

“Gravitational” scalar-tensor theory

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introduction

Einstein's General Relativity

- “**The classical theory of fields**” by Landau & Lifshitz,
GR is a **unique theory of gravity** provided
 - composed only by metric and its derivatives
 - 4D
 - covariant theory
 - EOM is (at most) 2nd order
- have to abandon one of assumptions above (or more ?)
metric ? > 4D ? covariance ? 2nd order EOM ?

$f(R)$ theory

- The action is given by a **non-linear** function of R
- **EOM is 4th order** because $R \supset \ddot{g} + \dot{g}^2$

$$\int d^4x \sqrt{-g} f(R) \sim \int d^4x f(\ddot{g}) \rightarrow \ddot{f}(\ddot{g}) \supset \dddot{\ddot{g}} : \text{EOM}$$

- Under a **Weyl (conformal) transformation**,

$$g_{\mu\nu} \rightarrow \Omega^2 g_{\mu\nu}$$

$f(R)$ theory = Einstein + **canonical scalar**

→ change of gravity law = introduction of matter !!

correspondence

metric + ϕ (scalar-tensor)
exotic matter ?

metric (purely gravitational)
modified gravity ?

canonical scalar

$f(R)$ theory



Horndeski's theory

- the most general S-T theory **with 2nd order EOM for g & φ** :

canonical : $\mathcal{L} = X - V(\phi) = -\frac{1}{2}\nabla_\mu\phi\nabla^\mu\phi - V(\phi)$

k-essence : $\mathcal{L}_2 = K(\phi, X)$ $X = -(\nabla\phi)^2/2$

KGB : $\mathcal{L}_3 = G(\phi, X)\square\phi$

Horndeski : $\mathcal{L}_4 = G_4(\phi, X)R - \frac{\partial G_4}{\partial X}[(\nabla\nabla\phi)^2 - (\square\phi)^2]$

$$\mathcal{L}_5 = G_5(\phi, X)G_{\mu\nu}\nabla^\mu\nabla^\nu\phi - \frac{1}{3}\frac{\partial G_5}{\partial X}[(\nabla\nabla\phi)^3 + \dots]$$

(cf. GLPV theory)

dual description ??

metric + ϕ (scalar-tensor)
exotic matter ?

metric (purely gravitational)
modified gravity ?

canonical scalar

f (R) theory

k-essence

?

KGB (\supset Box)

??

Horndeski (\supset Box²)

???

Horndeski (\supset Box³)

????

⚠ a special Horndeski \leftrightarrow f (Gauss-Bonnet) ⚠

Kobayashi et.al. (2011)

question

Question

What if we introduce derivatives of R ? cf. $f(R)$

$$f\left(R, (\nabla R)^2, \square R, \dots\right)$$

What is the corresponding S-T theory, if it exists ??

model

model

- The action is given by \mathbf{R} and **derivatives of \mathbf{R}** :

$$f(R, (\nabla R)^2, \square R, \dots) \quad (\nabla R)^2 = g^{\mu\nu} \nabla_\mu R \nabla_\nu R$$

c.f. **f (R, $\square R, \dots, \square^n R$)** theory Wands (1993)

f (Riemann) theory Deruelle et.al. (2009)

- We call this model “**gravitational** scalar-tensor theory” since the theory is constructed only in terms of **gravitational language**, namely metric & its derivatives.

Ostrogradski's ghost ??

- **non-degenerate** Lagrangian with **higher derivatives**

$$L = L(q, \dot{q}, \ddot{q}) \quad (\frac{d^2L}{d\dot{q}^2} \neq 0)$$

$$H = P_1 Q_2 + P_2 f(Q_1, Q_2, P_2) - L(Q_1, Q_2, P_2) \\ (Q_1 = q, Q_2 = \dot{q})$$

→ **Hamiltonian is unbounded below**

- Although $f(R, (\nabla R)^2, \square R, \dots) \supset f(g, \dot{g}, \ddot{g}, \dots)$
there is **no Ostrogradsky's ghost** in our model
because the Lagrangian is **degenerate**.

proof of healthiness

- consider a simple model : $f(R, (\nabla R)^2)$
 - rewrite \mathbf{R} with ϕ & λ : $\text{!! } \lambda \neq f_\phi \text{ !!}$
$$f(R, (\nabla R)^2) = f(\phi, (\nabla \phi)^2) - \lambda(\phi - R)$$
 - conformal tr. :
$$\frac{1}{2} \tilde{R} + f(\phi, 2\lambda (\tilde{\nabla} \phi)^2) - \frac{1}{2} (\tilde{\nabla} \chi)^2 \quad g^{\mu\nu} = 2\lambda \tilde{g}^{\mu\nu}$$
$$\lambda \equiv e^{\sqrt{2/3}} \chi$$
- 👉 # of d.o.f.s : **2 (GW) + 2 (scalar)** ($\neq \mathbf{2 + 1}$ in $f(R)$)
- 👉 sub class of bi-scalar k-essence \supset **healthy** domain

some remarks

- $(\square R)^2$ is **not allowed** in general -> **Ostrogradski's ghost** among models with the form of $f(R, (\nabla R)^2, \square R)$

$$\rightarrow K(R, (\nabla R)^2) + G(R, (\nabla R)^2) \times \square R$$

- $(\square R)^2$ **can be introduced** “with specific combinations”

$$Q(R, (\nabla R)^2) R + Q_X(R, (\nabla R)^2) [(\square R)^2 - (\nabla_\mu \nabla_\nu R)^2]$$

- Way to construct a healthy gravitational S-T theory :

$$\mathcal{L}^{\text{Horn}}(\phi, g_{\mu\nu}) \rightarrow \mathcal{L}^{\text{Horn}}(R, g_{\mu\nu})$$

dual description ??

metric + ϕ (scalar-tensor)
exotic matter ?

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modified gravity ?

canonical scalar

f (R) theory

k-essence

?

KGB (\supset Box)

??

Horndeski (\supset Box²)

???

Horndeski (\supset Box³)

????

to be continued ...

summary

summary

- We have considered a theory of gravity in which the action is given by R and derivatives of R .
- Despite the higher derivative nature of action, the theory is healthy (if f is properly chosen)
= no ghost & no Ostrogradsky's instabilities
- # of d.o.f.s = 2 (GW) + 2 (scalar)
 \Leftrightarrow 2 scalars-tensor theory \supset bi-scalar k-essence
- Way to construct a healthy gravitational S-T theory :
$$\mathcal{L}^{\text{Horn}}(\phi, g_{\mu\nu}) \rightarrow \mathcal{L}^{\text{Horn}}(R, g_{\mu\nu})$$

Thank you very much
for your attention !!

ありがとう (Arigatou)

how to construct GST

Horndeski (# : **2 + 1**)

k-essence : $K(\phi, (\nabla\phi)^2)$

KGB : $G(\phi, X) \square\phi$

Horn. (L4) :

$$G_4 R - G_{4,X} \left[(\nabla\nabla\phi)^2 - (\square\phi)^2 \right]$$

Horn. (L5) :

$$G_5 G_{\mu\nu} \nabla^\mu \nabla^\nu \phi + \dots$$

Gravitational S-T (# : **2 + 2**)

$K(R, (\nabla R)^2)$

$G(R, X_R) \square R$

$$G_4 R - G_{4,X} \left[(\nabla\nabla R)^2 - (\square R)^2 \right]$$

$$G_5 \dots + \dots$$

how to construct GST

Horndeski (**2 + 1**)

gravitational S-T (**2 + 2**)

$$\text{Horn : } \mathcal{L}_5 = G_5(\phi, X) G_{\mu\nu} \nabla^\mu \nabla^\nu \phi - \frac{1}{3} \frac{\partial G_5}{\partial X} [(\nabla \nabla \phi)^3 + \dots]$$



$$\begin{aligned} \mathcal{L}_5 \supset & f(R, (\nabla R)^2) \left[\nabla^\alpha (\nabla_\alpha \square - \square \nabla_\alpha) R - \frac{1}{2} (\nabla R)^2 - \frac{1}{2} R \square R \right] \\ & + \frac{1}{3} \frac{\partial f}{\partial (\nabla R)^2} \left(R, (\nabla R)^2 \right) \left[(\square R)^3 - 3(\square R)(\nabla_\mu \nabla_\nu R)^2 + 2(\nabla_\mu \nabla_\nu R)^3 \right] \end{aligned}$$

$$\text{identity : } G_{\mu\nu} \nabla^\mu \nabla^\nu \psi = \nabla(\nabla \square - \square \nabla) \psi - \frac{1}{2} \nabla R \nabla \psi - \frac{1}{2} R \square \psi$$