

The Kaluza Ansatz in Eddington inspired Born-Infeld Gravity

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Motivation

- Geometric deduction of EM coupling with gravity
- Non-minimally coupled actions in higher derivative theories.
- EiBI: Metric-affine theory; Determinantal action
- Infinite number of higher derivative terms.
- Smooths out singularities

The Eddington inspired Born-Infeld action

$$S_{EiBI} = \frac{1}{8\pi\kappa} \int d^4x \left[\sqrt{|\mathbf{g} + \kappa \mathbf{R}(\Gamma)|} - \lambda \sqrt{|g|} \right] + S_M$$

- Einstein-Hilbert action with c.c. and matter when

i) $\kappa R_{\mu\nu} \ll g_{\mu\nu}$, ii) $\lambda = \kappa\Lambda + 1$

- $\tilde{q}_{\mu\nu} = g_{\mu\nu} + \kappa R(\Gamma)_{\mu\nu}$

$$\sqrt{|\tilde{q}|} \tilde{q}^{\mu\nu} - (\kappa\Lambda + 1) \sqrt{|g|} g^{\mu\nu} = -8\pi\kappa \sqrt{|g|} T^{\mu\nu} \text{ (metric),}$$

$$\nabla_\alpha (\sqrt{|\tilde{q}|} \tilde{q}^{\mu\nu}) = 0 \text{ (connection).}$$

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- $R(\Gamma)_{\mu\nu} = \tau_{\mu\nu} + 64\pi^2\kappa \left[S_{\mu\nu} - \frac{1}{4} S g_{\mu\nu} \right] + \mathcal{O}(\kappa^2);$

where $S_{\mu\nu} = T_{\mu\alpha} T^\alpha_\nu - \frac{1}{2} T T_{\mu\nu}, \quad \tau_{\mu\nu} = \Lambda g_{\mu\nu} + 8\pi \left[T_{\mu\nu} - \frac{1}{2} T g_{\mu\nu} \right].$

- $\Gamma^\alpha_{\beta\gamma} = \left\{ \begin{smallmatrix} \alpha \\ \beta\gamma \end{smallmatrix} \right\} + \frac{1}{2} \kappa q^{\alpha\delta} (R(\Gamma)_{\delta\beta;\gamma} + R(\Gamma)_{\gamma\delta;\beta} - R(\Gamma)_{\beta\gamma;\delta})$

$$\Rightarrow R(g)_{\mu\nu} = \tau_{\mu\nu} + 64\pi^2\kappa \left[S_{\mu\nu} - \frac{1}{4} S g_{\mu\nu} \right] + \frac{1}{2} \kappa \left[\nabla_\mu \nabla_\nu \tau - 2 \nabla^\alpha \nabla_{(\mu} \tau_{\nu)\alpha} + \square \tau_{\mu\nu} \right] + \mathcal{O}(\kappa^2);$$

The Eddington inspired Born-Infeld action : Metric

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- $R_{\mu\nu} = \tau_{\mu\nu} + 64\pi^2\kappa \left(S_{\mu\nu} - \frac{1}{4}Sg_{\mu\nu} \right) + \frac{1}{2}\kappa \left[2R_{\alpha\mu\beta\nu}R^{\alpha\beta} - 2R_{\mu\beta}R_{\nu}^{\beta} + \square R_{\mu\nu} \right] + \mathcal{O}(\kappa^2).$
- $R(\Gamma)_{\mu\nu} = R(g)_{\mu\nu}$ (to lowest order)

The metric-affine theory and the metric theory agree to $\mathcal{O}(\kappa)$.

- Involves third order derivatives of matter fields (and above)

\Rightarrow Surface singularities [Pani,Sotiriou (2012)]

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The Kaluza Ansatz

Geometrically determines the coupling to gravity

- $\hat{g}_{AB} = \begin{pmatrix} g_{\mu\nu} + \alpha^2 A_\mu A_\nu & \alpha A_\mu \\ \alpha A_\nu & 1 \end{pmatrix}, \quad A, B, \dots = 0, 1, \dots, 4; \quad \sqrt{|\hat{g}|} = \sqrt{|g|}$

- Quantities independent of the 5th dimension.

- Need $\hat{q}_{AB} = \hat{g}_{AB} + \kappa \hat{R}_{AB}$. $(F^2 = F_{\mu\nu} F^{\mu\nu})$

$$\hat{q}_{\mu\nu} = g_{\mu\nu} + 4A_\mu A_\nu (1 + \kappa F^2) + \kappa [R_{\mu\nu} - 4A_{(\mu} \nabla_{|\beta|} F^{\beta}_{\nu)} - 2F_{\beta\mu} F^{\beta}_{\nu}],$$

$$\hat{q}_{5\nu} = 2A_\nu (1 + \kappa F^2) - \kappa \nabla_\beta F^{\beta}_{\nu},$$

$$\hat{q}_{55} = 1 + \kappa F^2.$$

The Kaluza Ansatz : EiBI action

- $$\det \begin{pmatrix} A & B \\ C & D \end{pmatrix} = \det(D) \det(A - BD^{-1}C)$$
- $$|\hat{q}| = q_{55} [q_{\mu\nu} - q_{\mu 5} (q_{55})^{-1} q_{5\nu}]$$

$$= \left[(1 + \kappa F^2) \left(g_{\mu\nu} + \kappa \left(R_{\mu\nu} + 2F_{\mu\beta} F^\beta_\nu \right) \right) - \kappa^2 \nabla_\delta F^\delta_\mu \nabla_\beta F^\beta_\nu \right] .$$
- $$\hat{S}_{EiBI} = \frac{1}{8\pi \hat{G}_5 \kappa} \int d^5x \left[\sqrt{|\hat{q}|} - (\kappa\Lambda + 1) \sqrt{|\hat{g}|} \right]$$
- $$S_{EiBI+EM} =$$

$$\frac{1}{8\pi\kappa} \int d^4x \left[\sqrt{\left| (1 + \kappa F^2) \left(g_{\mu\nu} + \kappa \left(R_{\mu\nu} + 2F_{\mu\beta} F^\beta_\nu \right) \right) - \kappa^2 \nabla_\delta F^\delta_\mu \nabla_\beta F^\beta_\nu \right|} \right.$$

$$\left. - (\kappa\Lambda + 1) \sqrt{|g|} \right] .$$

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$$\left. - (\kappa\Lambda + 1) \sqrt{|g|} \right] .$$

Equations of motion

- $\mathcal{O}(\kappa^0)$: $G_{\mu\nu} = -\Lambda g_{\mu\nu} + 8\pi T_{\mu\nu}$, $\nabla_\alpha F^{\alpha\nu} = 0$.
- $\mathcal{O}(\kappa)$: $G_{\mu\nu} = -\Lambda g_{\mu\nu} + 8\pi T_{\mu\nu} + \kappa P_{\mu\nu} + \kappa Q_{\mu\nu}$, where

$$\begin{aligned}
 P_{\mu\nu} = & R_{\mu\alpha} R_\nu^\alpha - \frac{1}{2} R R_{\mu\nu} - \frac{1}{4} R_{\alpha\beta} R^{\alpha\beta} g_{\mu\nu} + \frac{1}{8} R^2 g_{\mu\nu} + \frac{1}{2} \nabla_\mu \nabla_\nu R \\
 & - \frac{1}{2} g_{\mu\nu} \square R - \nabla_\alpha \nabla_{(\mu} R_{\nu)}^\alpha + \frac{1}{2} \square R_{\mu\nu} + \frac{1}{2} g_{\mu\nu} \nabla_\alpha \nabla_\beta R^{\alpha\beta} , \\
 Q_{\mu\nu} = & R F_{\mu\alpha} F_\nu^\alpha + \nabla_\alpha F_\nu^\alpha \nabla_\beta F_\mu^\beta + 4 R_{(\mu|\alpha|} F^{\alpha\beta} F_{|\beta|\nu)} + 2 F_{\mu\alpha} R^{\alpha\beta} F_{\beta\nu} \\
 & + 8 F_{\mu\alpha} F^{\alpha\beta} F_{\beta\gamma} F_\nu^\gamma - 2 \nabla_{(\mu} (F_{\nu)\beta} \nabla_\alpha F^{\alpha\beta}) - \frac{1}{8} F^4 g_{\mu\nu} - \frac{1}{4} g_{\mu\nu} R F^2 \\
 & - 2 \nabla_\alpha (F_\nu^\alpha \nabla_{|\beta|} F_\nu^\beta) + F^2 F_{\mu\beta} F_\nu^\beta + 2 (\nabla_{(\mu} F_{\nu)\beta}) \nabla_\alpha F^{\alpha\beta} \\
 & - 2 \nabla_\alpha \nabla_{(\mu} (F_{\nu)\beta} F^{\beta\alpha}) + \square (F_{\mu\beta} F_\nu^\beta) + g_{\mu\nu} \nabla_\alpha \nabla_\beta (F^{\alpha\gamma} F_\gamma^\beta) \\
 & - \frac{1}{2} \nabla_\mu \nabla_\nu F^2 + \frac{1}{2} g_{\mu\nu} \square F^2 + \frac{1}{2} F^2 R_{\mu\nu} - F_{\alpha\beta} F^{\beta\gamma} F_{\gamma\delta} F^{\delta\alpha} g_{\mu\nu} \\
 & - F_{\alpha\beta} F^{\beta\gamma} R_\gamma^\alpha g_{\mu\nu} + \nabla_\alpha (F_\beta^\alpha \nabla_\gamma F^{\gamma\beta}) g_{\mu\nu} - \frac{1}{2} (\nabla_\alpha F_\beta^\alpha) \nabla_\gamma F^{\gamma\beta} g_{\mu\nu}
 \end{aligned}$$

Ricci scalar and Electromagnetic EOM

$$R = 4\Lambda - \kappa \left[\frac{1}{2}F^4 + \frac{1}{2}RF^2 + 2R^{\alpha\beta}F_{\beta\gamma}F_{\alpha}^{\gamma} + 4F^{\alpha\beta}F_{\beta\gamma}F^{\gamma\delta}F_{\delta\alpha} \right. \\ \left. + \nabla_{\alpha}\nabla_{\beta}\left(R^{\alpha\beta} + 2F^{\alpha\delta}F_{\delta}^{\beta}\right) - \square\left(R - \frac{1}{2}F^2\right) + \nabla_{\beta}F^{\beta\gamma}\nabla^{\alpha}F_{\alpha\gamma} \right] .$$

$$\nabla_{\alpha}F^{\alpha\nu} = -\kappa \left[\nabla_{\alpha}\left(F^{\alpha\nu}\left(\frac{1}{2}R + \frac{1}{2}F^2\right)\right) - 4\nabla_{\alpha}\left(F^{\alpha\beta}F_{\beta\gamma}F^{\gamma\nu}\right) - \frac{1}{2}\square\left(\nabla_{\alpha}F^{\alpha\nu}\right) \right. \\ \left. - \nabla_{\alpha}\left(R^{\alpha\beta}F_{\beta}^{\nu} - R^{\nu\beta}F_{\beta}^{\alpha}\right) + \frac{1}{2}\nabla_{\alpha}\nabla^{\nu}\left(\nabla_{\beta}F^{\beta\alpha}\right) \right] .$$

- Regardless of appearances, the solutions are rather simple.

Solutions : Iterative procedure

$$\Lambda g_{\mu\nu} = -G_{\mu\nu} + 8\pi T_{\mu\nu} - \kappa C_{\mu\nu},$$

$$\nabla_\mu F^{\mu\nu} = -\kappa D^\nu.$$

- Solve via $g_{\mu\nu} = g_{\mu\nu}^0 + g_{\mu\nu}^1$ and $A_\mu = A_\mu^0 + A_\mu^1$
- $g_{\mu\nu}^0 = (f(r), f(r)^{-1}, r^2, r^2 \sin^2\theta); \quad f(r) = 1 - \frac{2m}{r} - \frac{\Lambda r^2}{3} + \frac{q^2}{r^2}$

$$A_\mu^0 = \left(\frac{q}{r}, 0, 0, 0\right)$$

- Next order can now be solved :

$$\Lambda g_{\mu\nu}^1 + G_{\mu\nu}^1 - 8\pi T_{\mu\nu}^1 = -\kappa C_{\mu\nu},$$

$$\nabla_\mu^0 F^{1\mu\nu} + \nabla_\mu^1 F^{0\mu\nu} = -\kappa D^\nu,$$

Solutions : Iterative procedure

$$\kappa C_{\mu\nu} = \begin{pmatrix} -f(r)\alpha(r) & 0 & 0 & 0 \\ 0 & \alpha(r)(f(r))^{-1} & 0 & 0 \\ 0 & 0 & r^2\beta(r) & 0 \\ 0 & 0 & 0 & r^2\sin^2\theta\beta(r) \end{pmatrix},$$

$$\kappa D^\nu = \left(\frac{12\kappa q^3}{r^7}, 0, 0, 0 \right);$$

$$\text{where } \alpha(r) = -\frac{9\kappa q^4}{2r^8} - \frac{\kappa\Lambda q^2}{r^4} \text{ and } \beta(r) = \frac{3\kappa q^4}{2\Lambda r^8} + \frac{\kappa q^2}{r^4}$$

- General solution for $g_{\mu\nu}$ and A_μ

$$g_{\mu\nu} = (-(f(r) + \kappa G(r)), (f(r) + \kappa G(r))^{-1}, 0, 0)$$

$$A_\mu = \left(\frac{q}{r} + \kappa B(r), 0, 0, 0 \right)$$

Solutions : Iterative procedure

- RNdS correction:

$$G(r) = \frac{3}{10} \frac{q^4}{r^6} - \frac{\Lambda}{r^2} \quad , \quad B(r) = \frac{3q^3}{5r^5}$$

- Trace : $R = 4\Lambda + \frac{6\kappa q^4}{r^8}$; finite for all finite r .
- No surface singularities present in the derived solution.
- Procedure can be repeated for higher order solutions.

Conclusions

- Derived a 4-d action of EM coupled to EiBI.
- Iterative procedure for solutions about a fixed background
 $\rightarrow \mathcal{O}(\kappa)$ correction of Reissner-Nördstrom de Sitter.
- Questions remain : Stability of solutions
EM wave propagation
Behaviour of geodesics.