

# Goldberg-Sachs and the alignment of Einstein-Maxwell fields

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Einstein-Maxwell +  $\Lambda$ -term:

$$R_{ab} - \frac{1}{2}Rg_{ab} + \Lambda g_{ab} = F_{ac}F_b{}^c - \frac{1}{4}g_{ab}F_{cd}F^{cd}$$

- “aligned case”: at least one PND of  $\mathbf{F}$  is parallel to a Debever<sup>1</sup>-Penrose vector
- “doubly aligned case”: both real PND's of  $\mathbf{F}$  are parallel to a Debever-Penrose vector
- “non-aligned case”: no PND's of  $\mathbf{F}$  are parallel to a Debever-Penrose vector

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<sup>1</sup>Westbridge 25.11.1915–Brussels 11.5.1998

Goldberg-Sachs theorem (1962):

$$\begin{aligned} \mathcal{M} \text{ (vacuum) algebraically special} \\ \iff \\ \mathcal{M} \text{ contains a shear-free geodesic null congruence} \end{aligned}$$

generalisation :

$$\begin{aligned} k \text{ shear-free, geodesic PND of } \boldsymbol{F} \\ \implies \\ k \text{ multiple Debever-Penrose vector} \end{aligned}$$

reverse?

Kundt-Trümper (1962:)

$k$  PND of a non-null  $F$  and multiple Debever-Penrose vector

$$\implies \kappa(3\Psi_2 - 2|\Phi_1|^2) = 0 \text{ and } \sigma(3\Psi_2 + 2|\Phi_1|^2) = 0$$

hence

$$|\kappa|^2 + |\sigma|^2 \neq 0 \implies \text{type II or D with } \frac{3}{2}\Psi_2 = \pm|\Phi_1|^2$$

# Doubly aligned type D

- $\frac{3}{2}\Psi_2 \neq \pm|\Phi_1|^2 \implies \kappa = \nu = \sigma = \lambda = 0 \implies$  'class  $\mathcal{D}$ '  
(Debever-McLenaghan 1981, Plebański Demiański 1976)
- $\frac{3}{2}\Psi_2 = |\Phi_1|^2$  ( $\sigma = 0$ )  $\implies$  (Plebański-Hacyan 1979) both  $k$  and  $\ell$  are shear-free, resp. non-geodesic and geodesic and "double Kundt" ( $\rho = \mu = 0$ );  $\Lambda < 0$
- $\frac{3}{2}\Psi_2 = -|\Phi_1|^2$  ( $\kappa = 0$ )  $\implies$  (García-Plebański 1982) both  $k$  and  $\ell$  are geodesic, shearing and twisting, but non-expanding ( $\Re\rho = \Re\mu = 0$ );  $\Lambda < 0$

What happens in the non-doubly aligned case?

[VdB, arXiv:1605.05830]

# Aligned multiple D-P vector with $|\kappa|^2 + |\sigma|^2 \neq 0$

$$|\kappa|^2 + |\sigma|^2 \neq 0$$

## 1. Theorem:

multiple Debever-Penrose vector with  $\sigma \neq 0$  and PND of non-null  $\mathbf{F} \implies \Lambda < 0 \implies$  García-Plebański (this corrects and generalizes Kozarzewski[Acta Phys. Pol. 27, 775, 1965]:  $\rho = \pm i|\sigma|$  cannot easily be dismissed).

## 2. Theorem:

multiple Debever-Penrose vector with  $\kappa \neq 0$ :  
 $\rho = 0$  (Kundt sub-family)  $\implies \Lambda < 0 \implies$  Plebański-Hacyan  
 $\rho \neq 0$ : open problem!

## 3. Theorem:

$k$  multiple Debever-Penrose vector with  $\kappa = \sigma = 0$  and  
 $k$  no PND of  $F \implies F$  non-null and  $R = 0$

Choosing  $\ell$  such that  $\Phi_1 = 0$ :

- $\pi = 0 \implies$  Griffiths 1986 (II)
  - $\tau = 0 \implies$  Griffiths 1983 (II)
  - $\mu = 0 \implies$  Cahen-Spelkens 1967 (III)
  - $\rho\nu = \tau\lambda \implies$  Cahen-Leroy 1966 (N)
- $\pi \neq 0 \implies ?$

## 4. Corollary:

type D with both Debever-Penrose vectors  $k$  and  $\ell$   
geodesic and shear-free ( $\kappa = \sigma = \lambda = \nu = 0$ )

- $\Lambda \neq 0 \implies$  doubly aligned  $\implies$  class  $\mathcal{D}$
- $\Lambda = 0 \implies ?$

(Geroch, Held and Penrose, JMP 1973)

$$\mathbf{e}_1, \mathbf{e}_2, \mathbf{e}_3, \mathbf{e}_4 \equiv \mathbf{m}, \bar{\mathbf{m}}, \ell, \mathbf{k} \text{ with } \mathbf{k} \cdot \ell = -1, \mathbf{m} \cdot \bar{\mathbf{m}} = 1$$

Under boosts

$$\mathbf{k} \rightarrow A\mathbf{k}, \ell \rightarrow A^{-1}\ell$$

and spatial rotations

$$\mathbf{m} \rightarrow e^{i\theta}\mathbf{m}$$

*well-weighted* variables  $\eta$  of weight  $[p, q]$  transform as

$$\eta \rightarrow A^{\frac{p+q}{2}} e^{i\frac{p-q}{2}\theta} \eta$$

( $\eta$  has *boost-weight*  $= \frac{p+q}{2}$  and *spin-weight*  $= \frac{p-q}{2}$ ).



Basic variables:

$$\begin{aligned}\kappa &= \Gamma_{414}, & \tau &= \Gamma_{413}, & \sigma &= \Gamma_{411}, & \rho &= \Gamma_{412}, \\ \nu &= \Gamma_{233}, & \pi &= \Gamma_{234}, & \lambda &= \Gamma_{232}, & \mu &= \Gamma_{231},\end{aligned}$$

$$(\Gamma_{abc} = -\Gamma_{bac} \equiv \mathbf{e}_a \nabla_c (\mathbf{e}_b)),$$

$$\Phi_{00}, \Phi_{22}, \Phi_{01}, \Phi_{12}, \Phi_{02}, \Phi_{11},$$

$$R, \Psi_0, \Psi_1, \Psi_2, \Psi_3, \Psi_4.$$

$\alpha, \beta, \epsilon, \gamma$  get absorbed in  $\mathbb{P}, \mathbb{P}', \bar{\partial}, \bar{\partial}'$ :

$$\mathbb{P}\eta = (D - p\epsilon - q\bar{\epsilon})\eta$$

$$\mathbb{P}'\eta = (\Delta - p\gamma - q\bar{\gamma})\eta$$

$$\bar{\partial}\eta = (\delta - p\beta - q\bar{\alpha})\eta$$

$$\bar{\partial}'\eta = (\bar{\delta} - p\alpha - q\bar{\beta})\eta$$

Symmetry transformations:

- complex conjugation
- prime transformation:

$$k \leftrightarrow \ell, m \leftrightarrow \bar{m},$$

$$\kappa \leftrightarrow -\nu, \quad \tau \leftrightarrow -\pi, \quad \sigma \leftrightarrow -\lambda, \quad \rho \leftrightarrow -\mu,$$

$$\Phi_{ij} \leftrightarrow \Phi_{2-i, 2-j}, \quad \Psi_i \leftrightarrow \Psi_{4-i}$$

- Sachs transformation:

$$k \rightarrow m, \ell \rightarrow \bar{m}, m \rightarrow \bar{k}, \bar{m} \rightarrow \ell$$

$$\rho^* = \tau, \tau^* = -\rho, \rho'^* = -\tau', \tau'^* = \rho'$$

Basic equations:

- 12 complex Ricci equations

$$\mathbb{P}\tau - \mathbb{P}'\kappa = (\tau - \bar{\tau}')\rho + (\bar{\tau} - \tau')\sigma + \Phi_{01} + \Psi_1,$$

$$\bar{\partial}\rho - \bar{\partial}'\sigma = (\rho - \bar{\rho})\tau + (\bar{\rho}' - \rho')\kappa + \Phi_{01} - \Psi_1,$$

$$\mathbb{P}\sigma - \bar{\partial}\kappa = (\rho + \bar{\rho})\sigma - (\tau + \bar{\tau}')\kappa + \Psi_0,$$

$$\mathbb{P}\rho - \bar{\partial}'\kappa = \rho^2 + \sigma\bar{\sigma} - \bar{\kappa}\tau - \kappa\tau' + \Phi_{00},$$

$$\mathbb{P}'\sigma - \bar{\partial}\tau = \sigma\rho' - \bar{\lambda}\rho - \tau^2 + \kappa\bar{\nu} - \Phi_{02},$$

$$\mathbb{P}'\rho - \bar{\partial}'\tau = \rho\bar{\rho}' - \lambda\sigma - \tau\bar{\tau} + \kappa\nu - \Psi_2 - \frac{1}{12}R$$

- Maxwell equations

$$\mathbb{P}\Phi_1 - \bar{\partial}'\Phi_0 = \pi\Phi_0 + 2\rho\Phi_1 - \kappa\Phi_2$$

$$\mathbb{P}\Phi_2 - \bar{\partial}'\Phi_1 = -\lambda\Phi_0 + 2\pi\Phi_1 + \rho\Phi_2$$

- 9 complex + 2 real Bianchi equations
- commutator relations

“doubly shearfree and geodesic  $\implies$  doubly aligned”

Define extension variables  $U, V$ :

$$\mathbb{P}'\Phi_0 = U, \quad \tilde{\delta}'\Phi_0 = V$$

Maxwell equations

$$\mathbb{P}\Phi_1 = \pi\Phi_0 + 2\rho\Phi_1 + V$$

$$\tilde{\delta}\Phi_1 = \mu\Phi_0 + 2\tau\Phi_1 + U$$

“doubly shearfree and geodesic  $\implies$  doubly aligned”

GHP equations

$$\mathbb{P}\rho = \rho^2 + |\Phi_0|^2$$

$$\mathbb{P}\tau = \rho(\tau + \bar{\pi}) + \Phi_0 \overline{\Phi_1}$$

$$\mathbb{P}\mu - \bar{\partial}\pi = \pi\bar{\pi} + \mu\bar{\rho} + \frac{1}{12}R + \Psi_2$$

$$\bar{\partial}\rho = \tau(\rho - \bar{\rho}) + \Phi_0 \overline{\Phi_1}$$

$$\bar{\partial}\tau = \tau^2 + \Phi_0 \overline{\Phi_2}$$

“doubly shearfree and geodesic  $\implies$  doubly aligned”

Bianchi equations  $\implies$

$$\mathbb{P}\Psi_2 = 3\rho\Psi_2 + 2\Phi_1(\rho\overline{\Phi_1} - \tau\overline{\Phi_0}) - U\overline{\Phi_0} + V\overline{\Phi_1}$$

$$\tilde{\partial}\Psi_2 = 3\tau\Psi_2 + 2\Phi_1(\rho\overline{\Phi_2} - \tau\overline{\Phi_1}) - U\overline{\Phi_1} + V\overline{\Phi_2}$$

and

$$\overline{\Phi_1}\mathbb{P}\Phi_0 - \overline{\Phi_0}\tilde{\partial}\Phi_0 = 0$$

$$\overline{\Phi_2}\mathbb{P}\Phi_0 - \overline{\Phi_1}\tilde{\partial}\Phi_0 = 0$$

hence

$$\mathbb{P}\Phi_0 = \tilde{\partial}\Phi_0 = 0$$

“doubly shearfree and geodesic  $\implies$  doubly aligned”

Integrability conditions  $\implies$

$$\bar{\partial}V = (\rho - \bar{\rho})U - 2\Phi_0(|\Phi_1|^2 - \psi_2 + \rho\mu + \frac{1}{24}R)$$

$$\mathbb{P}V = \rho V - 2\Phi_0(\rho\pi + \Phi_1\overline{\Phi_0})$$

$$\bar{\partial}U = \tau U - 2\Phi_0(\mu\tau + \Phi_1\overline{\Phi_2})$$

$$\mathbb{P}U = (\bar{\pi} - 2\tau)V + 3\rho U - \Phi_0(2\tau\pi + 2|\Phi_1|^2 - \psi_2 + \frac{1}{6}R)$$

after which  $[\mathbb{P}', \mathbb{P}]\Phi_0$  gives

$$\Phi_0(\psi_2 - \frac{1}{12}R) + \rho U - \tau V = 0$$

the derivatives of which imply

$$\tau\Phi_0 R = \rho\Phi_0 R = 0$$