

Gravitational lensing beyond geometric optics

Abraham Harte
(with Yi-Zen Chu)

Max Planck Institute for Gravitational Physics
Albert Einstein Institute
Potsdam, Germany

July 13, 2016

GR 21
Columbia University

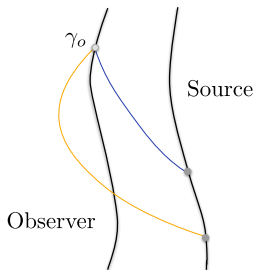
1 Geometric optics

2 Generalizing: Why and how?

3 Some consequences

Traditional gravitational lensing

- Light, high-freq. GWs, etc. move on null geodesics,



- Intensity follows from cross-sectional areas of null geodesic bundles.
- Polarization is parallel-transported, ...

Maxwell theory doesn't look anything like this:

$$dF = 0, \quad d \star F = 4\pi J.$$

How does one go from here to the rules of geometric optics?

Deriving geometric optics

Consider a 1-parameter family of fields [Ehlers (1967), ...]

$$F(x; \omega) \sim e^{i\omega S(x)} \sum_{n=0}^{\infty} \frac{f_n(x)}{\omega^n}.$$

Substituting into Maxwell's equations while using $\omega \rightarrow \infty$ gives

- Momentum fluxes point along $k_a := \nabla_a S$, and this is null.
- Amplitude transported along null geodesics: $2k \cdot \nabla f_0 + f_0 \nabla \cdot k = 0$.
- Amplitudes evolve via cross-sectional areas: $\nabla_a (|f_0|^2 k^a) = 0$.

When is geometric optics sufficient?

If there's one lengthscale r , fractional corrections look like

$$1 + \frac{(\dots)}{\omega r} + \frac{(\dots)}{(\omega r)^2} + \dots$$

So everything is fine if $\omega r \gg 1$ (i.e. $r \gg \lambda$).

Two lengthscales?

More dimensionless combinations are possible:

$$1 + \left[\frac{(\dots)}{\omega r} + \frac{(\dots)}{\omega l} + \frac{(\dots)r}{\omega l^2} + \dots \right] + \dots$$

Could some effects grow with distance r ?

Why generalize geometric optics?

- 1 Wave effects might accumulate over large distances.
- 2 Learn something using high precision measurements.
- 3 Caustics.

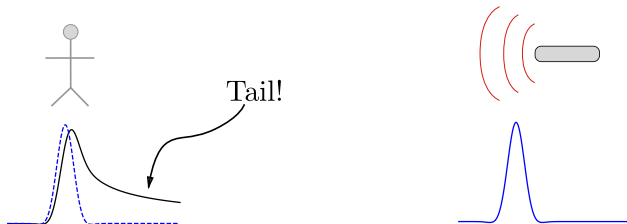
Straightforward to extend high-frequency ansatz to higher orders in ω^{-1}
[Anile (1976), ...]

But:

- 1 Is this ansatz appropriate? Too special?
- 2 Too general? Maybe equations allow too much.
- 3 How are field properties related to source properties?

Signaling experiments

Intervening spacetime distorts signals:



(output) = $\mathcal{T}[(\text{input})]$ with \mathcal{T} **nonlocal** and usually **linear**.

Find “transfer function” $\mathcal{T}[\cdot]$

Method: Green functions

For an EM field, there's some $G_{aa'}(x, x')$ such that

$$\nabla^b \nabla_b G_{aa'} - R_a{}^b G_{ba'} = -4\pi \delta(x, x').$$

Measured fields then look like

$$F_{ab}[J] = 2 \int \nabla_{[a} G_{b]} J^{b'} dV'$$

Lots known about Green functions. Can use well-developed bitensor results
[Hadamard, Friedlander, DeWitt, Poisson, ...]

Specialize for sources satisfying

$$J^a(x) = \int d\tau [\cancel{q\dot{z}^a} + \overset{\text{boring}}{\mu^{ab}(\tau)} \nabla_b + \dots] \delta(\gamma_s(\tau), x)$$

- Radiation depends on dipole moment $\mu^{ab} = \mu^{[ab]}(\tau)$.
- Find $F_{ab} = F_{ab}[\mu^{cd}]$ as μ^{cd} varies rapidly. . .

$$F_{ab}(x) = (\cdots) - 2 \int \nabla_{a'} \nabla_{[a} G_{b]b'}(z(\tau), x) \mu^{a'b'}(\tau) d\tau$$

So if $\mu(\tau) \sim e^{i\omega\tau}$,

$$F_{ab}(x) \sim \mathcal{F}[\nabla\nabla' G](\omega)$$

Everything reduces to **Fourier transforms** of $G(z(\tau), x)$ **wrt** τ .

High-frequency limits

Geometric optics should be recovered as $\omega \rightarrow \infty$.

\hookrightarrow Look at high- ω behavior of Fourier transforms.

“Almost-geometric” optics

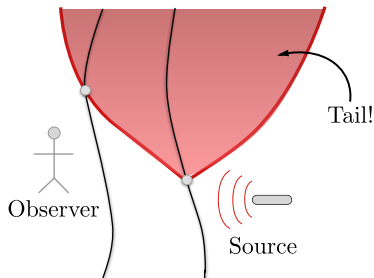
[High-frequency signals] \leftrightarrow [Most singular bits of $G(z(\cdot), x)$]
[Riemann-Lebesgue lemma, Paley-Wiener thms, etc.]

An immediate conclusion

Singularities of $G(z(\cdot), \gamma_o)$ concentrated on null geodesics from $\gamma_o \rightarrow \mathcal{S}$
 \hookrightarrow Geometric optics at leading order

All $O(\omega^{-n})$ effects come from neighborhood of observer's past light cone

What about tails?



Smooth tails away from null geodesics:

$\hookrightarrow \mathcal{F}[(\text{smooth})](\omega)$ decays faster than any polynomial, so
tails give nonperturbative corrections.

Detailed calculation gives

- (Polarization-dependent) phase shifts
- Frequency-dependent source directions
- Intensity shifts
- ...

All in terms of Hadamard-type transport equations along null geodesics.

Conclusions

- ① Framework to describe gravitational lensing beyond geometric optics
- ② Connects to bitensor constructions used in other contexts
- ③ Tail effects are “exponentially suppressed” away from null geodesics

Work in progress: Find interesting applications!