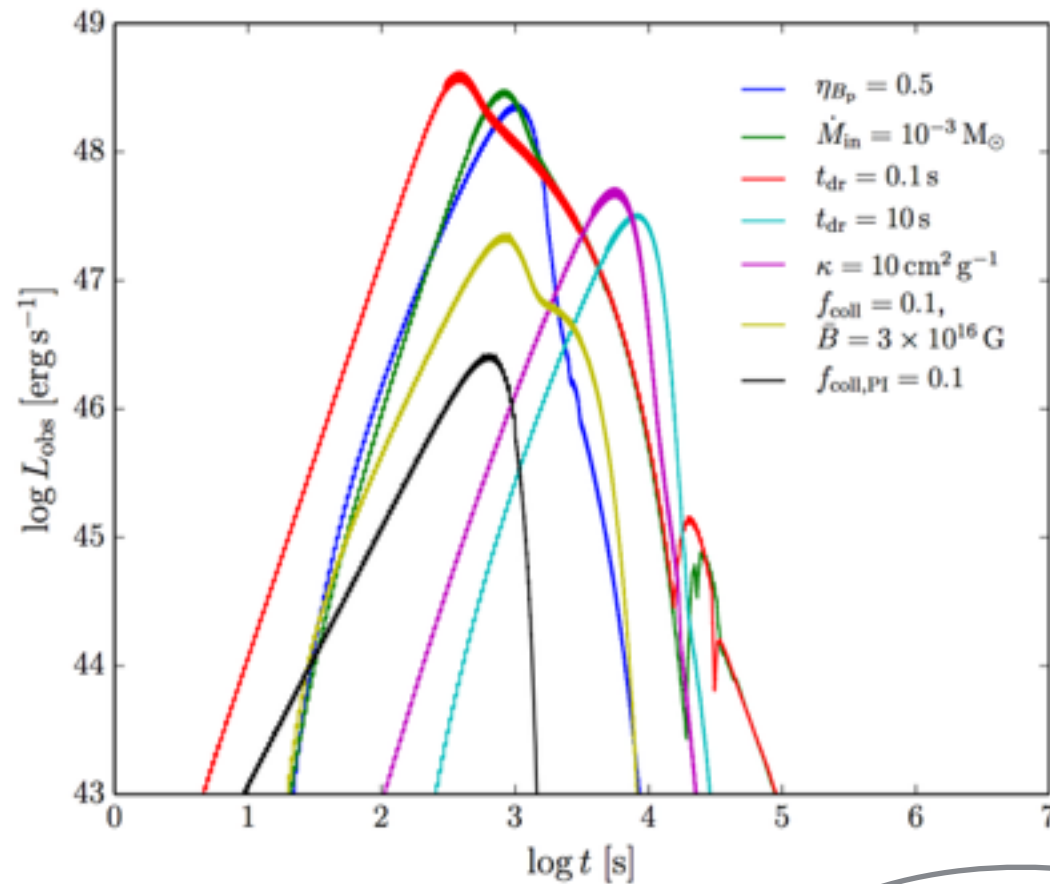
Ciolfi R. (2016) submitted to ApJ Letters, ArXiv:1606.01743

X-ray flashes powered by the spindown of long-lived neutron stars

BACKUP SLIDES

EM emission from the long-lived NS remnant

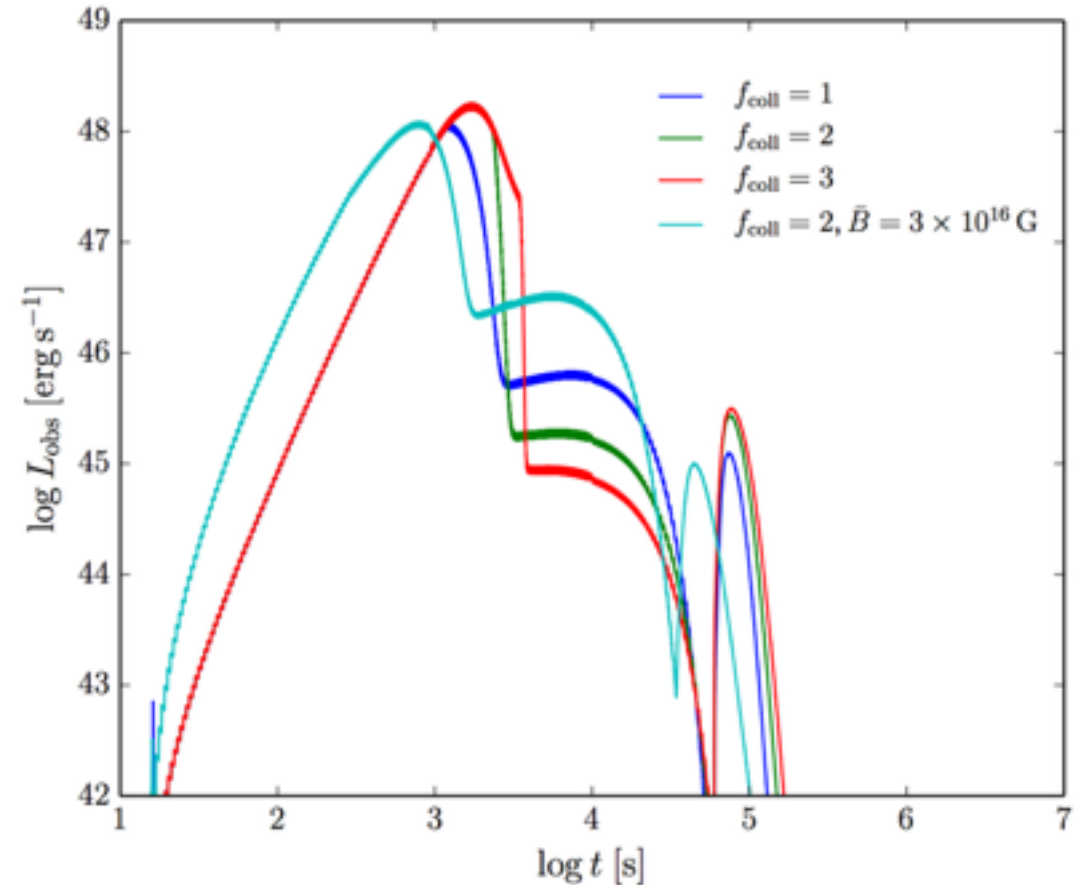
Siegel & Ciolfi 2016a, 2016b



non-collapsing
models

collapsing
models

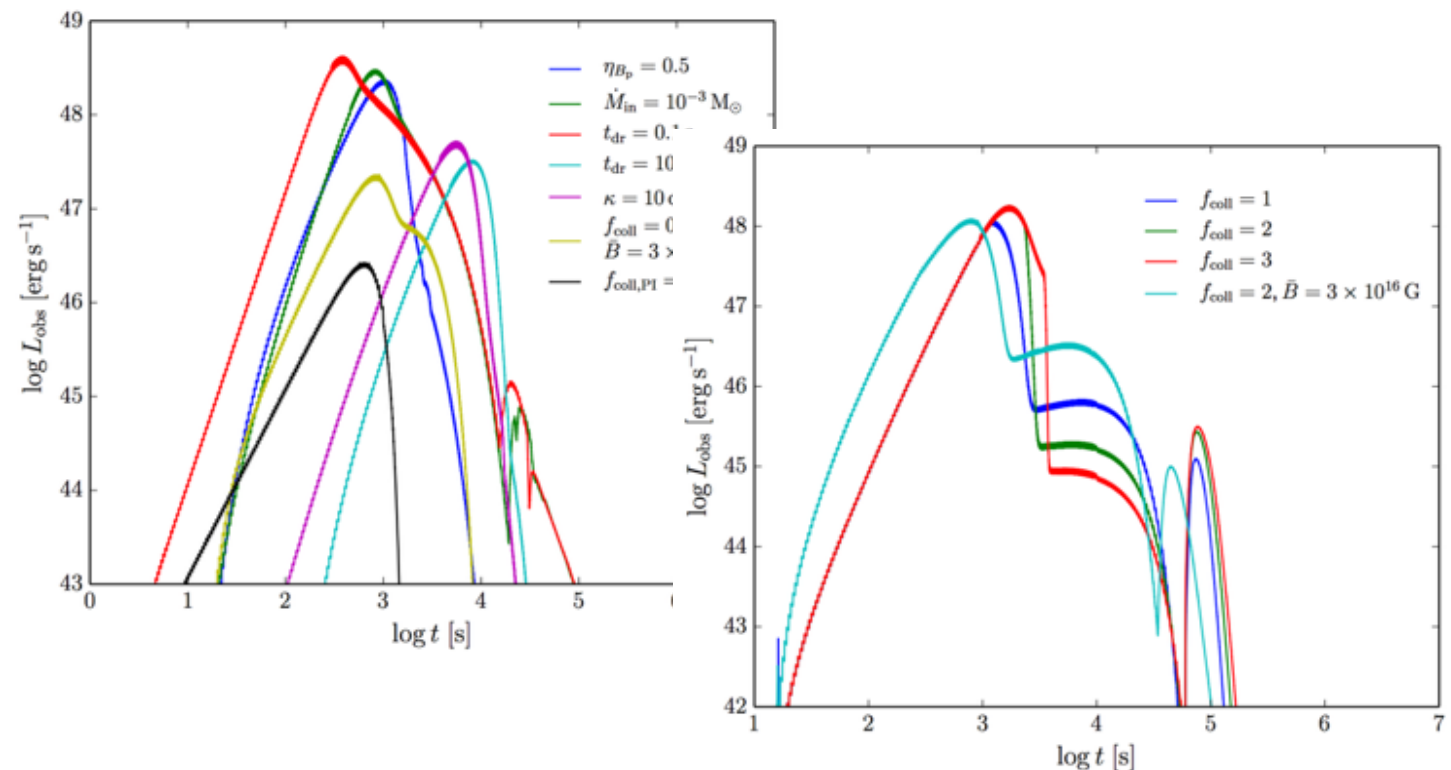
- signal peaks at 10^2 - 10^4 s (similar range for duration), with ~ 10 - 100 s delayed onset
- luminosities 10^{46} - 10^{49} erg/s
- mostly in the soft X-rays



Comparison with SGRB afterglows

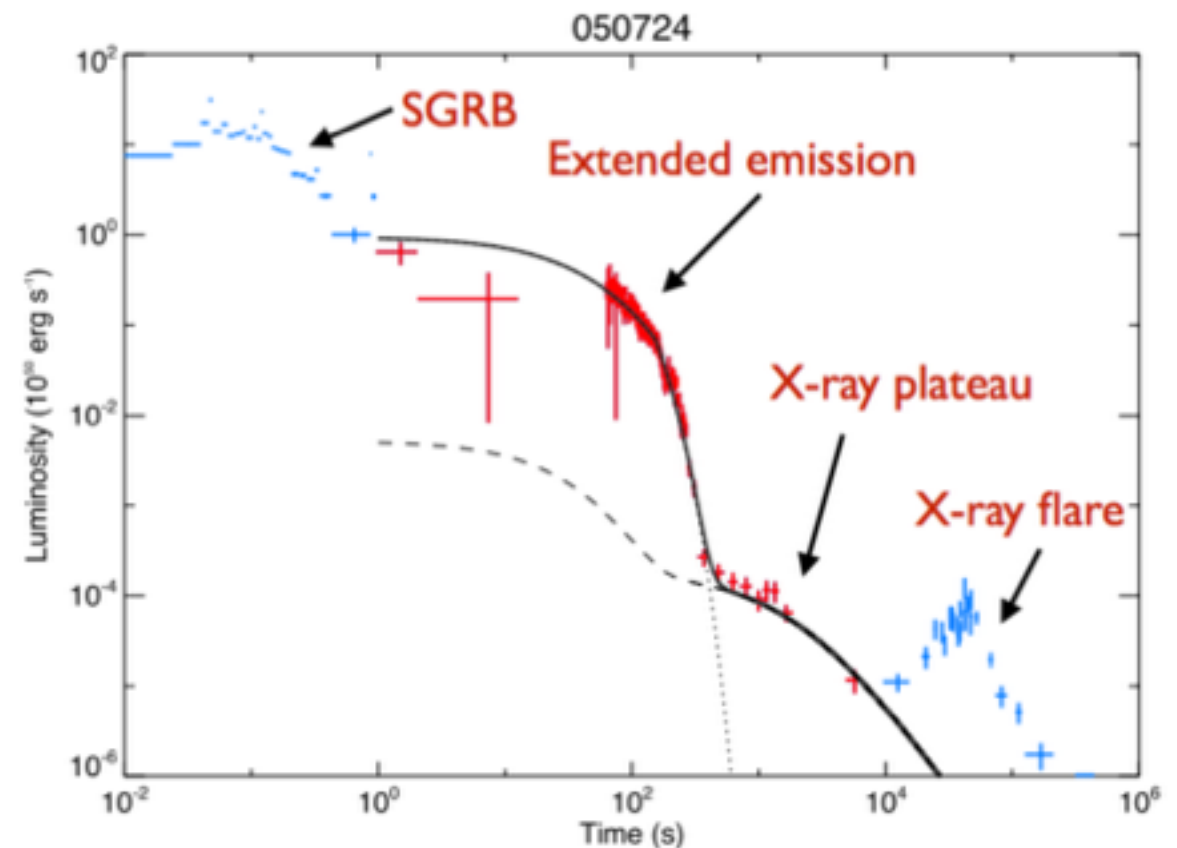
broad characteristics
in good agreement

- signals cover the right band (soft X-rays)
- very nice match with range of durations and luminosities

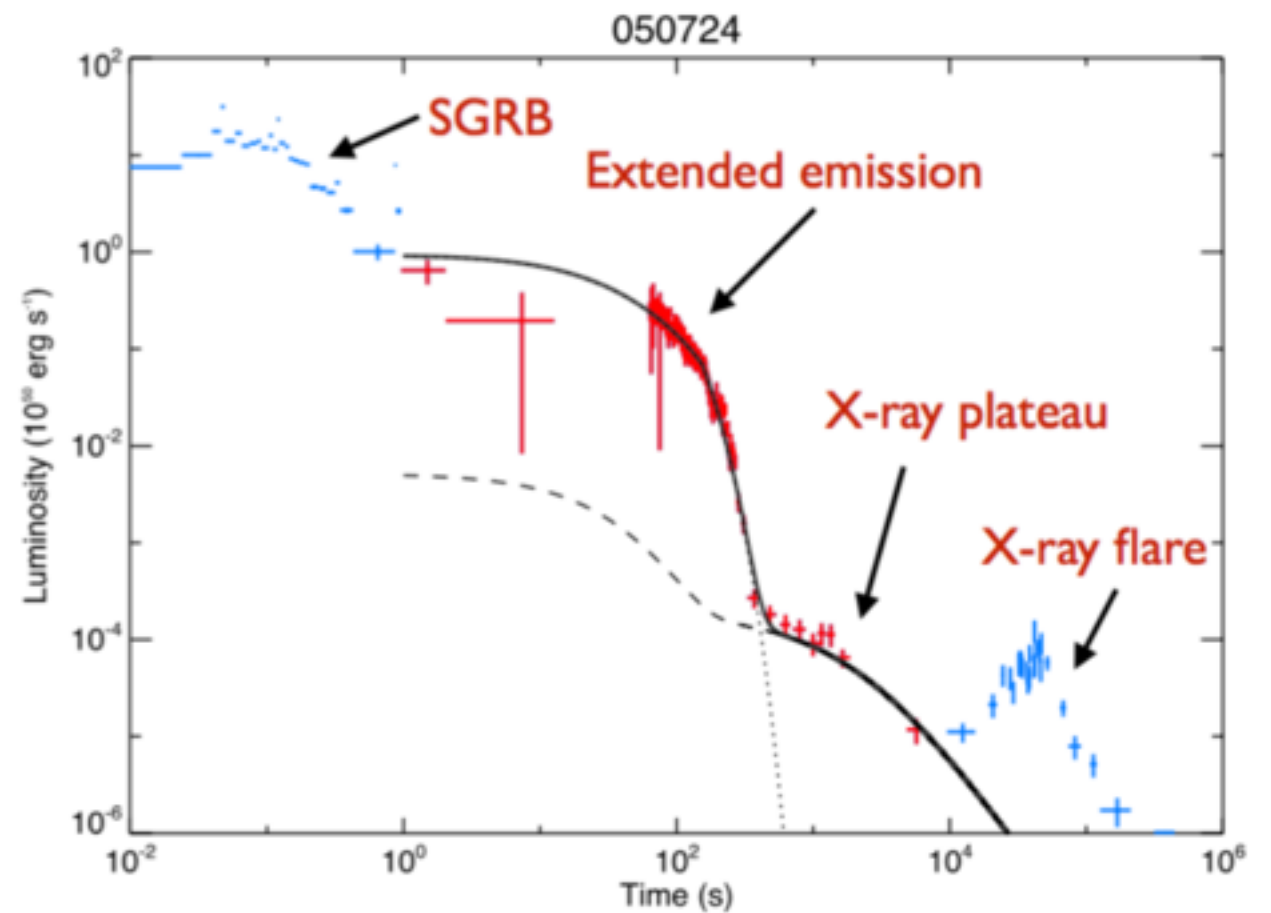
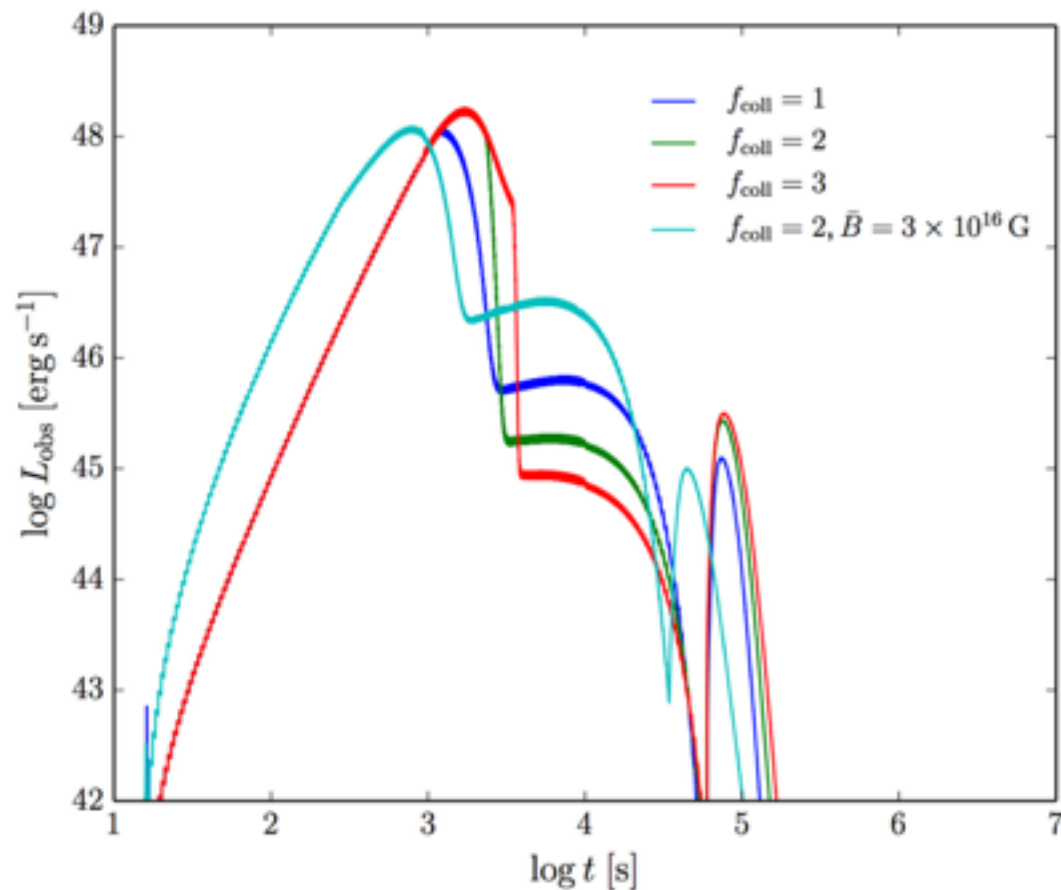


but with a closer look..

- no observations of early rising
- second plateau explained only for collapsing models
- flares explained as transition to optically thin ejecta?

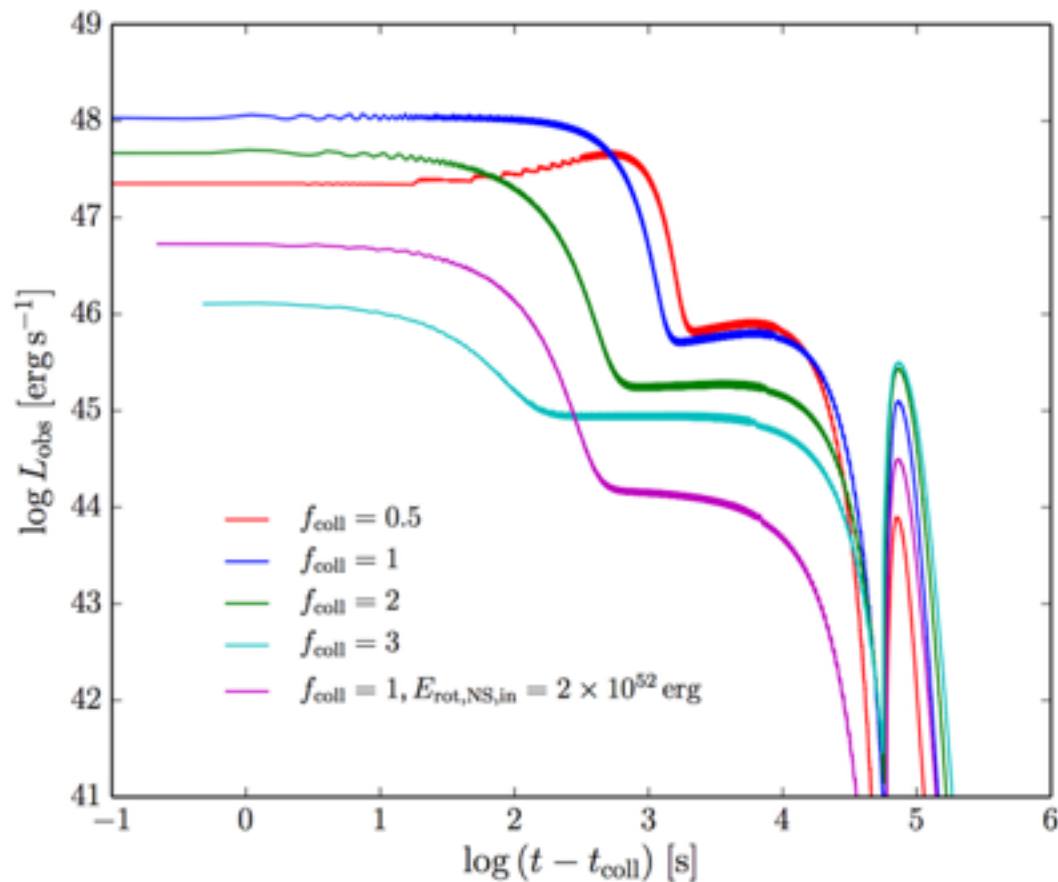


TR scenario - implications

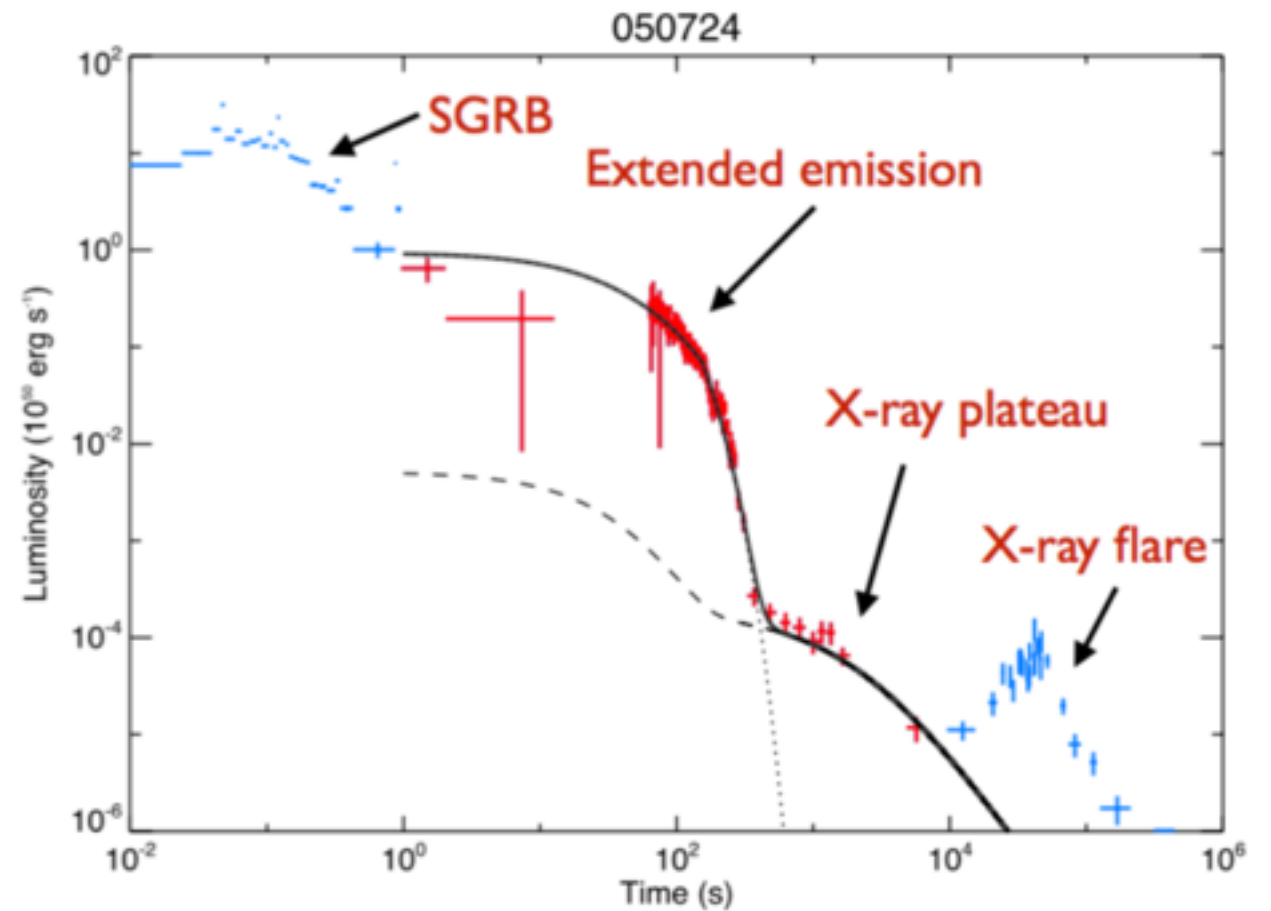


afterglows as seen by the
observer assuming **SGRB**
(trigger) at merger

TR scenario - implications

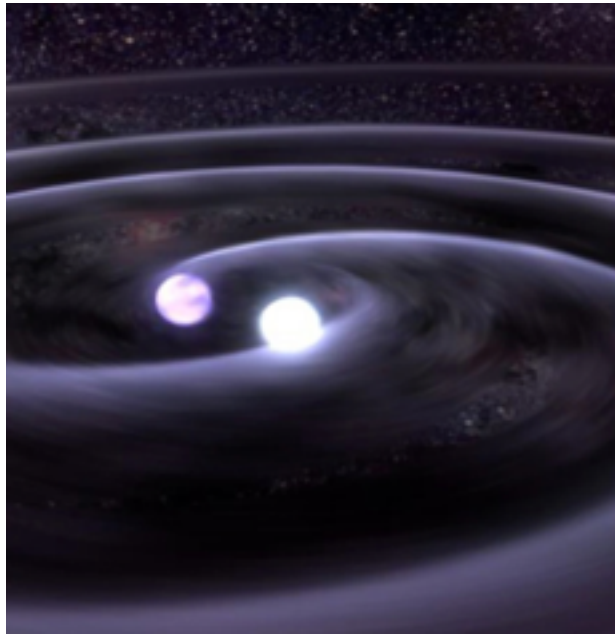


afterglows as seen by the
observer assuming **SGRB**
(trigger) at collapse



improved match assuming
the time-reversal scenario!

BNS mergers and EM counterparts



binary neutron star (BNS) mergers and
neutron star-black hole (NS-BH) binary mergers



most promising gravitational wave sources
for advanced LIGO and Virgo

	detection rate	best expectation
BNS	$\sim(0.4-400)/\text{yr}$	$\sim 40/\text{yr}$
NS-BH	$\sim(0.2-300)/\text{yr}$	$\sim 10/\text{yr}$

Abadie et al. 2010

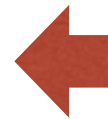
rewards of a combined
GW-EM detection:

- observed EM signals would incredibly **enhance the chances of GW detection**
- **EM follow-up observations** of a detected GW source is the ultimate way to unravel the nature of the system, by providing **crucial and complementary information**
- joint GW-EM signals can confirm the astrophysical origin of **short gamma-ray bursts**

SGRBs as EM counterparts

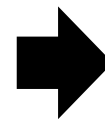
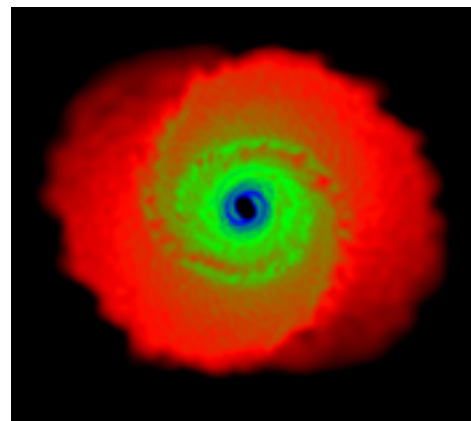
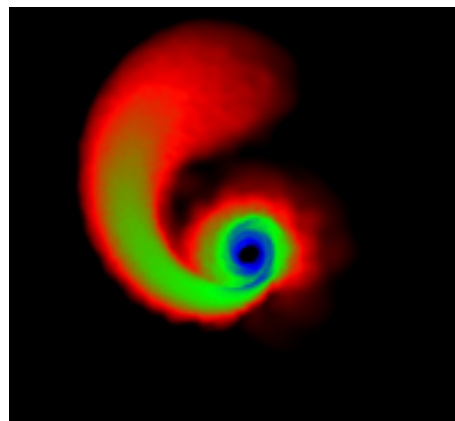
	BRIGHT	ISOTROPIC	LONG-LASTING	HIGH OCCURRENCE	DISTINGUISH BNS AND NS-BH
SGRB	✓	✗	✗	✗	✗

not very promising
as EM counterpart



Merger ejecta and r-process nucleosynthesis

ejecta in BNS and NS-BH mergers



r-process

capture rate much faster than decay
more than one neutron capture at a time
requires very special conditions:

- High T ($T > 10^9$ K)
- High neutron density ($n_n > 10^{22}$ cm $^{-3}$)



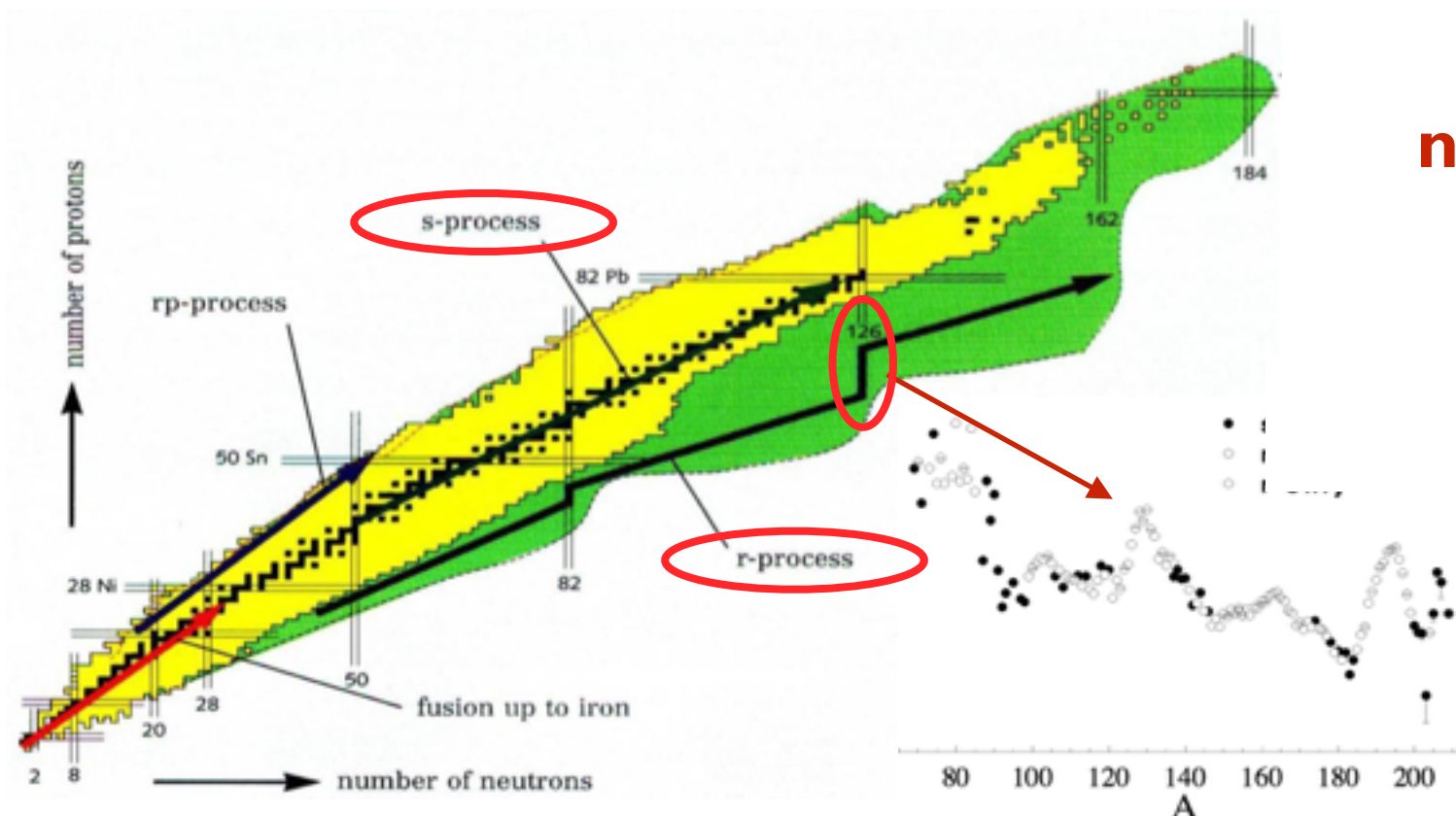
nucleosynthesis of heavy nuclei



initially unstable
radioactive decay on
timescales of >days



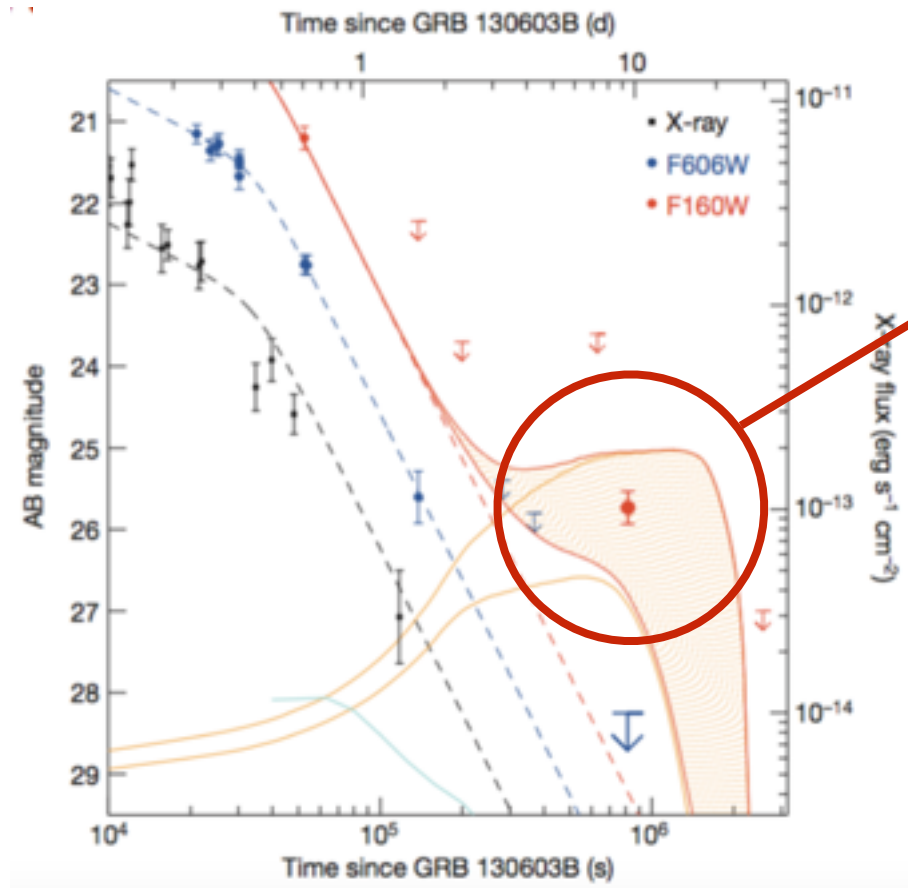
optical signal!



curtesy of A.Arcones

heavy element abundances

Kilonovae as EM counterparts



optical rebrightening in GRB 130603B
tentatively interpreted as a kilonova

connection

SGRB \longleftrightarrow BNS or NS-BH mergers

Tanvir et al. 2013, Berger et al. 2013

BRIGHT

ISOTROPIC

LONG-LASTING

**HIGH
OCCURRENCE**

**DISTINGUISH
BNS AND NS-BH**

SGRB

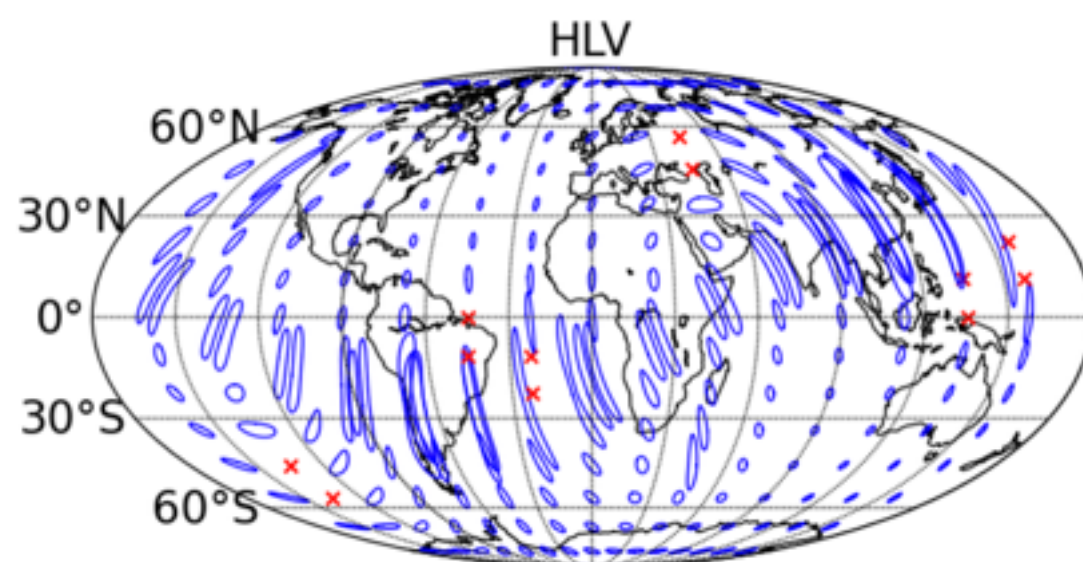


KILONOVA

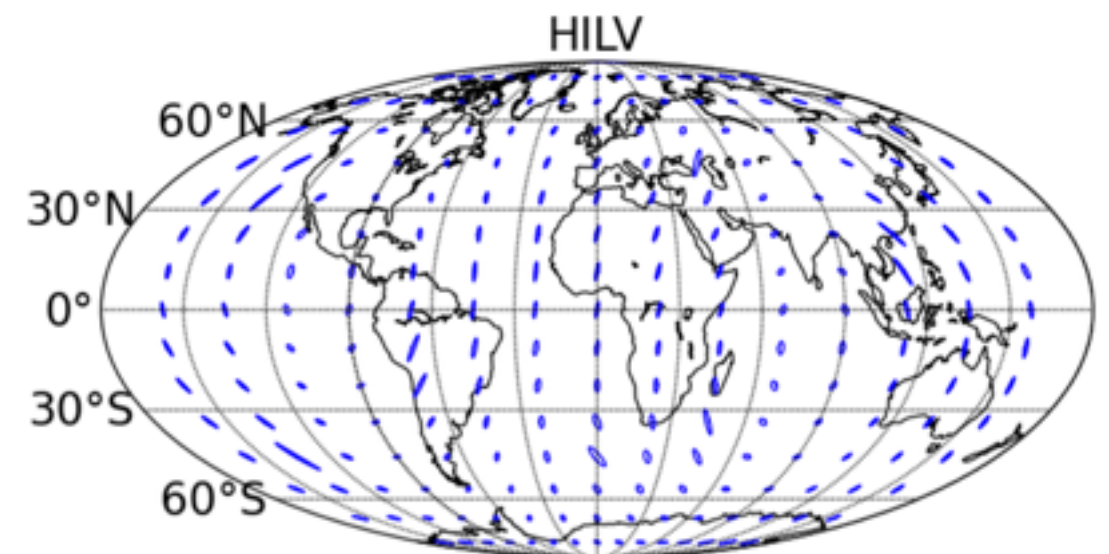


nice EM counterpart!

GW detector network



3 detectors



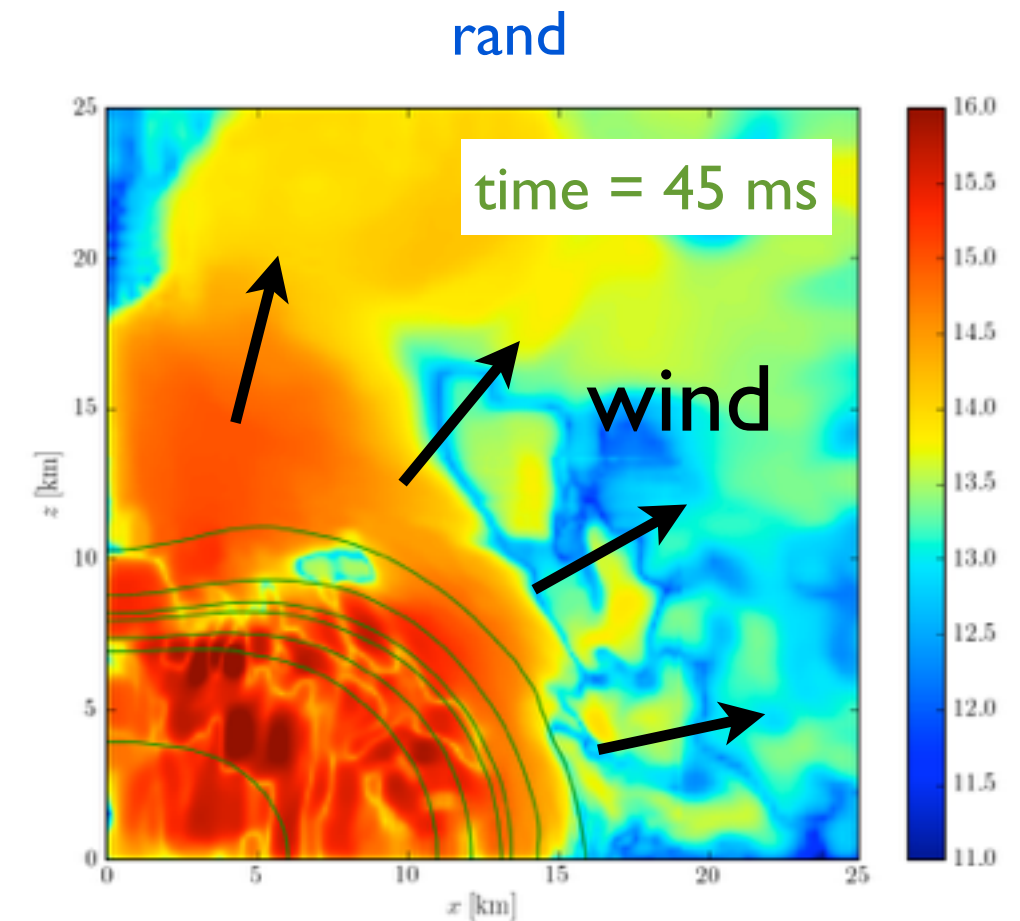
4 detectors

sky localization (90% confidence level)

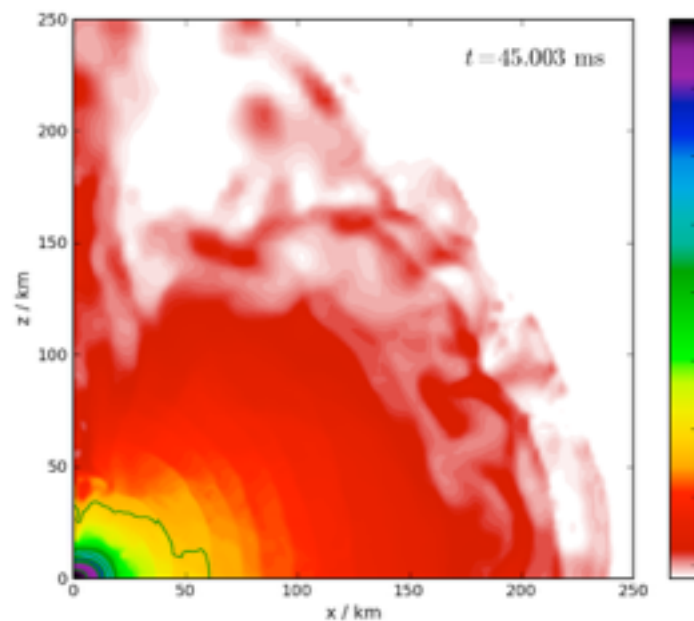
Baryon-loaded wind

- rest-mass density of the wind $\rho \sim 10^8 \text{ g/cm}^3$
- ejection speed $v \lesssim 0.1 \text{ c}$
- mass loss rate $\dot{M} \sim 10^{-3} \text{ M}_\odot/\text{s}$
- **mostly isotropic!**

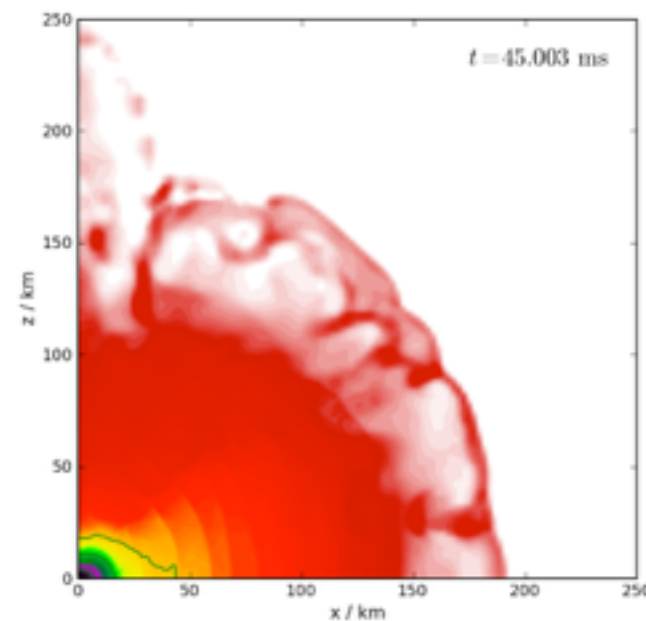
rest-mass density evolution ↓



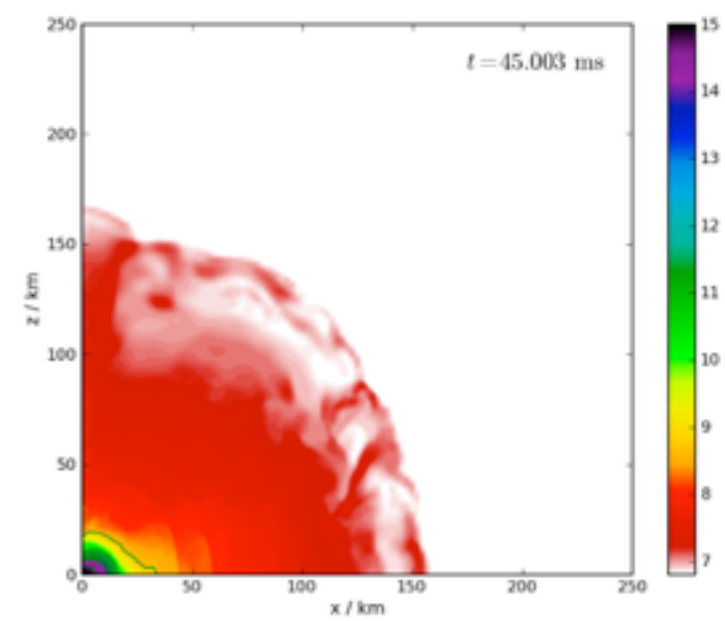
dipole 60



dipole 6

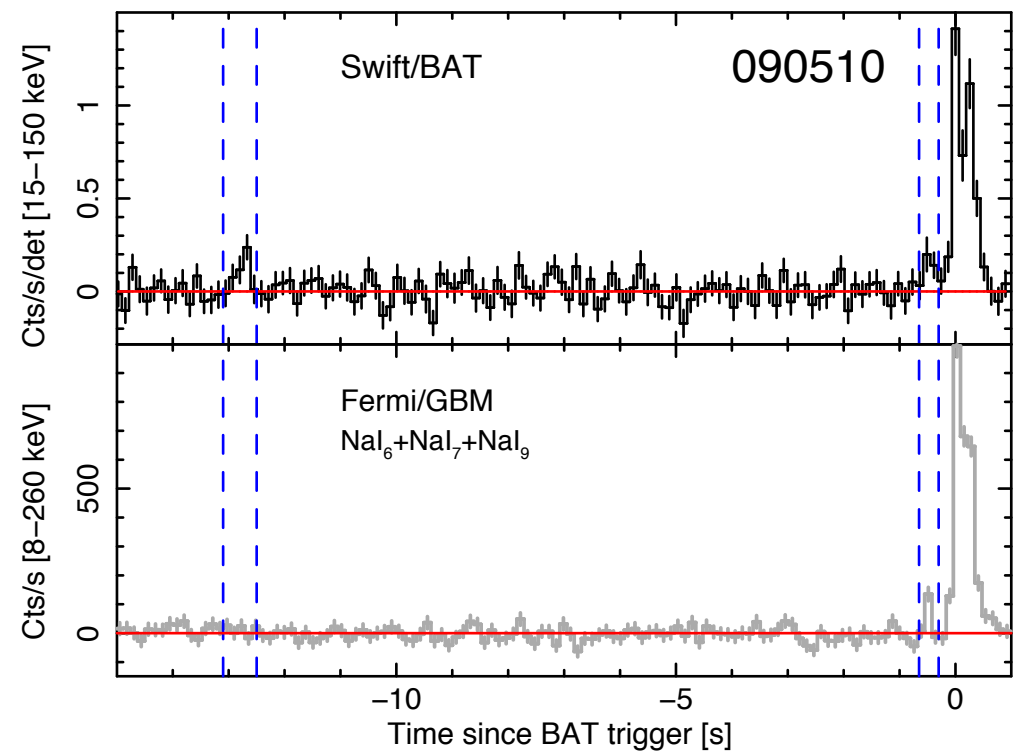
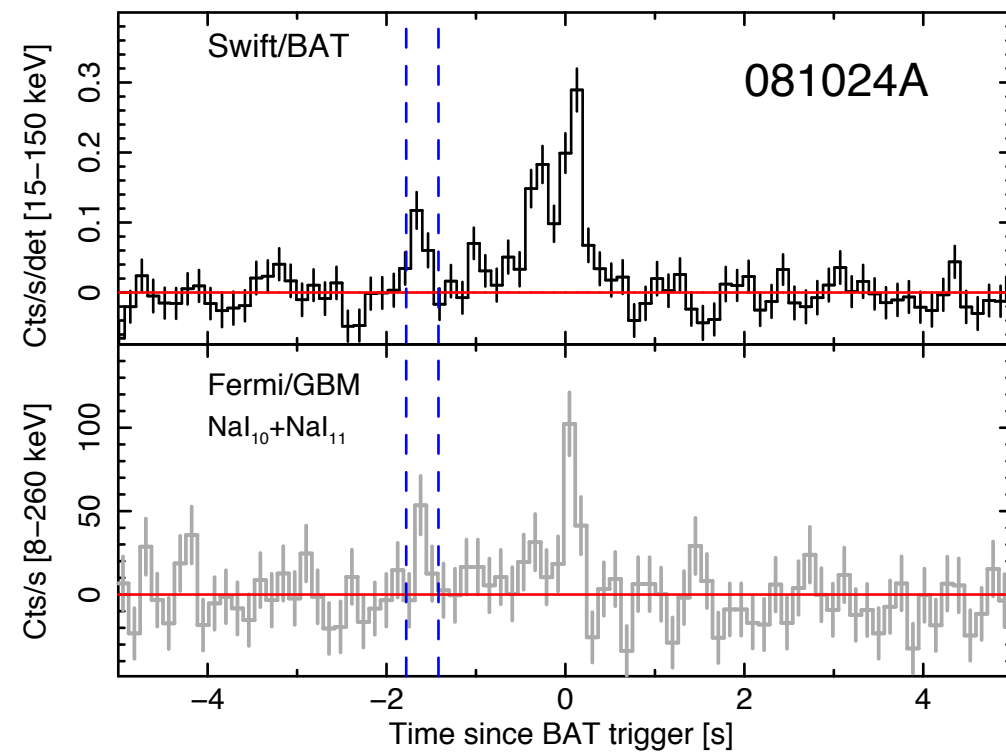
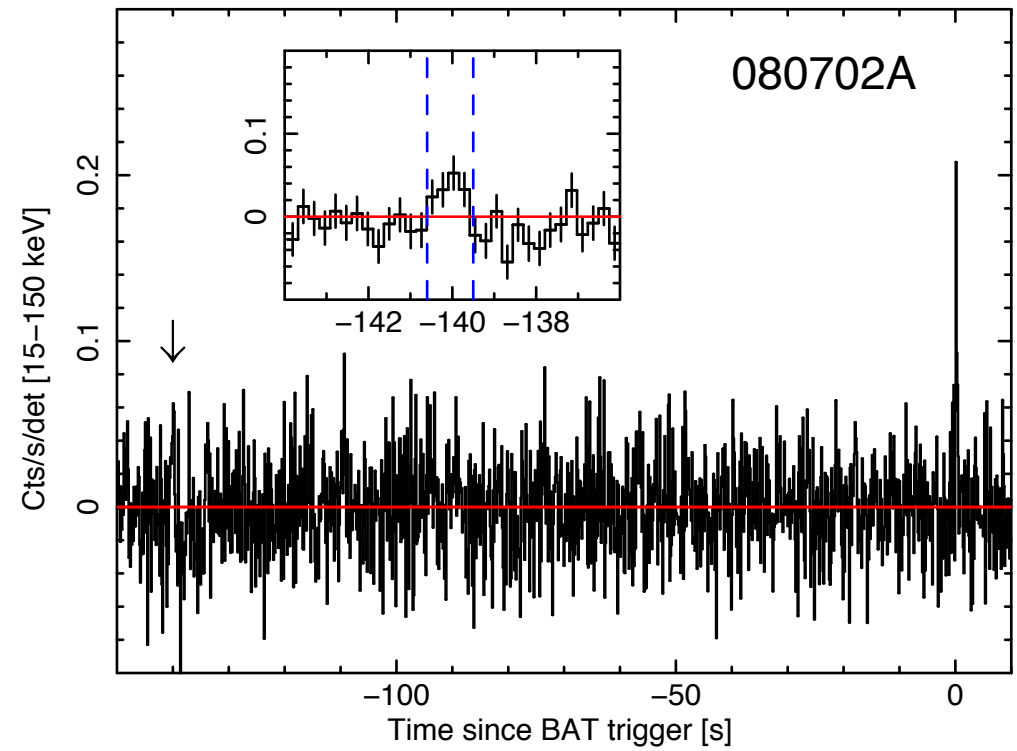
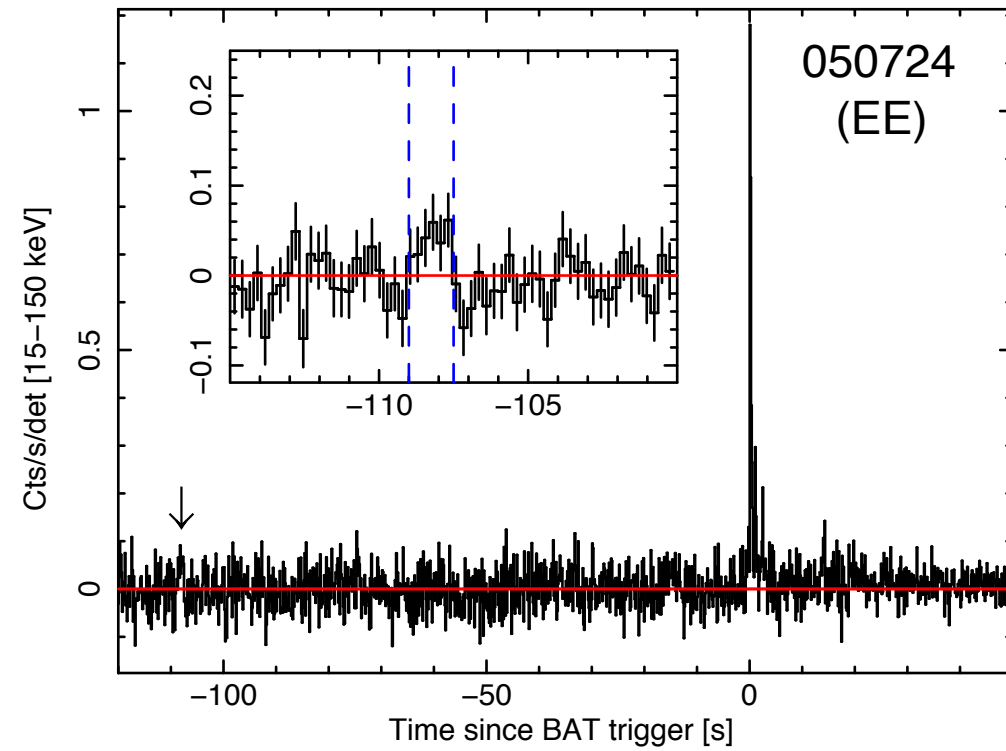


rand



SGRB precursors

Troja et al. 2010

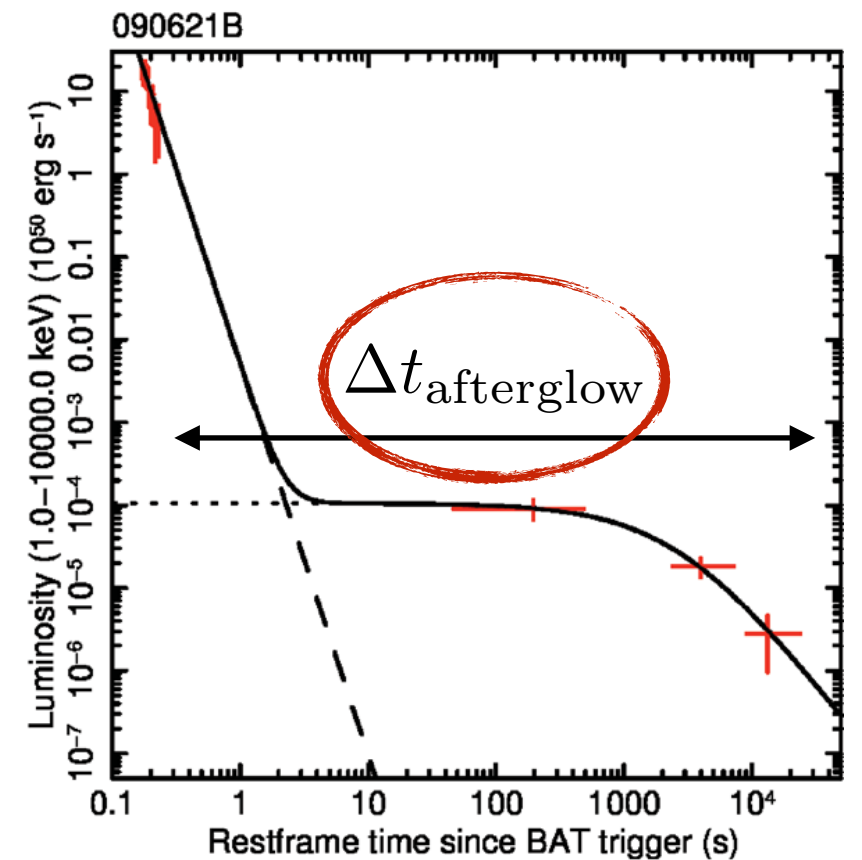
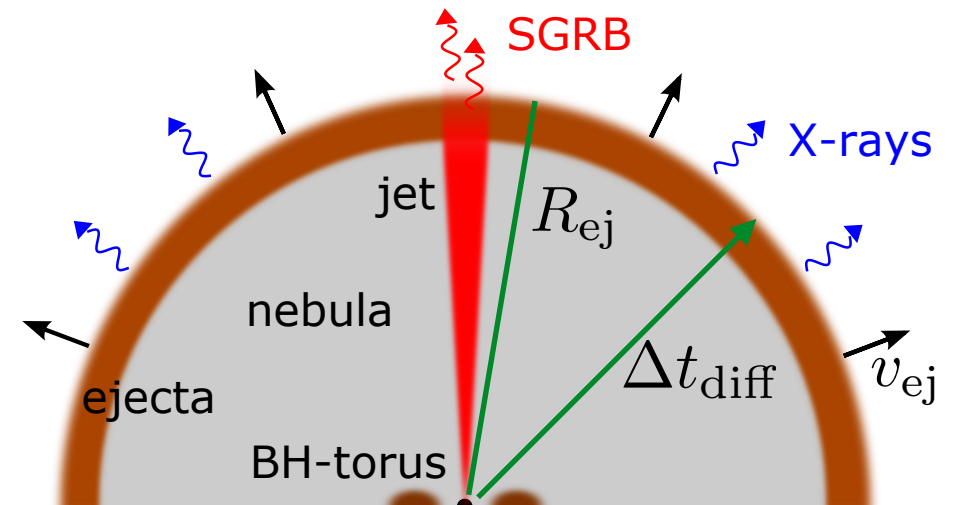


Timing argument

The scenario cannot hold unless the maximum delay is at least as large as the observed afterglow duration

- from observations: $t_{\text{coll}} \gtrsim t_{\text{sd}}$
- typically: $t_{\text{sd}} \gg t_{\text{dr}} + \Delta t_{\text{shock}}$
 \rightarrow at t_{coll} ejecta matter is swept up into thin shell
- delay for a photon emitted just before collapse (“last spin-down photon”):

$$t_{\text{NS}}^{\text{delay}} = \Delta t_{\text{diff}} - \frac{R_{\text{ej}}(t_{\text{coll}} + \Delta t_{\text{diff}})}{c}$$
- Delay of jet is negligible (very low densities at t_{coll})



Timing criterion

$$t_{\text{NS}}^{\text{delay}} \simeq t_{\text{NS}}^{\text{delay}} - t_{\text{jet}}^{\text{delay}} \gtrsim \Delta t_{\text{afterglow}}$$

Diffusion timescales

- for static ejecta:

$$t_{\text{diff}}^{\text{ej, stat}}(t) = \frac{\Delta_{\text{ej}}}{c} (1 + \kappa \rho_{\text{ej}}(t) \Delta_{\text{ej}}) \propto t^{-2}$$

- for static nebula:

$$t_{\text{diff}}^{\text{n, stat}}(t) = \frac{R_{\text{n}}(t)}{c} \left(1 + \sqrt{\frac{4Y \sigma_{\text{T}} L_{\text{sd}}(t)}{\pi R_{\text{n}}(t) m_{\text{e}} c^3}} \right) \propto t^{-1/2}$$

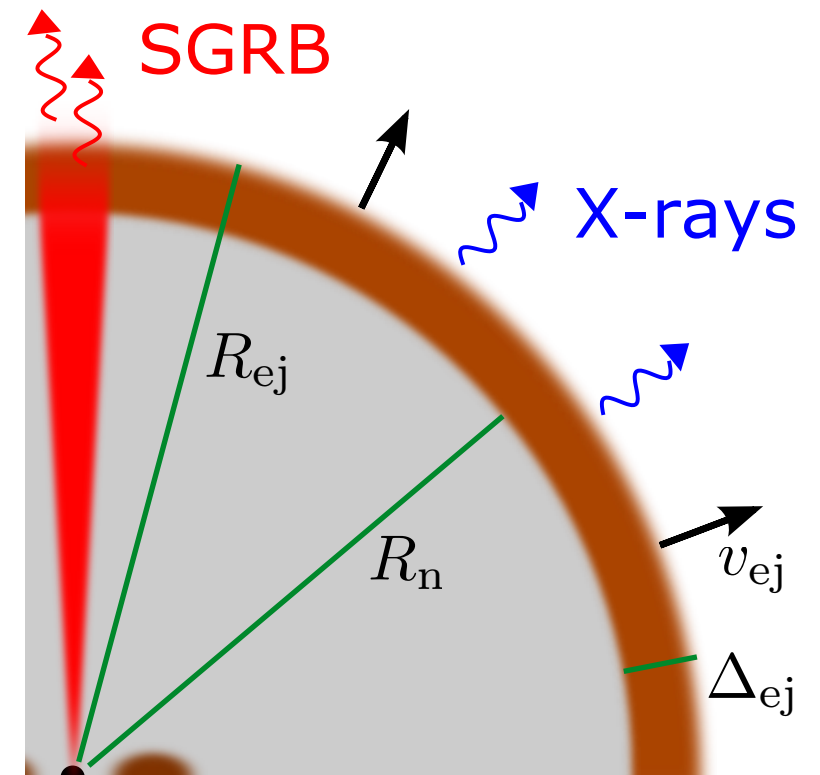
$$t_{\text{NS}}^{\text{delay}} \lesssim t_{\text{diff}}^{\text{ej, stat}}(t_{\text{coll}}) + t_{\text{diff}}^{\text{n, stat}}(t_{\text{coll}}) - R_{\text{ej}}(t_{\text{coll}})/c$$

$$t_{\text{NS}}^{\text{delay}} \gtrsim t_{\text{diff}}^{\text{ej, stat}}(t_{\text{coll}}^*) + t_{\text{diff}}^{\text{n, stat}}(t_{\text{coll}}^*) - R_{\text{ej}}(t_{\text{coll}}^*)/c$$

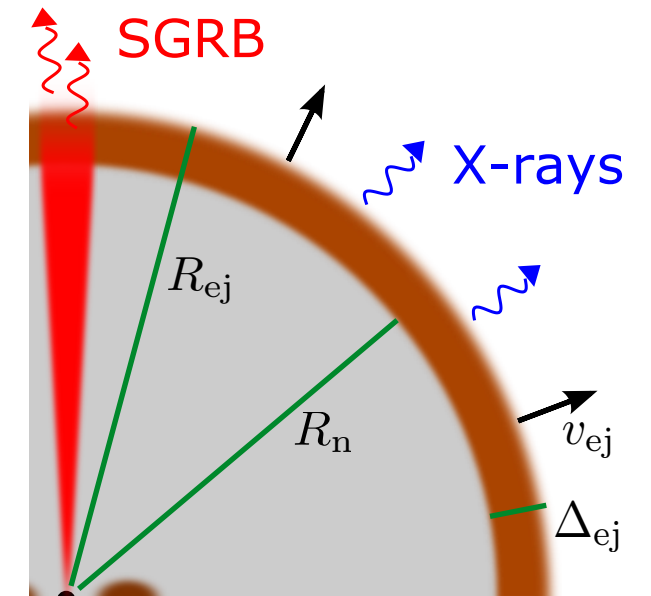
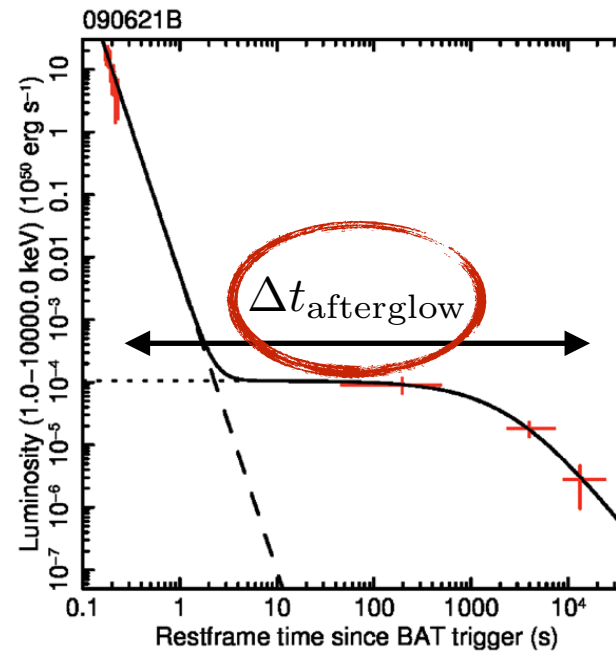
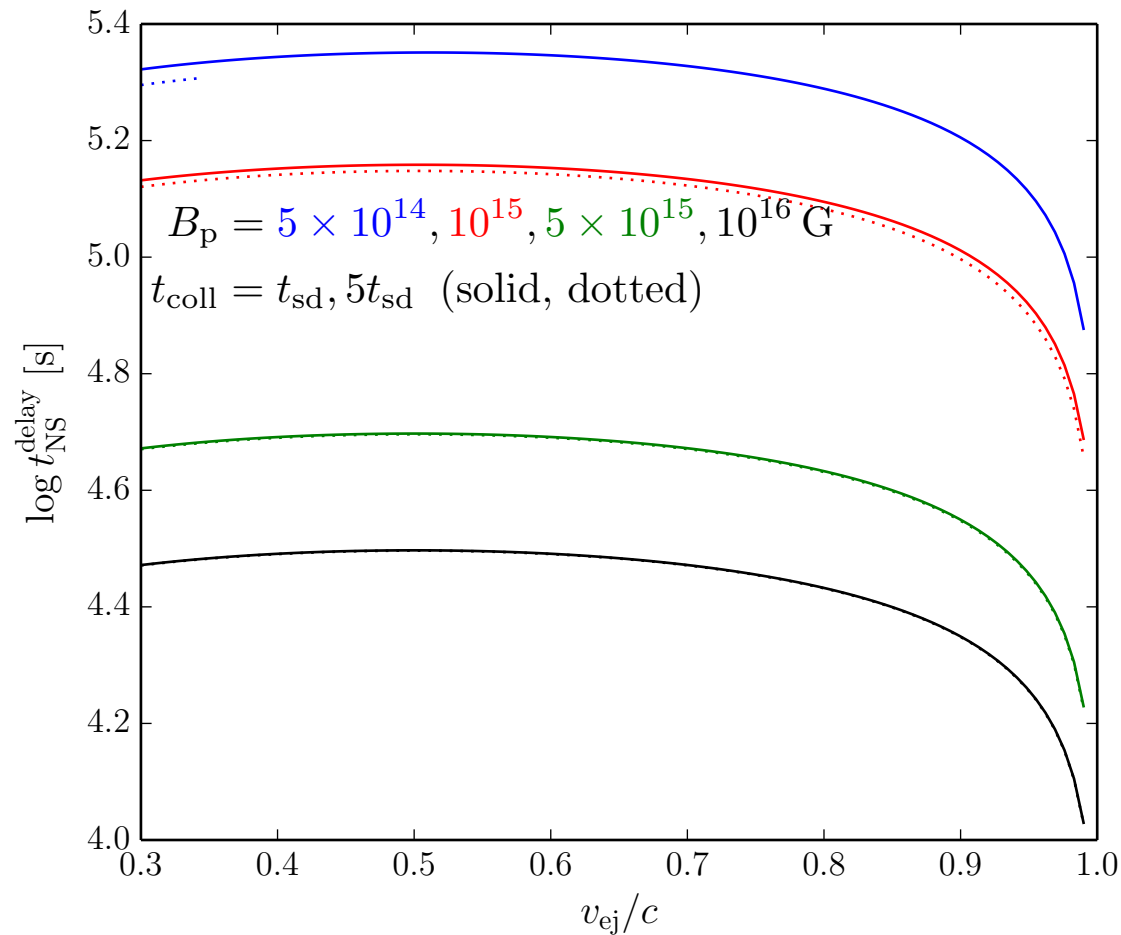
solve iteratively for t_{coll}^* :

$$t_{\text{coll}}^* = t_{\text{coll}} + t_{\text{diff}}^{\text{ej, stat}}(t_{\text{coll}}^*) + t_{\text{diff}}^{\text{n, stat}}(t_{\text{coll}}^*)$$

→ use lower limit to check the timing criterion



Results on delay estimation



parameter ranges:

$$B_p \sim 10^{14} - 10^{16} \text{ G}$$

$$P_{\text{in}} \sim 0.5 - 5 \text{ ms}$$

$$t_{\text{dr}} \sim 0.1 - 10 \text{ s}$$

$$\dot{M} \sim 10^{-4} - 10^{-2} M_{\odot} \text{ s}^{-1}$$

$$\Delta t_{\text{shock}} \sim 0 - 100 t_{\text{dr}}$$

$$v_{\text{ej}}^0 \sim 0.01 - 0.1 c$$

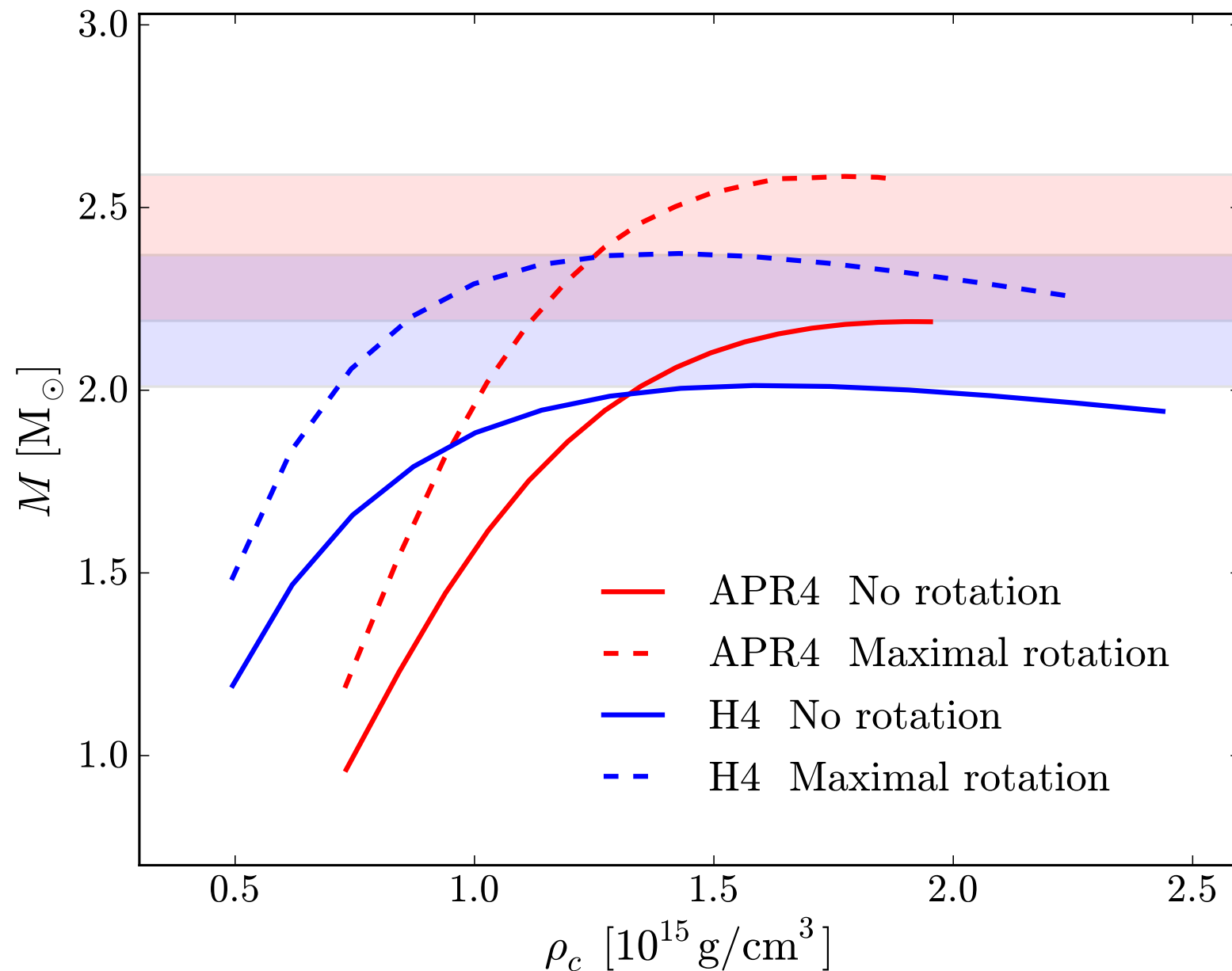
- for parameter ranges considered: $t_{\text{NS}}^{\text{delay}} > 3 \times 10^4 \text{ s}$
 $t_{\text{NS}}^{\text{delay}} \gtrsim 10^5 \text{ s} \quad (B_p \lesssim 10^{15} \text{ G})$

generally: $t_{\text{NS}}^{\text{delay}} \gtrsim \Delta t_{\text{afterglow}}$

→ “time-reversal” scenario compatible with observations

EOS constraint for a SMNS

Ciolfi & Siegel 2015b



most likely progenitor
mass combination

$$1.3 - 1.4 M_{\odot}$$

Belczynski et al. 2008



$$M \sim 0.9(M_1 + M_2 - 0.1)$$



$$M \sim 2.34 M_{\odot}$$

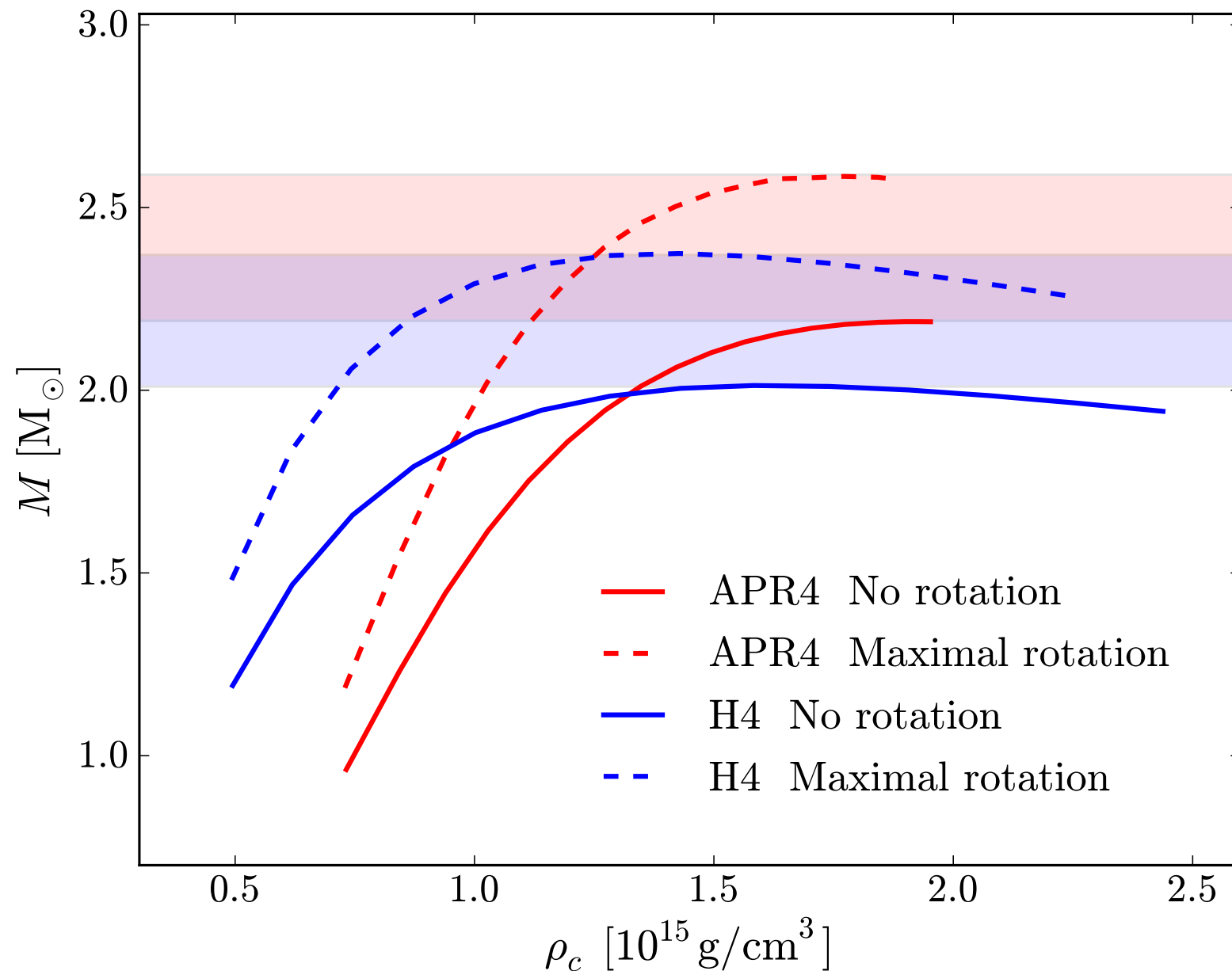
$$(M_b \sim 2.87 M_{\odot})$$

APR4 ✓

H4 ✓

EOS constraint for a SMNS

Ciolfi & Siegel 2015b



progenitor
mass combination

$$1.4 - 1.4 M_{\odot}$$



$$M \sim 0.9(M_1 + M_2 - 0.1)$$



$$M \sim 2.43 M_{\odot}$$

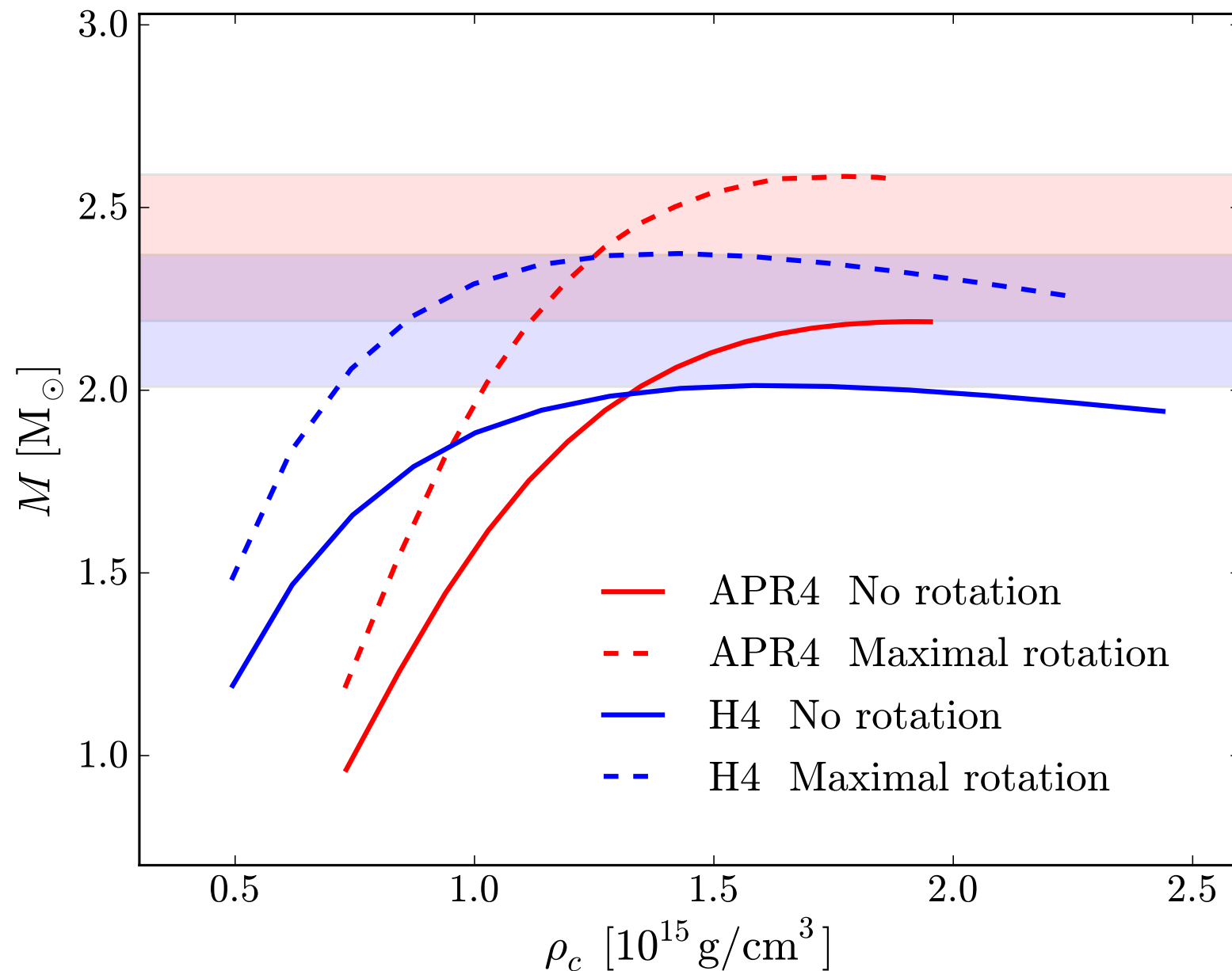
$$(M_b \sim 2.98 M_{\odot})$$

APR4 ✓

H4 ✗

EOS constraint for a SMNS

Ciolfi & Siegel 2015b



progenitor
mass combination

$$1.5 - 1.5 M_\odot$$



$$M \sim 0.9(M_1 + M_2 - 0.1)$$



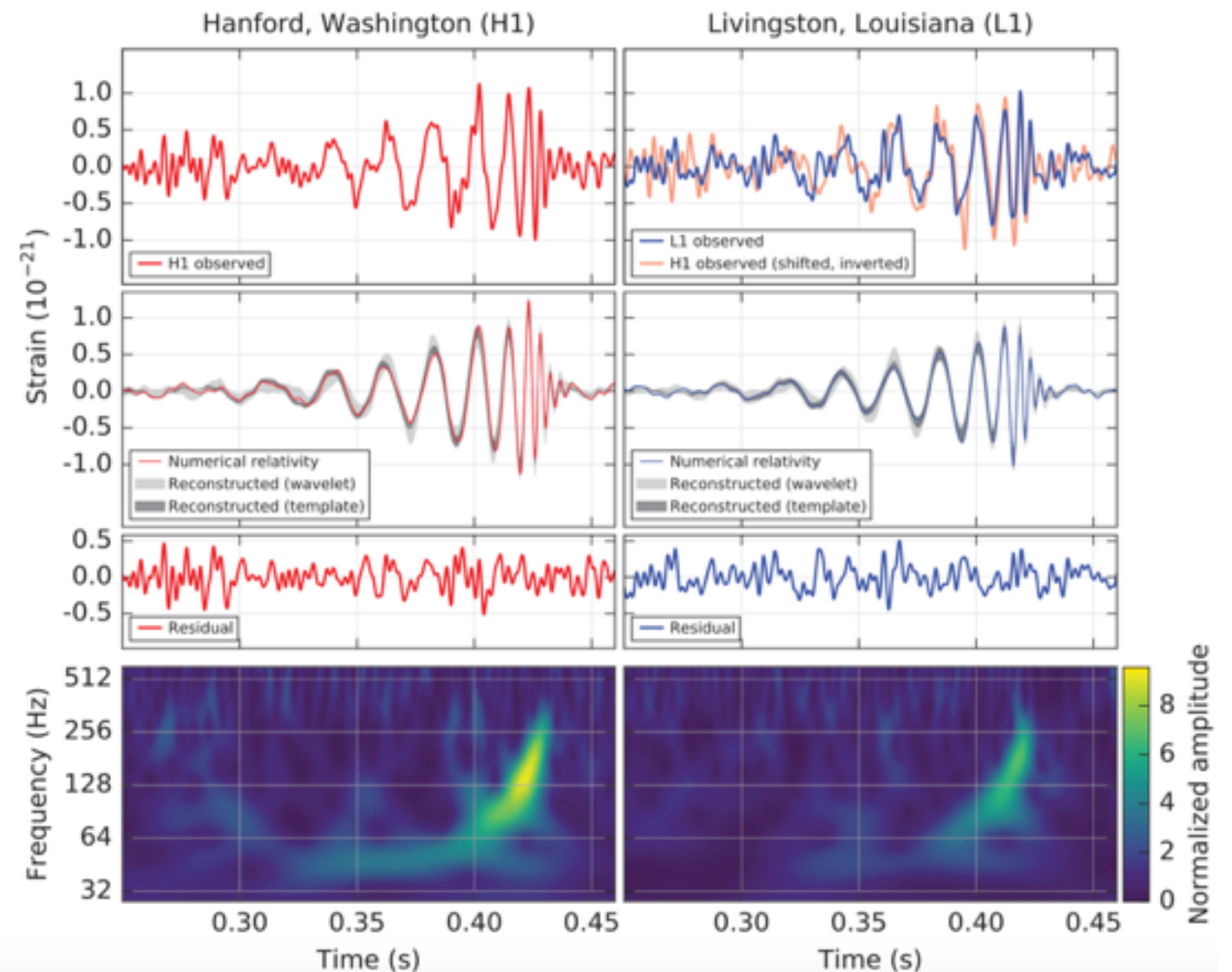
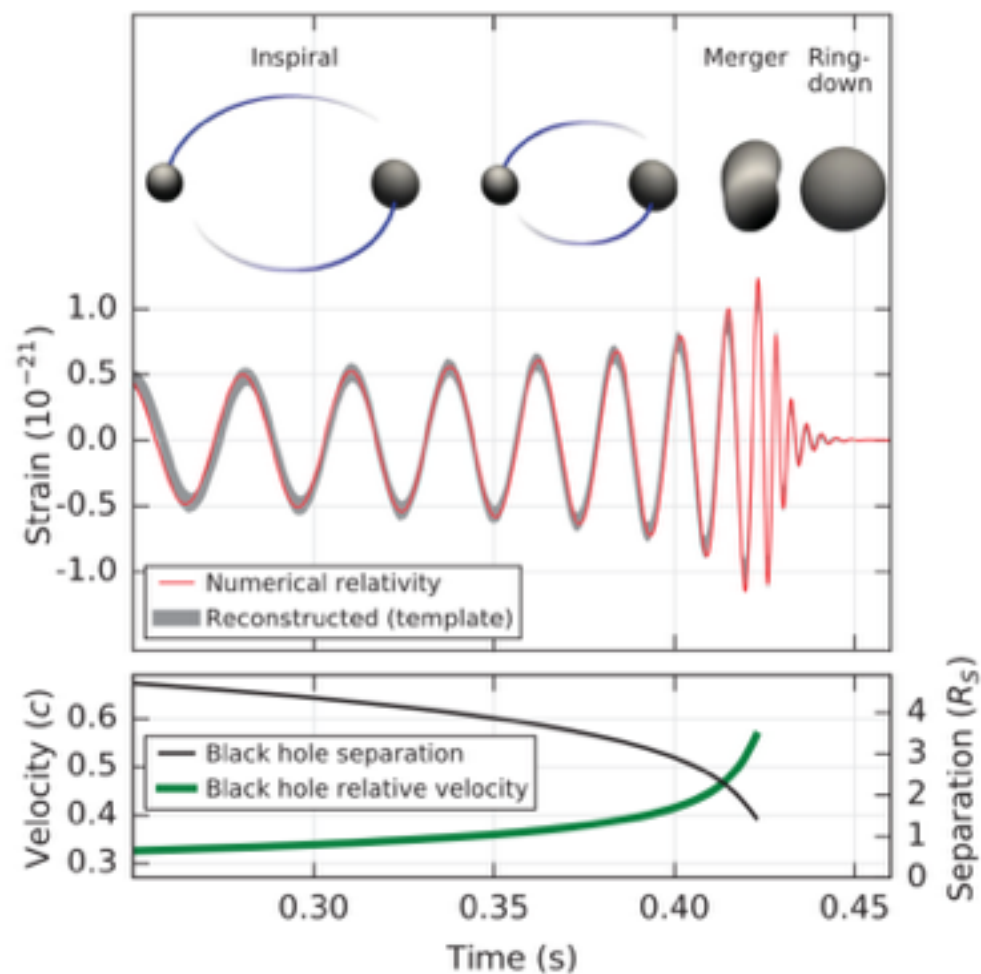
$$M \sim 2.61 M_\odot$$
$$(M_b \sim 3.2 M_\odot)$$

APR4 ✗

H4 ✗

FIRST GRAVITATIONAL WAVE DETECTION

LIGO Scientific Collaboration & Virgo Collaboration, PRL 116, 061102 (2016)



a black hole binary merger!

Phenomenology - phase I

Siegel et al. 2014

Siegel & Ciolfi 2015a

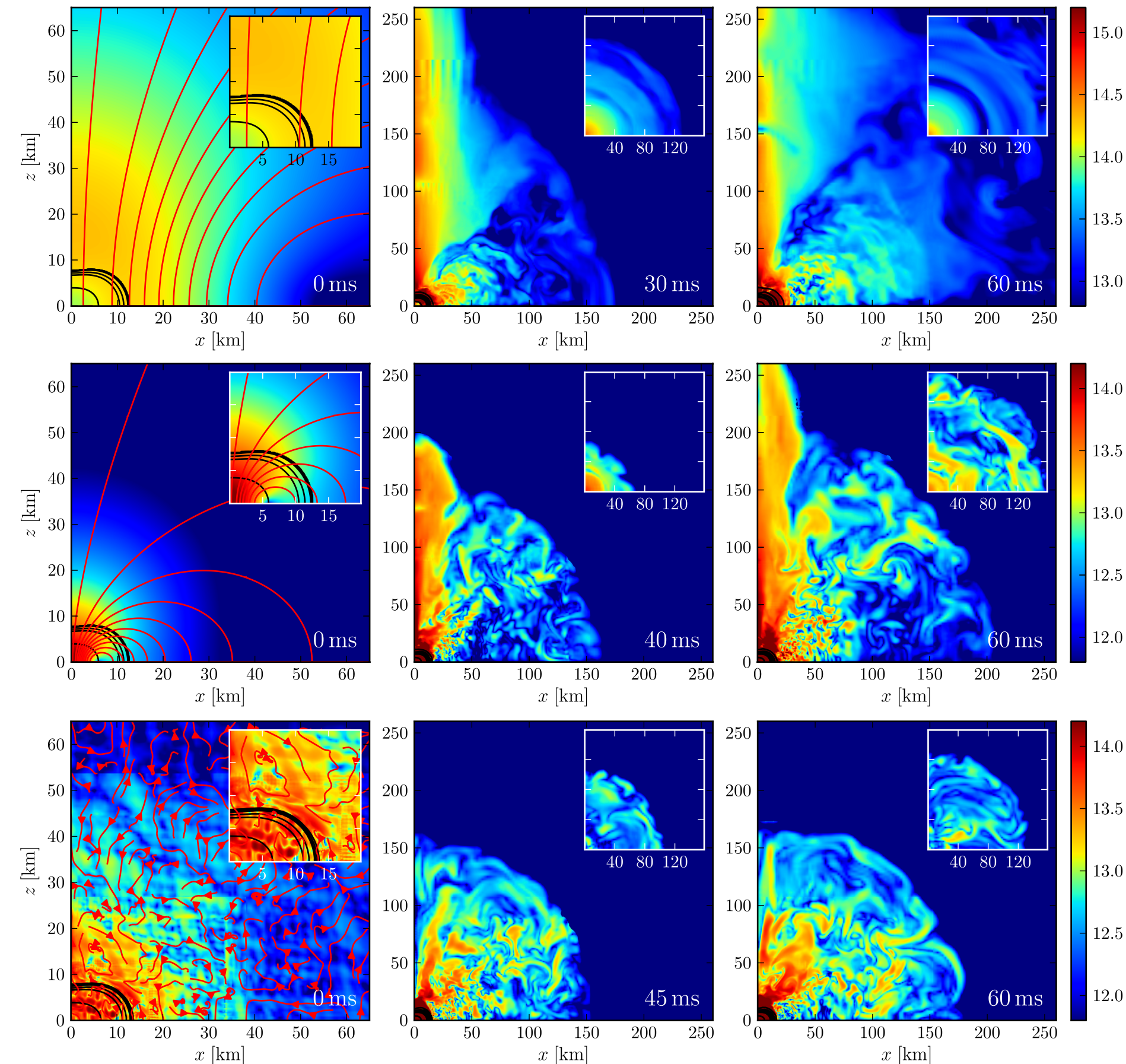
60 ms evolution
for 3 geometries

dipole 60
dipole 6
random

differential rotation
powers baryon-loaded
and magnetized outflow

for all MF geometries
the outflow has an
isotropic component

collimation depends
strongly on MF
geometry



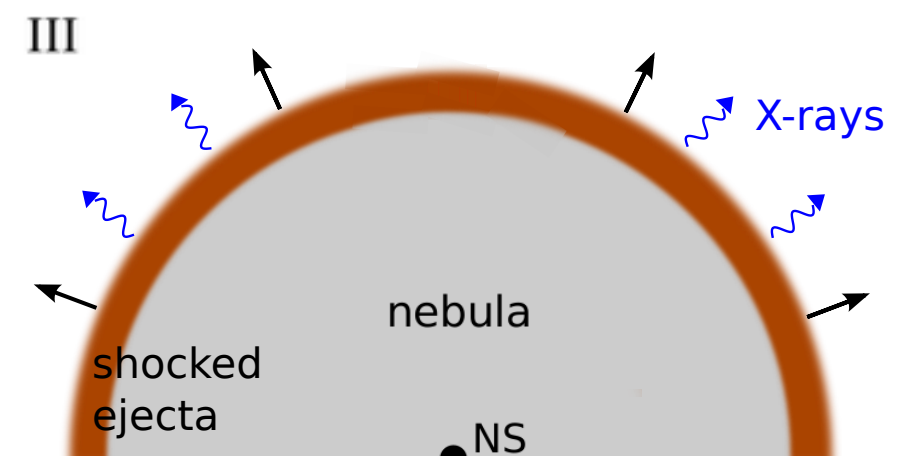
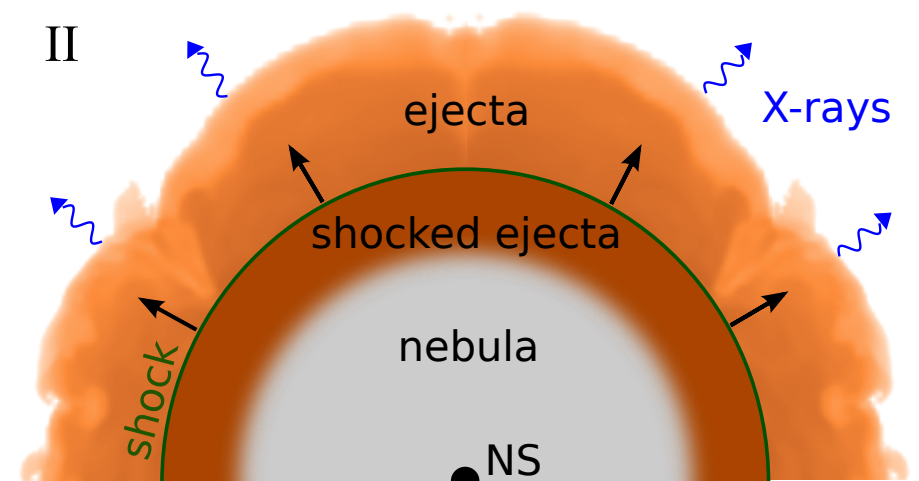
Phenomenology - phase II-III

- uniformly rotating NS emits spin-down radiation and inflates a photon-pair nebula

$$L_{\text{sd}} \simeq 1.5 \times 10^{49} B_{\text{p},15}^2 R_6^3 P_{\text{in},-3}^{-4} (1 + t/t_{\text{sd}})^{-2} \text{ erg s}^{-1}$$

$$t_{\text{sd}} \simeq 2.7 \times 10^3 B_{\text{p},15}^{-2} R_6^{-3} P_{\text{in},-3}^2 \text{ s}$$

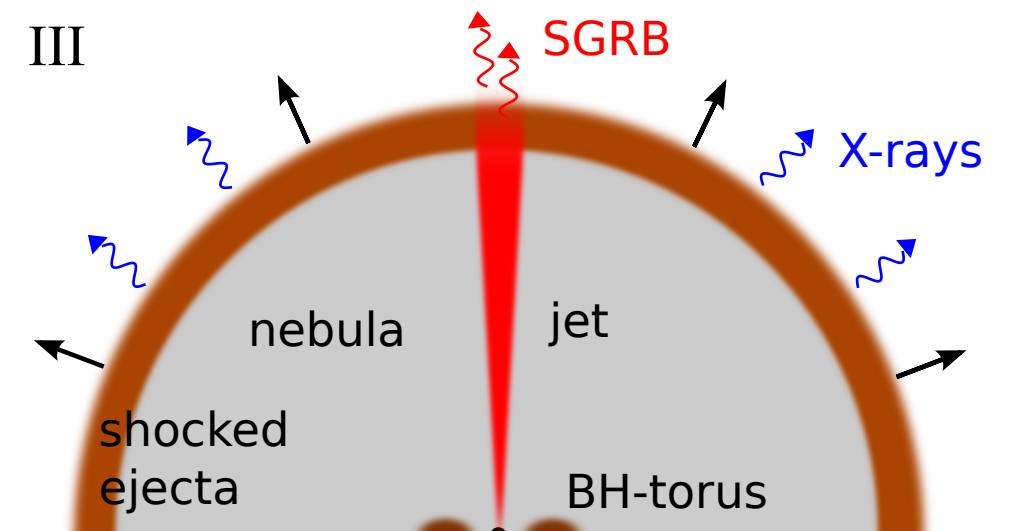
- high photon pressure drives a strong shock through the ejecta, sweeps up material into a thin shell
 - nebula energy rapidly heats up and accelerates the ejecta shell (up to mildly relativistic speeds)
- Metzger & Piro 2014



“Time-reversal” scenario

Ciolfi & Siegel 2015a, ApJ Letters 798, L36

- at $t_{\text{coll}} \sim t_{\text{sd}}$ the NS collapses to a **BH-torus system**
 - **transient jet** is formed in $\lesssim 0.01 - 1$ s
drills through the ejecta and generates the **SGRB**



- nebula and ejecta represent an **optically thick environment**
 - large fraction of **spin-down energy** is still trapped and diffuses outwards on much longer timescale
 - spin-down energy acquires **substantial delay** before emerging and producing the **X-rays**

EM emission from the long-lived NS remnant

Siegel & Ciolfi 2016a, 2016b

isotropy \rightarrow 1D model

$$\frac{dR_{ej}}{dt} = v_w(R_{ej}(t), t)$$

$$\frac{dE_{th}}{dt} = L_{EM}(t) + \frac{dE_{th,NS}}{dt} - L_{rad}(t)$$

set of coupled ODEs
for the evolution

+

balance equation for
photons and particles



$$\frac{dR_{ej}}{dt} = v_w(R_{ej}(t), t)$$

$$\frac{dR_{sh}}{dt} = v_{sh}(t)$$

$$\frac{dR_n}{dt} = \frac{dR_{sh}}{dt} - \frac{d\Delta_{sh}}{dt}$$

$$\frac{dE_{th,sh}}{dt} = \frac{dE_{sh}}{dt} + \frac{dE_{th,vol}}{dt} + \frac{dE_{PWN}}{dt} - L_{rad,in}(t)$$

$$\frac{dE_{th,ush}}{dt} = -\frac{dE_{th,vol}}{dt} - L_{rad}(t)$$

$$\frac{dE_{th}}{dt} = \frac{dE_{th,sh}}{dt} + \frac{dE_{th,ush}}{dt}$$

$$\frac{dE_{nth}}{dt} = -\frac{E_{nth}}{R_n} \frac{dR_n}{dt} - \frac{dE_{PWN}}{dt} + L_{rad,in}(t) + \eta_{TS}[L_{sd}(t) + L_{rad,pul}(t)]$$

$$\frac{dE_B}{dt} = \eta_{B_s}[L_{sd}(t) + L_{rad,pul}(t)]$$

$$0 = Q(\gamma) + P(\gamma) + \dot{N}_{C,syn}(\gamma)$$

$$\frac{dv_{ej}}{dt} = a_{ej}(t)$$

$$\frac{dR_{ej}}{dt} = v_{ej}(t) + \frac{1}{2}a_{ej}(t)dt$$

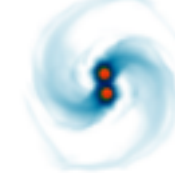
$$0 = \dot{n}_0 + \dot{n}_A + \dot{n}_C^{NT} + \dot{n}_C^T + \dot{n}_{syn} - \frac{c}{R_n}n(\Delta\tau_C^{NT} + \Delta\tau_{\gamma\gamma}) - \dot{n}_{esc}$$

$$\frac{dR_n}{dt} = \frac{dR_{ej}}{dt}$$

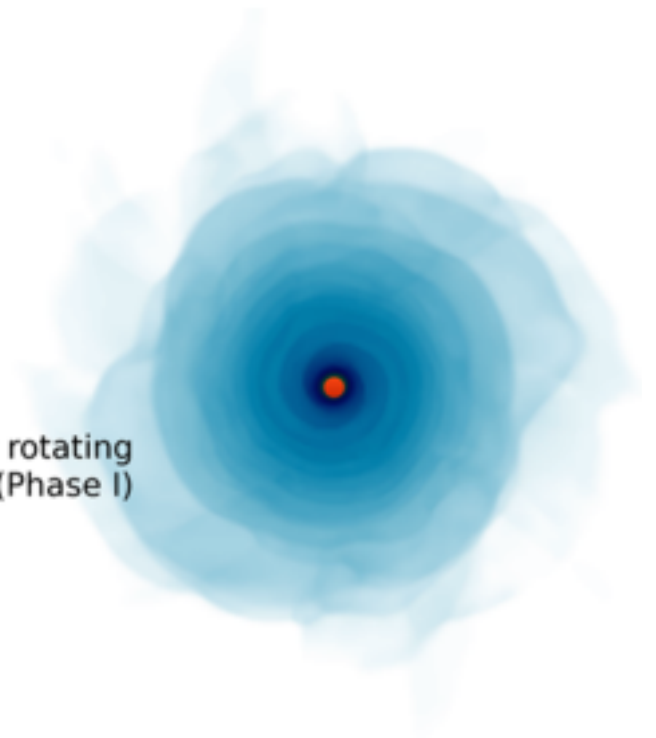
$$\frac{dE_{th}}{dt} = [1 - f_{ej}(t)] \frac{dE_{PWN}}{dt} - L_{rad}(t) - L_{rad,in}(t)$$

$$\frac{dE_B}{dt} = \eta_{B_s}[L_{sd}(t) + L_{rad,pul}(t)]$$

BNS merger



differentially rotating
NS remnant (Phase I)



shock and PWN
(Phase II-III)

